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COMPUTER PROGRAMS
IN
EARTHQUAKE SEISMOLOGY
VOLUME 2: SURFACE WAVE PROGRAM

EDITED BY
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<p>16. Abstract (Limit: 200 words)</p> <p>This volume contains twelve documented copies of computer programs dealing with surface wave propagation in a plane layered medium, to be used in earthquake seismology. The programs have been designed so that the output of one is compatible with the input of another program. These programs are presently operating on a Honeywell 6023 and are free of machine dependent peculiarities, with a few exceptions listed in the introduction. Although most programs require an offline CALCOMP drum plotter, in more recent programs plotting calls are performed in modular subroutine units to facilitate conversion to other plotting systems. The CALCOMP subroutines are listed. Most of the programs consist of a statement of purpose, input/output, program units, description or theory, the actual program, sample data, and sample output. The program titles are: SURFACE, REIGEN, LEIGEN, EIGEN, MODEL, NONPRP, EXCIT, SRFWVPLT, SRFWVSRC, RADPAT, QUESTION, and WIGGLE.</p>																													
<p>17. Document Analysis a. Descriptors</p> <table> <tr> <td>Surface waves</td> <td>Earthquakes</td> <td>Programming</td> </tr> <tr> <td>Seismic waves</td> <td>Computer applications</td> <td>Data processing</td> </tr> <tr> <td>Seismology</td> <td>Computer software</td> <td></td> </tr> </table> <p>b. Identifiers/Open-Ended Terms</p> <table> <tr> <td>SURFACE</td> <td>MODEL</td> <td>SRFWVSRC</td> <td>CALCOMP</td> </tr> <tr> <td>REIGEN</td> <td>NONPRP</td> <td>RADPAT</td> <td></td> </tr> <tr> <td>LEIGEN</td> <td>EXCIT</td> <td>QUESTION</td> <td></td> </tr> <tr> <td>EIGEN</td> <td>SRFWVPLT</td> <td>WIGGLE</td> <td></td> </tr> </table> <p>c. COSATI Field/Group</p>					Surface waves	Earthquakes	Programming	Seismic waves	Computer applications	Data processing	Seismology	Computer software		SURFACE	MODEL	SRFWVSRC	CALCOMP	REIGEN	NONPRP	RADPAT		LEIGEN	EXCIT	QUESTION		EIGEN	SRFWVPLT	WIGGLE	
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Introduction to Volume 2

The purpose of this volume is to provide documented copies of computer programs which are useful in earthquake seismology. This volume contains programs dealing with surface wave propagation in a plane layered medium. As with any computer program, there is always room for improvement. The program REIGEN could be made more accurate by using compound matrix techniques for computing Rayleigh wave eigenfunctions. The program WIGGLE could be modified slightly so that a Q model could also be entered. The programs have been designed so that the output of one is compatible with the input of another program. To help understand the theory used, a translation of a Russian article is included in Appendix A.

The author has had programming experience with an IBM 1620, a CDC 3300, a CDC 6400, and most recently on a Honeywell 6023 on which these programs are presently operating. The programs should be free of machine dependent peculiarities, with a few exceptions:

1. there are six alphanumeric characters to a word;
2. a floating or fixed point number occupies one word;
3. the largest and smallest floating point numbers on the Honeywell 6023 are about 1.0 E + 39 and 1.0 E - 39 respectively;
4. there is no PROGRAM statement;
5. the arcsine and arccosine functions are ARSIN and ARCOS, respectively;
6. multiple files can be used for input or output.

Many programs contain calls for an offline CALCOMP drum plotter. In the more recent programs, plotting calls are performed in modular subroutine units so that conversion to other plotting routines can be facilitated. The CALCOMP subroutines used are as follow:

PLOTS	initiates plot tape
PLOT	move pen
FACTOR	scales entire plot
SYMBOL	plots character string and special symbols
NUMBER	plots decimal equivalent of a floating point number
SCALE	determines starting value and scale for an array of data to be plotted on a graph
AXIS	draws an annotated axis line for a graph
LINE	scales and plots a set of data points defined by X and Y coordinate arrays.

The author is interested in any comments, corrections or improvements to the programs. The computer programs hopefully will be of use to other researchers. An acknowledgment of their source is all that is requested.

I. SURFACE

PROGRAMMER: R. B. HERRMANN / Nov 1971

PURPOSE:

This program solves the Rayleigh and Love wave period equations in order to find the dispersion curves for phase and group velocity, the surface wave amplitude factors, and the Rayleigh wave ellipticity. The phase velocity is determined to an accuracy of 0.00000001. One liquid layer is allowed at the surface.

INPUT/OUTPUT

No tapes are required. Input is on card. Output is on printer with option for punched card output of dispersion curve data.

PROGRAM DESCRIPTION

PROGRAM SURFACE: This is the main control link. The model and control cards are read in here.

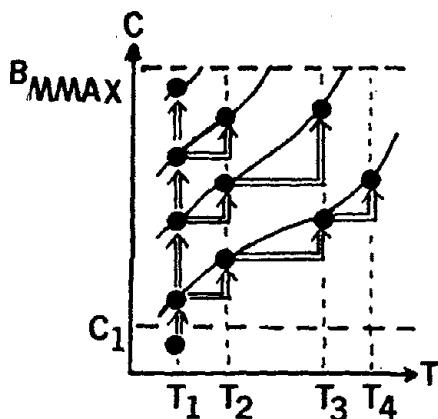
SUBROUTINE GPHDIS: This subroutine evaluates the Love and/or Rayleigh wave period equations and plots the sign of the period equation. The true phase velocity is the curve for which the period equation is zero, e.g., the curve between the positive and negative values of the period equation. The ordinate

C2	-5-5-5 5 5 5-6-6 6 -5-5-5 5 5-5-6-6 6 -5-5-5 5 5-5-5 6 6 -4-5 5 5-5-5 5 6 6 -4-5 5 5-5-5 5 5 5 -4 4 4-4-4 4 4 4 4 4 4-4-4 4 4 4 4 4 4-4-4 4 4 4 4 4 4	is a set of phase velocities between C1 and C2 in increments of DC. This plot is helpful in determining which mode a particular solution of the period equation belongs. This is important because the modes might lie very close together and the parameter DC inputted into SUBROUTINE DISPER may be too large, in
C1	-4-4 4 4 4 4 4 4	

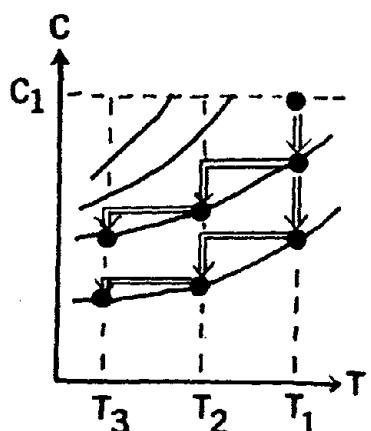
which case a particular mode may be jumped over in the root finding process. The plot also has a number to the right of the sign of the period equations. This number indicates the number of layers used in the root finding process. The number is the units value of the number of layers used (e.g., if 23 layers were used and the sign of the period equation was negative, then -3 would be plotted). A layer dropping procedure is used to increase the speed of computation and to decrease roundoff error. The abscissa of the graph is an array of up to 59 specified periods.

SUBROUTINE CUTOFF: This determines the periods corresponding to a given phase velocity by an interval halving technique to an accuracy of 0.001 sec. The routine starts with an initial period T_1 at velocity C_1 . $T_2 = T_1 + DT$ is generated and the period equation is again evaluated at phase velocity C_1 . The signs of the period equation values are compared and the period varies until a zero crossing is found. At this point interval halving is used to determine the period of the zero crossing. The procedure starts over again and searches for the next zero crossing until KMAX zero crossings, or cutoffs, are found.

SUBROUTINE DISPER: This routine determines the phase velocities of up to ten modes at period $T(1)$ and then follows these modes until the phase velocities have been determined for all periods. A mode is followed until its phase velocity becomes greater than $B(MMAX)$, the cutoff phase velocity. The routine brackets a phase velocity and calls SUBROUTINE NEVILLE to refine the value. This routine also determines the group velocity, amplitude factors and Rayleigh wave ellipticity.



The diagram at the upper left shows the operation of the mode following technique for DT positive, C_1 an underestimate, and DC positive. There will be no mode jumping if DC is small enough and if C increases monotonically.



The diagram at lower left shows the operation of the mode following technique for DT negative, C_1 an overestimate, and DC negative. Mode jumping may be a problem if DC or DT are too large. This is not a problem for DT and DC positive, since the program has been written to watch for this.

SUBROUTINE NEVILLE: The desired root of the period equation has been bracketed by DISPER. This subroutine refines it to an accuracy of 0.0000001. The method is as follows. The root refining proceeds by interval halving until a suitable region is found where the Neville iteration method can be used. The Neville iteration method starts out as a linear interpolation, then in the next iteration becomes quadratic interpolation, and proceeds up to a 10 degree interpolation. If at any time the Neville method becomes unsuitable, e. g. when it would lead to slow or improper convergence to the correct solution, the interval halving procedure is returned to. This combination of interval halving and polynomial interpolation has been designed for maximum efficiency.

FUNCTION DLTAR: The layer dropping procedure is performed here. Layer dropping will not occur in the middle of root refinement by Neville, since layer dropping affects the magnitude of the period equation but not its sign. This function also serves to call the period equation, ellipticity, and amplitude response routines.

FUNCTION DLTAR1: This is the Love wave formulation. Haskell's matrix is used, but working up from the bottom layer to the free surface. This routine is valid for C = B(MMAX).

FUNCTION DLTAR4: This is the Rayleigh wave function. The Dunkin-Thrower algorithm is used because it is not affected by roundoff and truncation error as much as the Haskell formulation. Ellipticity is positive for retrograde elliptical motion because the coordinate system has been chosen such that radial = positive and z-axis = positive downward. The function is valid at C = B(MMAX).

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- Haskell, N. A. (1953). The dispersion of surface waves on multi-layered media, Bull. Seism. Soc. Am. 43, 17-34,

Isaacson, I., and H. B. Keller (1966). Analysis of Numerical Methods, John Wiley, New York.

Schwab, F., and L. Knopoff (1970). Surface-wave dispersion computations, Bull. Seism. Soc. Am. 60, 321-344.

Watson, T. H. (1970). A note on fast computation of Rayleigh wave dispersion in the multi-layered elastic halfspace, Bull. Seism. Soc. Am. 60, 161-166.

INPUT DATA

Card Sequence	Column	Name	Format	Explanation
A.	1-4	MMAX	I4	LT.0 end program
				EQ.0 use previously inputted model with new options
				GT.0 number of layers in model including halfspace
	5-8	MODE	I4	Number of modes for which dispersion curves are desired

B. Earth Model (MMAX cards)

1-10	D(I)	F10.4	Layer thickness in km
11-20	A(I)	F10.4	P wave velocity in km/sec
21-30	B(I)	F10.4	S wave velocity in km/sec
31-40	RHO(I)	F10.4	Density in gm/sec

C. Control Card

1-4	IGPHL I4	GT.O calls Love wave period equation plot in SUBROUTINE GPHDIS
		LE.O skip Love plot
11-14	IGPHR I4	GT.O calls Rayleigh wave period equation plot in SUBROUTINE GPHDIS
		LE.O skip Rayleigh plot
21-24	ICUT I4	GT.O calls CUTOFF to determine the modal cutoff periods of Love and Rayleigh waves.
		LE.O skip CUTOFF
31-34	IDISPL I4	GT.O calls DISPER and determines phase and group velocity and amplitude factor for Love waves.
		LE.O skip Love wave dispersion determination
41-44	IDISPR I4	GT.O calls DISPER to determine phase and group velocity, ellipticity and amplitude factor for Rayleigh waves
		LT.O skip Rayleigh wave dispersion determination

C. Control Card

51-54	IPUNCH I4	GE.1 output of DISPER also on cards
		LT.1 no card output from DISPER
61-70	FACT F10.5	If zero or negative program sets FACT = 4.0.

D. If IGPHL or IGPHR GT 0, SUBROUTINE GPHDIS is called which requires the following data cards at this point

a.	1-4	KK	I4	number of periods along abscissa. KK .LE. 59.
	11-20	C1	F10.5	lower phase velocity limit
	21-30	C2	F10.5	upper phase velocity limit of ordinate
	31-40	DC	F10.5	phase velocity increment along ordinate
b.	1-80	T(I)	16F5.0	I = 1, KK. These periods form the abscissa of the plot

E. If ICUT GT 0, SUBROUTINE CUTOFF is called which requires the following data cards at this point

1-5	KMAX	I5	number of cutoff periods to be found for phase velocity C1
11-10	T1	F10.4	initial period of search
21-30	DT	F10.4	period increment (negative if going from long period toward short period)
31-40	C1	F10.4	phase velocity for which KMAX cutoffs are desired; can equal B(MMAX)

F. If IDISPL or IDISPR GT 0, SUBROUTINE DISPER is called which uses the following data:

1-4	KMAX	I4	number of periods for which phase velocity is to be determined
-----	------	----	--

F.

11-20	T1	F10.4	initial period. If T1 is zero or negative, the array T(I) I = 1, MMAX are read in on the next set of cards. If positive T(I) are generated by $T(I) = T1 + (I - 1)*DT$. (I) = 1, MMAX.
21-30	DT	F10.4	period increment. Only has meaning if T1 is positive.
31-40	C1	F10.4	initial phase velocity guess. May not be less than 0.8 B(1). For T(I) increasing with I, C1 must be an underestimate.
41-50	DC	F10.4	increment to phase velocity for determining roots. If C1 is an underestimate, DC must be positive. Next guess is $C2 = C1 + DC$. DC must be small enough that modes are not skipped.

(This set of cards is read next if T1 is not positive.)

1-80	T(I)	8F10.0	I = 1. KMAX array of periods at which phase velocities are to be determined.
------	------	--------	--

If all routines are to be used, the data are assembled in the above order. If a particular option is to be skipped, just skip the data cards listed under that subroutine.

At this point the program returns to Point A to terminate program, run other options, or to consider a new earth model.

```

C      PROGRAM SURFACE
C      THIS PROGRAM WILL ACCEPT ONE LIQUID LAYER AT THE SURFACE.
C      PROGRAM DEVELOPED BY ROBERT B HERRMANN SAINT LOUIS UNIVERSITY
C      NOV 71 IN WHICH CASE ELLIPTICITY OF RAYLEIGH WAVE IS THAT AT THE T
C      TOP OF SOLID ARRAY. LOVE WAVE COMMUNICATIONS IGNORE LIQUID LAYER
C      COMMON D(100),A(100),B(100),RHO(100),NMAX,MMAX,IDROP,FACT,IPUNCH
C      COMMON T(100),C(100,10),U(100,10),RATIO(100,10),MODE
   4 FORMAT(1H ,40X,4F10.4)
  11 FORMAT(14,14,2X,F10.5)
   5 FORMAT(1H ,50X,3F10.4)
   44 FORMAT(4F10.4)
1212 FORMAT(14,6X,4F10.4)
   55 FORMAT(10X,3F10.4)
  801 READ 11,NMAX,MODE
C      NMAX = NUMBER OF LAYERS TO BE READ IN, INCLUDING HALFSPACE
C      MODE = NUMBER OF MODES FOR WHICH DISPERSION CURVES ARE DESIRED
  IF(NMAX) 777,779,776
778 PRINT 20
  20 FORMAT(1H1,/,1H ,54X,13HCRUSTAL MODEL ,/1H )
  PRINT 21
  21 FORMAT(1H ,42X,40H THICK P VEL S VEL DENSITY /1H0)
L=MMAX-1
D = THICKNESS OF LAYER IN KILOMETERS
A = COMPRESSIONAL WAVE VELOCITY IN KM/SEC
B = TRANSVERSE WAVE VELOCITY IN KM/SEC
RHO = DENSITY IN GM/CC
DO 3 I=1,L
  READ 44, D(I),A(I),B(I),RHO(I)
  3 PRINT 4, D(I),A(I),B(I),RHO(I)
  READ 55, A(MMAX),B(MMAX),RHO(MMAX)
  PRINT 5, A(MMAX),B(MMAX),RHO(MMAX)
779 IF(NMAX.EQ.0) NMAX = NMAX
IGPHL GT 0 CALLS LOVE WAVE DISPERSION PLOT
IGPHR GT 0 CALLS RAYLEIGH WAVE DISPERSION PLOT
ICUT GT 0 CALLS SEARCH FOR LOVE AND RAYLEIGH HIGHER MODE CUTOFF
IDISPL GT 0 LOVE WAVE DISPERSION CURVE
IDISPR GT 0 RAYL WAVE DISPERSION CURVE
IPUNCH .GE. 1 PUNCHED OUTPUT OF DISPERSION CURVES ALSO
FACT = NUMBER OF WAVELENGTHS BELOW FIRST LAYER WHERE PHASE
          VELOCITY IS LESS THAN SHEAR VELOCITY THAT WE MAY CONSIDER
          TO BE EFFECTIVE HALFSPACE. SEE KNOPOFF SCHWAB-SURFACE WAVE
          COMPUTATIONS- BSSA PH321-344 VOL 60 NO 2 APR 70
READ 12, IGPHL,IGPHR,ICUT,DISPL,DISPR,IPUNCH,FACT
12 FORMAT(6(I4,6X),F10.5)
  IF(FACT)100-100+101
100 FACT = 4.0
101 CONTINUE
NMAX = MMAX
IF(IGPHL.GT.0.OR.IGPHR.GT.0) CALL GPHDIS(IGPHL,IGPHR)
IF(ICUT.GT.0) CALL CUTOFF

```

SURFACE PAGE 2

```
IF(IDISPL.GT.0.OR.IDISPR.GT.0) CALL DISPER(IDISPL,DISPR)
GO TO 801
777 CONTINUE
END

SUBROUTINE GPHDIS(IGPHL,IGPHR)
COMMON E(100),A(100),B(100),RHQ(100),NMAX,MMAX,IDROP,FACT,IPUNCH
COMMON T(100),C(100,10),U(100,10),RATIO(100,10),MODE
DIMENSION KK(62)
DIMENSION XE(62)
C THIS SUBROUTINE GRAPHICALLY DISPLAYS THE SIGN OF THE LOVE OR
C RAYLEIGH WAVE PERIOD EQUATION IN THE C-T PLANE. THE DISPERSION
C CURVE IS THE LINE OF ZEROES.
11 FORAT(16F5.0)
C GPHDIS READS IN UP TO 59 DIFFERENT PERIODS TO FORM ABSCISSA OF
C PLOT
C THE ORDINATE VARIES FROM C1 TO C2 IN INCREMENTS OF DC
C KK IS THE NUMBER OF ABSCISSA VALUES
C C1 IS LESS THAN C2
C DC IS POSITIVE
READ 9,    KK,C1,C2,DC
9 FORMAT(I4,6X,3F10.5)
READ 10,   (T(I),I=1,KK)
IF(KK.GT.59) KK = 59
IF(CC.LT.0.0) DC = - DC
IF(C1.LT.C2) GO TO 2000
DUM = C1
C1 = C2
C2 = DUM
2000 CONTINUE
PRINT 9998
9998 FORMAT(1H1,,0X,40HPERIOD FOR ABSCISSA OF FOLLOWING GRAPH /1H0)
PRINT 9999,((I,T(I)),I=1,KK)
9999 FORMAT(1H ,12(13'F7.2))
DO 801 IFUNC = 1,2
IF(CC.EQ.0.0) GO TO 801
GO TO (811,812),IFUNC
811 IF(IGPHL.LE.0) GO TO 801
PRINT 7
7 FORMAT(1H1, 54X,25H PLOT OF LOVE FUNCTION/1H0)
GO TO 813
812 IF(IGPHR.LE.0) GO TO 801
PRINT 8
8 FORMAT(1H1, 54X,25H PLOT OF RAYLEIGH FUNCTION/1H0)
813 CC = C2
1004 DO 1010 I=1,KK
IDROP = 0
DEL = DLTAR(CC,T(I),IFUNC)
L = MMAX
L10 = L/10
```

SURFACE PAGE 3

```
L = L - L1*10
KX(1) = IABS(L)
IF(JEL)1011,1012,1013
1011 KS(1) = 1H-
GO TO 1010
1012 KS(1) = 1H
GO TO 1010
1013 KS(1) = 1H
1010 CONTINUE
PRINT 1100, CC, ((KS(L),KX(L)),L=1,KK)
1100 FORMAT(1H *E7.3,59(A1,I1))
CC = CC + 00
IF(CC = C1) 801,1004,1004
801 CONTINUE
RETURN
END
```

SUBROUTINE CUTOFF

```
COMMON U(100),A(100),B(100),RHO(100),NMAX,MMAX,IRDOP,FACT,IPUNCH
COMMON T(100),C(100,10),U(100,10),RATIO(100,10),MODE
C KMAX = NUMBER OF CUTOFF PERIODS TO BE FOUND FOR PHASE VELOCITY C1
C T1 = INITIAL PERIOD IN SEARCH
C DT = PERIOD INCREMENT NEGATIVE IF STARTING AT HIGH PERIOD AND
C       AND GOING TOWARD SHORTER PERIOD
C C1 = PHASE VELOCITY FOR WHICH KMAX CUTOFFS ARE BEING FOUND
C THIS ROUTINE FINDS BOTH LOVE AND RAYLEIGH CUTOFFS
READ 2000,KMAX,T1,DT,C1
2000 FORMAT(I5,5X,3F10.4)
PRINT 996
996 FORMAT(1H1/1H0/1H0)
PRINT 998
998 FORMAT(1H0/1H0,51X,26HHIGHER MODE CUTOFF PERIODS /1H0,34X,
1 10H LOVE ,40X,10H RAYLEIGH /1H0,2(28X,6HPERIOD,10X,6HPH VEL
2 )//)
LL = 0
TT = T1
DO 47 IFUNC = 1,2
K = 1
IRDOP = 0
T1 = TT
999 DEL1 = DLTAR(C1,T1,IFUNC)
80 T2 = T1 + DT
IRDOP = 0
IF(T2)3000,3000,3001
3000 J = K + LL
T(J) = 0.0
KMAX = K
3001 CONTINUE
11 DEL2=DLTAR(C1,T2,IFUNC)
IF(ABS(DEL2-DEL1)-ABS(DEL2+DEL1))81,81,54
```

```

51 T1 = T2
    DEL1 = DEL2
    GO TO 80
54 IF(ABS(T1 - T2) = 0.0001) 120,120,130
130 T3 = (T1 + T2) * 0.5
    DEL3 = DELTAR(C1,T3,IFLNC)
    IF(ABS(DEL3 - DEL1) = ABS(DEL3 + DEL1))101,120,102
101 T1 = T3
    DEL1 = DEL3
    GO TO 54
102 T2 = T3
    DEL2 = DEL3
    GO TO 54
120 TF = (T1 + T2) * 0.5
    J = K + LL
    T(J) = TF
    T1 = TF + DT
    IF(K-KMAX)76,46,46
75 K = K + 1
    IDROP = 0
    GO TO 999
46 LL = 50
47 CONTINUE
997 PRINT 997, ((T(J),C1,T(J+50),C1),J=1,KMAX)
FORMAT(1H,24X,F10.4,6X,F10.4,24X,F10.4,6X,F10.4)
RETURN
END

```

```

SUBROUTINE DISPER(IDISPL,DISPR)
COMMON D(100),A(100),B(100),RHO(100),NMAX,MMAX,IDROP,FACT,IPUNCH
COMMON T(100),C(100,10),U(100,10),RATIO(100,10),MODE
DIMENSION IMAX(10)
DIMENSION AR(100,10)

```

C THE ROOT DETERMINATION SECTION IS ONE OF INTERVAL HALVING ONCE A
C ZERO CROSSING HAS BEEN FOUND.

C TO FOLLOW MODES THE PROGRAM INITIALLY FINDS THE PHASE VELOCITIES
C OF THE MODE NUMBER OF MODES FOR PERIOD T1. THEN IT FINDS THE
C PHASE VELOCITIES FOR PERIOD T1 + DT USING PREVIOUS RESULTS TO STAY
C ON THE SAME MODE. IF STARTING AT SHORT PERIODS DC MUST BE SMALL
C ENOUGH SO MODE IS NOT JUMPED. IT IS PREFERABLE TO HAVE DC POS
C SINCE THE PROGRAM WILL FOLLOW A MODE UP TO ITS CUTOFF VALUE. FOR
C DC NEG THE PROGRAM CANNOT PICK UP ANY MODES.

```
1000 FORMAT(14.6X,4F10.4)
```

```
1001 FORMAT(8F10.0)
```

C T1 = INITIAL STARTING PERIOD

C KMAX = NUMBER OF PERIODS FOR WHICH PHASE VELOCITY IS TO BE
C DETERMINED

C IF T1 = 0 PROGRAM READS IN ARRAY OF T(I) PERIODS INSTEAD OF
C COMPUTING THEM

C DT = PERIOD INCREMENT. NEXT PERIOD T2 = T1 + DT

```

C      C1 = INITIAL PHASE VELOCITY GUESS. MAKE SURE IT IS OUTSIDE
C      DESIRED RESULT
C      DC = PHASE VELOCITY INCREMENT = POS IF GOING FROM LOWER TO HIGHER
C          NEXT PHASE VELOCITY GUESS IS C2 = C1 + DC
READ 1000,KMAX,T1,BT,C1,DC
IF(T1)300,300,301
300 READ 1001,(T(J),J=1,KMAX)
GO TO 302
301 T(1) = T1
DO 303 I = ?,KMAX
303 T(I) = T(I-1) + DT
302 CONTINUE
CC = C1
DO 9999 IF(IFNC = 1,2
GO TO 10,10,IFUNC
10 IF(I0ISPL,LB,0) GO TO 9999
GO TO 30
20 IF(I0ISPR,LB,0) GO TO 9999
30 CONTINUE
C1 = CC
PRINT 1
1 FORMAT(1H1,
KNOWF = MODE
DO 2 I=1,1C
2 IA^(I)=0
DO 9998 K = 1,KMAX
T1 = T(K)
DO 9997 IQ = 1,KMODE
IF(K-1)605,605,599
599 IF(IQ - 2)600,601,601
600 C1 = C(K-1,1) - 0.01 * ABS (DC)/DC
GO TO 605
601 IF(DC*(C(K-1,IQ) - C(K,IQ-1)))602,603,604
602 C1 = C(K,IQ-1) + 0.01*ABS (DC)/DC
GO TO 605
603 C1 = C(K,IQ-1) + 0.01*ABS (DC)/DC
GO TO 605
604 C1 = C(K-1,IQ)
605 CONTINUE
IDRUP = 0
999 DEL1 = DLTAR(C1,T1,IFUNC)
80 C2=C1+DC
IDRUP = 0
DEL2 = DLTAR(C2,T1,IFUNC)
IF (SIGN(1.,DEL1).NE.SIGN(1.,DEL2)) GO TO 54
81 C1=C2
DEL1=DEL2
C      CHECK THAT C1 IS IN REGION OF SOLUTIONS
IF(C1-0.8*B(1))250,251,251
251 IF(C1-(B(MMAX)+0.3))252,250,250

```

SURFACE PAGE 6

```
252 GO TO 80
54 CALL NEVILLE(T1,C1,C2,DEL1,DEL2,IFUNC,CN)
C1 = CN
IF(C1 - B(MMAX))121,121,250
121 C(K,IQ) = C1
P = DLTAR(C1,T1+0.001,IFUNC)
Q = DLTAR(C1,T1-0.001,IFUNC)
R = BLTAR(C1+0.001,T1,IFUNC)
S = BLTAR(C1-0.001,T1,IFUNC)
DCET = -(P-Q)/(R-S)
U(K,IQ) = C1/(1. + T1*DCET/C1)
GO TO (6001,6002),IFUNC
6002 CONTINUE
C RATIO = RAYLEIGH WAVE ELLIPTICITY
RATIO(K,IQ) = DLTAR(C1,T1,3)
C AR = RAYLEIGH WAVE AMPLITUDE RESPONSE. SEE HARKRIDER ESSA DEC 70.
C AR = (2.*U*C*10)**(-1)
AR(K,IQ) = ABS(DLTAR(C1,T1,4)*7.01256637/(C1*C1*T1*(R-S)))
GO TO 6003
C HERE AR = LLVE WAVE AMPLITUDE RESPONSE
6001 CONTINUE
AR(K,IQ) = ABS(DLTAR(C1,T1,5)*0.01256637/(C1*C1*T1*(R-S)))
6003 CONTINUE
C1 = C1 + 0.01 * ABS(DC)/DC
IMAX(IQ) = K
GO TO 9997
250 IF(K*IQ-1)256,256,255
256 PRINT 258
258 FORMAT(1H ,40H IMPROPER INITIAL VALUE NO ZERO FOUND )
GO TO 9999
9997 CONTINUE
GO TO 9998
255 KMODE = IQ - 1
IF(KMODE) 9996,9996,9998
9998 CONTINUE
9996 DO 9995 IQ = 1,MODE
J=IMAX(IQ)
IF(J)9995,9995,9994
9994 GO TO (7001,7002),IFUNC
7001 PRINT 103
105 FORMAT(1Hu,55X,10HLCVE WAVE /1H ,45X,30H PERIOD PHASE VEL GROUP
LVEL ,7X,2HAL /1H )
PRINT 106, ((T(I),C(I,IQ),U(I,IQ),AR(I,IQ)),I=1,J)
106 FORMAT(1H ;45X,0PF7.3,1X,0PF10.6,2X,0PF8.4,4X,1PE12.5)
IF(IPUNCH,LT.1) GO TO 9995
PUNCH 1106, ((T(I),C(I,IQ),U(I,IQ),AR(I,IQ)),I=1,J)
1106 FORMAT(4(1PE14.7,1X),15X,3HLCV,12)
GO TO 9995
7002 PRINT 103
103 FORMAT(1H0,53X,14HRAYLEIGH WAVE /1H ,30X,
```

```

131H PERIOD   PHASE VEL GROUP VEL ,6X,5HUR/UZ, 12X,2HAR/1H )
PRINT 104, ((T(I),C(I,IQ),U(I,IQ),RATIO(I,IQ),AR(I,IQ)),I=1,J)
104 FORMAT(1H ;30X,0PF7.3,1X,0PF10.6,2X,0PF8.4,4X,1PE12.5,4X,1PE12.5)
IF(IPUNCH,LT,1) GO TO 9995
PUNCH 1104, ((T(I),C(I,IQ),U(I,IQ),AR(I,IQ),RATIO(I,IQ),IQ),
  I=1,J)
1104 FORMAT(5(1PE14.7,1X),3HRAY,I2)
9995 CONTINUE
9999 CONTINUE
      RETURN
      END

SUBROUTINE NEVILLE(T,C1,C2,DEL1,DEL2,IFUNC,CC)
C HYBRID METHOD FOR REFINING ROOT ONCE IT HAS BEEN BRACKETTED
C BETWEEN C1 AND C2. INTERVAL HALVING IS USED WHERE OTHER SCHEMES
C WOULD BE INEFFICIENT. ONCE SUITABLE REGION IS FOUND NEVILLE'S
C ITERATION METHOD IS USED TO FIND ROOT.
C THE PROCEDURE ALTERNATES BETWEEN THE INTERVAL HALVING AND NEVILLE
C TECHNIQUES USING WHICHEVER IS MOST EFFICIENT
COMMON B(100),A(100),B(100),RHO(100),NMAX,MMAX,IRDROP,FACT,IPUNCH
COMMON C1(100),E2(100,10),D3(100,10),D4(100,10),MODE
DIMENSION X(20),Y(20)
C3 = 0.5*(C1 + C2)
DEL3 = DELTAU(C3,T,IFUNC)
NEV = 1
1310 IF(C1 - C3)1320,1320,1330
1320 IF(C2 - C3) 1344,1344,1000
1330 IF(C2 - C3) 1000,1344,1344
1000 S13 = DEL1 - DEL3
S32 = DEL3 - DEL2
IF(SIGN(1.,DEL3)*SIGN(1.,DEL1)) 1441,1441,1443
1441 C2 = C3
DEL2 = DEL3
GO TO 1444
1443 C1 = C3
DEL1 = DEL3
GO TO 1444
1444 CONTINUE
IF(ABS(C1 - C2) -0.0000001) 20,20,22
22 CONTINUE
IF(SIGN(1.,S13).NE.SIGN(1.,S32)) NEV = 0
SS1=ABS(DEL1)
S1=0.1*SS1
SS2=ABS(DEL2)
S2=0.1*SS2
IF(S1.GT,SS2.OR,S2.GT,SS1) GO TO 1344
IF(NEV.EQ.0) GO TO 1344
IF(NEV.EQ.2) GO TO 1350
X(1) = C1
Y(1) = DEL1

```

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```
X(2) = C2
Y(2) = DEL2
M = 1
GO TO 1355
1344 C3 = 0.5*(C1 + C2)
DEL3 = DLTAR(C3,T,IFUNC)
N2V = 1
M = 1
GO TO 1310
1350 X(M+1) = C3
Y(M+1) = DEL3
GO TO 1355
1355 DO 1360 KK=1,M
J = M-KK+1
IF (ABS(Y(M+1)-Y(J)).LE.1.E-10) GO TO 1344
X(J) = (-Y(J)*X(J+1) + Y(M+1)*X(J))/(Y(M+1)-Y(J))
1360 CONTINUE
21 CONTINUE
C3 = X(1)
DEL3 = DLTAR(C3,T,IFUNC)
N2V = 2
M = M+1
IF(M.GT.10) M = 10
GO TO 1310
20 CONTINUE
CC = C3
RETURN
END

FUNCTION DLTAR(CC,TT,KK)
COMMON D(100),A(100),B(100),RHO(100),NMAX,MMAX,IDRCP,FACT,IPUNCH
COMMON T(100),C(100,10),L(100,10),PAT(100,10),MODE
C LAYER DROPPING PROCEDURE--HALFSPACE DEFINED AS FACT WAVELENGTHS
C BELOW REGION WHERE C IS LESS THAN R
IF(IDRCP)899,899,905
899 CONTINUE
DMAX = FACT * CC * TT
MMAX = NMAX
SUM = 0
DO 900 II = 1,NMAX
IF(CC-B(II))901,900,900
901 SUM = SUM + D(II)
IF(SUM - DMAX) 900,900,902
902 MMAX = II
GO TO 904
900 CONTINUE
904 IDRCP = 1
IF(MMAX.LT.2) MMAX = 2
905 CONTINUE
GO TO (1,2,3,4,5),KK
```

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```
C      LOVE WAVE PERIOD EQUATION
1  DLTAIR = DLTAIR1(CC,TT,1)
  RETURN
C      RAYLEIGH WAVE PERIOD EQUATION
2  DLTAIR = DLTAIR4(CC,TT,1)
  RETURN
C      RAYLEIGH WAVE ELLIPTICITY
3  DLTAIR = DLTAIR4(CC,TT,2)
  RETURN
C      RAYLEIGH WAVE AMPLITUDE RESPONSE COMPONENT
4  DLTAIR = DLTAIR4(CC,TT,3)
  RETURN
C      LOVE WAVE AMPLITUDE RESPONSE COMPONENT
5  DLTAIR = DLTAIR1(CC,TT,2)
  RETURN
END
```

```
FUNCTION DLTAIR1(C,T,MUP)
HASKELL-THOMPSON LOVE WAVE FORMULATION FROM HALFSPACE TO SURFACE
COMMON D(100),A(100),B(100),RHO(100),NMAX,MMAX,IDROP,FACT,IPUNCH
COMMON D1(100),D2(100,10),D3(100,10),D4(100,10),MODE
WVNO = 6.2931853/(C*T)
CCVB = C/B(MMAX)
H = RHO(MMAX)*B(MMAX)*B(MMAX)
RB = SQRT (ABS (CCVB**2-1.))
UT = 1
TT = H * RB
MM1 = MMAX - 1
DO 1340 K = 1,MM1
M = MMAX - K
IF (K.EQ.0) GO TO 1340
CCVB = C/B(K)
RB = SQRT (ABS (CCVB**2-1.))
H = RHO(M)*B(M)*B(M)
7001 S = -WVNO*D(H)*RB
IF (C-B(M))1209,1221,1231
1231 SINC = SIN(S)
Y = SINC/RB
Z = RB*SINC
COSQ = COS(P)
GO TO 1242
1221 Y = -WVNO*D(M)
Z = 0
COSQ = 1
GO TO 1242
1209 EXQP = EXP(C)
EXGM = 1./EXQP
Y = (EXQP - EXGM)/(2.*RB)
Z = -RB*RB*Y
COSQ = (EXQP + EXGM)/2.
```

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```
1242 EUT = COSQ#LT - Y#TT/H
      EIT = H#Z#UT + COSQ#TT
      UT = EUT
      TT = ETT
1340 CONTINUE
      GO TO (1,2),MUP
1     DLTARI1 =-ETT
      RETURN
2     DLTARI1 = UT
      RETURN
      END

FUNCTION DLTARI4(C,T,MUP)
C MUP = 1   DLTARI4 = RAYLEIGH WAVE PERIOD EQUATION
C MUP = 2   DLTARI4 = RAYLEIGH WAVE ELLIPTICITY
C MUP = 3   DLTARI4 = RAYLEIGH WAVE AMPLITUDE FACTOR
C UP TO ONE LIQUID LAYER IS ALLOWED AT THE SURFACE
C MODIFIED DUNKIN-THROWER ALGORITHM FOR RAYLEIGH WAVES
C C.F. J. DUNKIN BSSA VOL 55 NO 2 PP 335-358 APR 65
C C.F. C. WATSON BSSA VOL 60 NO 1 PP 161-166 FEB 70
C C.F. D. HARKRIDER BSSA VOL 60 NO 6 PP 1937-1987 DEC 70
C HIGHER MODE CUTOFF MODIFICATION BY R. HERRMANN
C COMMON D(100),P(100),S(100), RHO(100),NMAX,MMAX, IDROP,FACT,IPUNCH
C COMMON D1(100),D2(100,10),D3(100,10),D4(100,10),MCDE
99 NVNG = 6.26318531/(C*T)
CSQ = C*C
JUMP=1
1000 IF(JUMP.GT.1) JUMP=2
      B1 = 0.0
      B2 = 0.0
      B3 = 0.0
      B4 = 0.0
      B5 = 0.0
      IF(JUMP = 2)1,2,3
1     B1 = 1.0
      GO TO 4
2     B2 = 1.0
      GO TO 4
3     B3 = 1.0
      GO TO 4
4     CONTINUE
      DO 50 M=1,NMAX
100 ARGA = 1. - CSQ/P(M)**2
      RA = SQRT(ABS(ARGA))
      IF(ARGA)202,202,201
201 RA = -RA
202 IF(S(M))101,103,101
C LIQUID SURFACE LAYER
103 P=MVNO*RA*D(M)
      IF(MUP.GT.1) GO TO 50
```

```

RHOC=RHO(M)*CSQ
IF(RA)313,312,314
312 SINPR = WVNO*D(M)
RSINP = 0.0
COSP = 1.0
GO TO 315
313 SINPR = (EXP(PM)-EXP(-PM))/(2.*RA)
RSINP = -RA * RA * SINPR
COSP = 0.5*(EXP(PM)+EXP(-PM))
GO TO 315
314 SINPR = SIN(PM)/RA
RSINP = RA * SIN(PM)
COSP = COS(PM)
315 CONTINUE
A11 = COSP
A21 = RHOC*SINPR
A31 = 0.0
A41 = 0.0
A51 = 0.0
A12 = 0.0
A22 = 0.0
A32 = 0.0
A42 = 0.0
A13 = 0.0
A23 = 0.0
A33 = 0.0
A14 = 0.0
A24 = 0.0
A15 = 0.0
GO TO 1001
101 ARGB=1.-CSQ/S(M)**2
RB = SQRT(AWS(ARGB))
IF(ARGB)204,204,203
203 RB = -RB
204 CONTINUE
G=2.*S(M)**2/CSQ
G1=G-1.0
IF(XMAX-M)40,52,40
40 RHOC=RHO(M)*CSQ
PM = WVNO*RA*D(M)
QR = WVNO * RB * D(M)
IF(RA)213,212,214
212 RSINP = 0.0
SINPR = WVNO*D(M)
COSP = 1.0
GO TO 215
213 RSINP = -RA*0.5*(EXP(PM)-EXP(-PM))
SINPR=-RSINP/(RA**2)
COSP = 0.5*(EXP(PM)+EXP(-PM))
GO TO 215

```

```

214 RSINP = RA*SIN(PM)
    SINPR=RSINP/(RA**2)
    COSP = COS(PM)
215 IF(RB) 216,218,217
216 RSING = -RB*.5*(EXP(GM)-EXP(-GM))
    SINQR=-RSING/(RB**2)
    COSQ = .5*(EXP(GM)+EXP(-GM))
    GO TO 219
217 RSIQ = RB*SIN(QM)
    SINQR=RSIQ/(RB**2)
    COSQ = COS(QM)
    GO TO 219
218 RSIQ=0.0
    SINQR = WVND*C(%)
    COSQ = 1.0
    GO TO 219
219 RR=RSINP*RSING
    SS=SINPR*SINQR
    CC=COSP*COSQ
    A11 = -2.*G*G1+(2.*G**2-2.*G+1.)*CC-G**2*RR-G1**2*SS
    A12 = -(1./RHOC )*(RSINP*COSQ+SINQR*COSP)
    A13 = -(2./RHOC )*((2.*G-1.)*(1.-CC)+G1*SS+G*RR)
    A14 = ( 1.0/RHOC )*(SINPR*COSQ+RSING*COSP)
    A15 = (RHOC*(-2.))*(2.*(1.-CC)+RR+SS)
    A21 = RHOC*(G1**2*SINPR*COSQ+G**2*RSING*COSP)
    A22 = CC
    A23 = 2.*G*(G*RSING*COSP+G1*SINPR*COSQ)
    A24 = SINPR*RSING
    A31 = RHOC*(G*G1*(2.*G-1.)*(1.-CC)+G1**3*SS+G**3*RR)
    A32 = G1*SINQR*COSP+G*RSINP*COSQ
    A33 = 1.0+2.4*(2.*G*G1*(1.-CC)+G**2*RR+G1**2*SS)
    A41 = -RHOC*(G**2*RSINP*COSQ+G1**2*SINQR*COSP)
    A42 = RSINP*SINQR
    A51 = RHOC**2*((2.*G**2*G1**2)*(1.-CC)+G1**4*SS+G**4*RR)
1001 CONTINUE
C MATRIX MULTIPLICATION
C EFFECT OF IMAGINARY ELEMENTS INCLUDED
B61 = A11*B1 + A12*B2 + A13*B3 + A14*B4 + A15*B5
B62 = A21*B1 + A22*B2 + A23*B3 + A24*B4 - A14*B5
B63 = A31*B1 + A32*B2 + A33*B3 -.5*A23*B4 + .5*A13*B5
B64 = A41*B1 + A42*B2 -2.*A32*B3 + A22*B4 - A12*B5
B65 = A51*B1 - A41*B2 +2.*A31*B3 - A21*B4 + A11*B5
B1 = BB1
B2 = BB2
B3 = BB3
B4 = BB4
B5 = BB5
50 CONTINUE
C THE FOLLOWING EXPRESSION IS VALID FOR RB = 0
52 A11 = -2.*RB*(S(M)/P(M))**2*(CSQ*G1**2)/(P(M)**2*G*RA)

```

```

A12 = -( RHO(M) *P(M)**2*G )**(-1)
A13 = -RB/(RHO(M)*P(M)*P(M))+G1/(RHO(M)*P(M)*P(M)*G*RA)
A14 = RR/(G*RHO(M)*P(M)*P(M)*RA)
A15 = ((RH0(M)*P(M))**2*CSQ*G)**(-1)*(RB-1./RA)
B61 = A11*B1 + A12*B2 + 2.*A13*B3 + A14*B4 + A15*B5
GO TO 501,502,503),MLP
501 DLTAR4=-B61
RETURN
502 IF(JUMP.EQ.2) R12 = B61
JUMP = JUMP + 1
IF(JUMP.EQ.3) GO TO 1000
DLTAR4 = 0.5*B61/R12
RETURN
503 DLTAR4 = ABS(B61)
IF(S(1).GT.0.0) RETURN
RA = C/P(1)
RAD = WENO * D(1) * SQRT(ABS(RA*RA - 1.))
DLTAR4 = ABS(B61 * COS(RAD))
RETURN
END

```

DATA PAGE 1

5 2
1.0 5.0 2.89 2.5
9.0 6.1 3.52 2.7
10. 6.4 3.7 2.9
20. 6.7 3.87 3.0
8.15 4.7 3.4
1 1 1 1 0 3.0
24 2.6 4.7 0.1
2.0 3.0 4.0 5.0 6.0 7.0 8.0 9.0 10. 11. 12. 13. 14. 15. 16. 17
18. 19. 20. 22. 24. 26. 28. 30.
4 18. -0.8 4.7
19 2.0 1.6 2.6 0.07
-1

CRUSTAL MODEL

THICK	P VEL	S VEL	DENSITY
1,0000	5,0000	2,8900	2,5000
9,0000	6,1000	3,5200	2,7000
10,0000	6,4000	3,7000	2,9000
20,0000	6,7000	3,8700	3,0000
	8,1500	4,7000	3,4000

PERIOD FOR ABSISSA OF FOLLOWING GRAPH

1	2.00	2	3.00	3	4.00	4	5.00	5	6.00	6	7.00	7	8.00	8	9.00
13	14.00	14	15.00	15	16.00	16	17.00	17	18.00	18	19.00	19	20.00	20	22.00

9	10.00	10	11.00	11	12.00	12	13.00
21	24.00	22	26.00	23	28.00	24	30.00

PLOT OF LOVE FUNCTION

4,700 5 5-5 5 5-5-5-5-5-5-5
4,600 5 5 5 5-5-5-5-5
4,500-5-5 5 5-5-5-5-5
4,400-5-5 5 5-5-5-5
4,300-5-5 5-5-5-5-5
4,200 5 5 5-5-5-5
4,100 5 5-5-5-5 5-5-5
4,000-5 5-5-5 5-5+5-5+5
3,900 5-5-5 5+5-5+5
3,800-5 5+5-5+5
3,700 4 5 5 5 5 5 5-5-5-5-5+5-5-5+5-5-5+5-5-5+5-5+5-5+5-5+5-5+5-5+5-5+5
3,600 4 5 5 5 5-5-5-5-5-5-5+5-5-5-5-5-5-5-5-5-5+5-5+5-5+5-5+5-5+5-5+5-5+5
3,500 4 4-5
3,400-4-4-5
3,300-4-4-5
3,200-4-4-4-5
3,100-3-4-4-5
3,000-3-4-4-5
2,900-3-4-4-5
2,800-3-4-4-5
2,700-3-4-4-5

PLOT OF RAYLEIGH FUNCTION

4,700 5 5-5 5 5 5-5-5-5-5-5-5-5+5+5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5
4,600 5 5-5 5 5-5-5-5-5-5
4,500-5-5 5 5-5-5-5-5-5
4,400-5-5 5-5-5-5-5
4,300 5-5 5-5-5-5
4,200 5 5 5-5-5-5
4,100-5 5-5-5-5
4,000-5 5-5-5
3,900 5-5-5
3,800-5-5
3,700 4 5
3,600 4 5
3,500 4 4 5
3,400 4 4 5 5 5 5 5 5 5-5-5-5+5-5-5-5-5-5-5+5-5-5-5-5+5-5-5
3,300 4 4 5 5 5 5-5-5-5-5-5-5-5+5-5-5-5-5-5-5-5+5-5-5-5
3,200 4 4 4-5
3,100-3-4-4-5
3,000-3-4-4-5
2,900-3-4-4-5
2,800-3-4-4-5
2,700-3-4-4-5

HIGHER MODE CUTOFF PERIODS

LOVE

RAYLEIGH

PERIOD	PH VEL	PERIOD	PH VEL
12,9806	4,7000	16,4834	4,7000
6.5576	4,7000	7,4149	4,7000
4.3681	4,7000	4,7535	4,7000
3,2668	4,7000	3,4268	4,7000

LOVE WAVE			
PERIOD	PHASE VEL	GROUP VEL	AL
2.000	3.407477	3.2085	7.38827E-03
3.000	3.480797	3.3257	3.89512E-03
4.000	3.526896	3.3727	2.68719E-03
5.000	3.563732	3.4016	2.07591E-03
6.000	3.595662	3.4234	1.69963E-03
7.000	3.624460	3.4409	1.44338E-03
8.000	3.651184	3.4547	1.25884E-03
9.000	3.676586	3.4650	1.12094E-03
10.000	3.701213	3.4725	1.01479E-03
11.000	3.725450	3.4776	9.30848E-04
12.000	3.749553	3.4811	8.62696E-04
13.000	3.773685	3.4835	8.05973E-04
14.000	3.797935	3.4854	7.57680E-04
15.000	3.822342	3.4872	7.15653E-04
16.000	3.846911	3.4893	6.78389E-04
17.000	3.871615	3.4920	6.44754E-04
18.000	3.896413	3.4955	6.13988E-04
19.000	3.921246	3.5001	5.85475E-04
20.000	3.946051	3.5060	5.58787E-04

LOVE WAVE			
PERIOD	PHASE VEL	GROUP VEL	AL
2.000	3.706053	3.4703	9.98171E-04
3.000	3.810811	3.5588	6.44528E-04
4.000	3.893335	3.6029	4.79195E-04
5.000	3.973622	3.5934	4.32769E-04
6.000	4.063835	3.5554	4.32677E-04
7.000	4.167551	3.5162	4.37390E-04
8.000	4.283090	3.5002	4.27207E-04
9.000	4.404134	3.5341	3.91930E-04
10.000	4.519581	3.6480	3.27520E-04
11.000	4.615236	3.8675	2.36907E-04
12.000	4.678198	4.2091	1.27248E-04

PERIOD	PHASE VEL	BAYLEIGH	WAVE	AR
		GROUP VEL	UR/UZ	
2,000	3.114265	3.0052	7.75331E-01	6.65303E-03
3,000	3.151168	3.0684	8.00347E-01	4.06700E-03
4,000	3.179362	3.0665	7.96662E-01	3.04409E-03
5,000	3.210016	3.0588	7.86004E-01	2.44811E-03
6,000	3.242074	3.0607	7.75376E-01	2.02923E-03
7,000	3.273481	3.0722	7.66864E-01	1.71251E-03
8,000	3.303160	3.0884	7.60712E-01	1.46746E-03
9,000	3.330947	3.1049	7.56531E-01	1.27642E-03
10,000	3.357235	3.1182	7.53791E-01	1.12696E-03
11,000	3.382662	3.1262	7.52007E-01	1.00926E-03
12,000	3.407898	3.1287	7.50802E-01	9.15366E-04
13,000	3.433525	3.1260	7.49919E-01	8.39774E-04
14,000	3.459986	3.1195	7.49199E-01	7.77146E-04
15,000	3.487566	3.1104	7.48571E-01	7.24097E-04
16,000	3.516395	3.1007	7.48024E-01	6.77879E-04
17,000	3.546462	3.0922	7.47594E-01	6.36441E-04
18,000	3.577626	3.0865	7.47349E-01	5.98293E-04
19,000	3.609635	3.0848	7.47373E-01	5.62368E-04
20,000	3.642155	3.0884	7.47756E-01	5.27983E-04
PERIOD	PHASE VEL	RAYLEIGH	WAVE	AR
		GROUP VEL	UR/UZ	
2,000	3.702188	3.5109	6.22385E-01	2.04879E-04
3,000	3.791024	3.5610	6.16642E-01	1.52631E-04
4,000	3.869216	3.5827	5.86135E-01	1.40361E-04
5,000	3.951656	3.5479	5.54225E-01	1.50242E-04
6,000	4.052316	3.4753	5.19288E-01	1.73177E-04
7,000	4.175212	3.4192	4.77813E-01	1.91032E-04
8,000	4.309079	3.4664	4.30754E-01	1.81464E-04
9,000	4.427527	3.6621	3.86174E-01	1.41061E-04
10,000	4.513492	3.9042	3.51858E-01	9.87478E-05
11,000	4.572207	4.0913	3.28316E-01	7.09548E-05
12,000	4.614150	4.2182	3.13487E-01	5.43378E-05
13,000	4.645877	4.3102	3.06065E-01	4.33073E-05
14,000	4.670371	4.3906	3.05473E-01	3.40924E-05
15,000	4.688310	4.4814	3.11537E-01	2.38601E-05
16,000	4.698593	4.6170	3.24455E-01	8.87479E-06

II. REIGEN

PROGRAMMER: R. B. HERRMANN / DEC 1971

PURPOSE:

This program accepts the card output for Rayleigh wave dispersion curves determined by the program SURFACE. The output of SURFACE consists of the period, phase velocity, group velocity, amplitude factor, and Rayleigh wave ellipticity. The program REIGEN will determine the Rayleigh wave eigenfunctions as a function of depth as well as phase velocity partial derivatives as a function of layer compressional and shear velocities, as well as density. The program will accept one liquid layer at the surface.

INPUT/OUTPUT

Input is from card. The eigenfunctions and phase velocities are listed on the printer. The eigenfunctions and phase velocity partial derivatives may also be listed on FILE 1. This file can then be used by other programs to generate synthetic seismograms and amplitude spectra.

THEORY

The coordinate system used is taken such that the z-coordinate is positive downwards, and the radial coordinate is positive away from the center of a cylindrical coordinate system. In this coordinate system, the ellipticity of Rayleigh waves at the surface of an elastic halfspace will be a positive quantity.

The system of differential equations to be solved is

$$\begin{vmatrix} \frac{d}{dz} UZ \\ \frac{d}{dz} TZ \\ \frac{d}{dz} UR \\ \frac{d}{dz} TR \end{vmatrix} = \begin{vmatrix} 0 & \frac{1}{\lambda+2\mu} & \frac{k\lambda}{\lambda+2\mu} & 0 \\ -\rho\omega^2 & 0 & 0 & k \\ -k & 0 & 0 & \frac{1}{\mu} \\ 0 & \frac{-k\lambda}{\lambda+2\mu} & -\rho\omega^2 + \frac{4k^2\mu(\lambda+\mu)}{\lambda+2\mu} & 0 \end{vmatrix} \begin{vmatrix} UZ \\ TZ \\ UR \\ TR \end{vmatrix}$$

where $k = 2\pi/CT$ is the wavenumber, λ and μ are the Lame constants, $\omega = 2\pi/T$ is the natural frequency, and ρ is the density. Now T , C , and the ellipticity e_0 are available from SURFACE.

The boundary conditions to the above system are that $UZ = 1$, $UR = e_0$, $TZ = TR = 0$ at $z = 0$ and that UZ and UR tend to zero as z goes to positive infinity.

The above system of equations can be derived using variational techniques and the following relation for the system Lagrangian L :

$$L = \omega^2 I_0 - k^2 I_1 - 2kI_2 - I_3 ,$$

where

$$I_0 = \int_0^\infty \rho [UZ^2 + UR^2] dz$$

$$I_1 = \int_0^\infty [(\lambda+2\mu)UR^2 + \mu UZ^2] dz$$

$$I_2 = \int_0^\infty [\mu UZ \frac{d UR}{dz} - \lambda UR \frac{d UZ}{dz}] dz$$

$$I_3 = \int_0^\infty [(\lambda+2\mu) \left(\frac{d UZ}{dz} \right)^2 + \mu \left(\frac{d UR}{dz} \right)^2] dz$$

The partials of the phase velocity with respect to layer shear velocity, compressional velocity, and density at constant frequency for the m 'th layer are defined as

$$\begin{aligned} \left(\frac{\partial c}{\partial \beta_m} \right)_{\omega, \rho, \alpha} &= 2 \rho_m \beta_m \left(\left(\frac{\partial c}{\partial \mu_m} \right)_{\rho, \lambda, \omega} - 2 \left(\frac{\partial c}{\partial \lambda_m} \right)_{\rho, \mu, \omega} \right) \\ \left(\frac{\partial c}{\partial \alpha_m} \right)_{\omega, \rho, \beta} &= 2 \rho_m \alpha_m \left(\frac{\partial c}{\partial \lambda_m} \right)_{\omega, \rho, \mu} \\ \left(\frac{\partial c}{\partial \rho_m} \right)_{\omega, \alpha, \beta} &= \left(\frac{\partial c}{\partial \rho_m} \right)_{\omega, \lambda, \mu} + (\alpha_m^2 - 2\beta_m^2) \left(\frac{\partial c}{\partial \lambda_m} \right)_{\omega, \rho, \mu} \\ &\quad + \beta_m^2 \left(\frac{\partial c}{\partial \mu_m} \right)_{\omega, \rho, \lambda} \end{aligned}$$

where

$$\left[\frac{\partial c}{\partial \mu_m} \right]_{\omega, \rho, \lambda} = \frac{c}{k} \frac{\left[\frac{\partial L}{\partial \mu_m} \right]_{\omega, \rho, \lambda, k}}{\left[\frac{\partial L}{\partial k} \right]_{\omega, \rho, \lambda, \mu}}$$

$$\left[\frac{\partial c}{\partial \lambda_m} \right]_{\omega, \rho, \mu} = \frac{c}{k} \frac{\left[\frac{\partial L}{\partial \lambda_m} \right]_{\omega, \rho, \mu, k}}{\left[\frac{\partial L}{\partial k} \right]_{\omega, \rho, \mu, \lambda}}$$

$$\left[\frac{\partial c}{\partial \rho_m} \right]_{\omega, \lambda, \mu} = - \frac{c}{k} \frac{\left[\frac{\partial L}{\partial \rho_m} \right]_{\omega, \mu, \lambda, k}}{\left[\frac{\partial L}{\partial k} \right]_{\omega, \mu, \lambda, \rho}}$$

with

$$\left[\frac{\partial L}{\partial k} \right]_{\omega, \rho, \lambda, \mu} = - 2kI_1 - 2I_2$$

$$\left[\frac{\partial L}{\partial \omega} \right]_{k, \rho, \mu, \lambda} = 2\omega I_0$$

$$\left[\frac{\partial L}{\partial \rho_m} \right]_{\omega, k, \mu, \lambda} = \omega^2 \int_{z_m}^{z_{m+1}} (UZ^2 + UR^2) dz$$

$$\begin{aligned} \left[\frac{\partial L}{\partial \mu_m} \right]_{\omega, k, \rho, \lambda} &= - \int_{z_m}^{z_{m+1}} (k^2(2UR^2 + UZ^2) + 2k UZ \frac{dUR}{dz} \\ &\quad + (2 \frac{dUZ}{dz} \frac{dUR}{dz} + \frac{dUR}{dz} \frac{dUR}{dz})) dz \end{aligned}$$

$$\left[\frac{\partial L}{\partial \lambda} \right]_{\omega, p, k, \mu} = - \int_{z_m}^{z_{m+1}} \left(k^2 UR^2 - 2k \frac{dUZ}{dz} UR + \frac{dUZ}{dz} \frac{dUZ}{dz} \right) dz$$

z_m and z_{m+1} are the depths to the top and bottom of layer m . The group velocity $U = U(ENER)$ can be defined in terms of the energy integrals as

$$U = \frac{d\omega}{dk} = \frac{(\partial L / \partial k)_{\omega}}{(\partial L / \partial \omega)_k} = \frac{kI_1 + I_2}{\omega I_0}$$

The amplitude factor of the Rayleigh wave $A_R = ARE$ can be defined in terms of the energy integrals as

$$ARE = \frac{1}{2 C U I_0}$$

The numerical integration performed in REIGEN uses a fourth order Runge-Kutta technique. Since T and C are already known, two independent solutions are integrated upward from the halfspace to the free surface and are combined in such a way that the surface conditions that $UZ = 1$ and $UR = e$ are met. If no numerical error occurs and the T and C pair are exact, then TZ and TR must equal zero at the free surface. When the proper combination of the two independent solutions is found, then the eigenfunctions are determined as a function of depth. Simultaneously, the energy integrals and phase velocity partial derivatives are formed.

Several parameters are available to test the numerical accuracy of the results. The group velocity $U = U(DC/DT)$, determined by SURFACE, can be compared to $U = U(ENER)$ determined in REIGEN. The Lagrangian L should be zero (this is usually a poor parameter to use for testing goodness of fit). The two amplitude factors AR (from SURFACE) and ARE (from REIGEN) should be equal. At the bottom of the listing, under $M = 0$, the values of TZ and TR , determined at the free surface in the process of satisfying the surface boundary conditions on UZ and UR , are displayed. These would be equal to zero if no

numerical error were encountered.

The partial derivatives for the last row are for changes in the parameters of the halfspace.

PROGRAM DESCRIPTION

PROGRAM REIGEN : This program performs the desired numerical integration and lists the results on the printer or tape.

SUBROUTINE TAPEIN : This subroutine makes it possible to update the eigenfunction tape.

INPUT DATA

There are two slightly different arrangements to the input data for the cases when the output is or is not written on FILE 01.

Case 1 - No output written on FILE 01.

Card	Sequence	Column	Name	Format	Explanation
A.		1-4	MMAX	I4	LE 0 End program
					GT 0 Number of layers in earth model including the halfspace
	11-14		ITAPE	I4	LE 0 No tape output (In this case, ITAPE must be LE 0)

B. Earth model (MMAX cards).

1-10	D(I)	F10.5	Layer thickness in km (Leave blank for halfspace, e.g. I=MMAX)
11-20	A(I)	F10.5	P velocity of layer in km/sec

Card Sequence	Column	Name	Format	Explanation
21-30	B(I)		F10.5	S velocity of layer in km/sec
31-40	RHO(I)		F10.5	Density of layer in gm/cm ³ .

C. Period data

1-14	T		E14.7	T LT 0 read new MMAX, ITAPE card (Point A).
				EQ 0 does nothing in this case except read in a new T, C, U, AR, RATIO card
				GT 0 Period in seconds
15-29	C		E14.7	Phase velocity in km/sec
31-44	U		E14.7	Group velocity in km/sec
46-59	AR		E14.7	Rayleigh wave amplitude factor
51-74	RATIO		E14.7	Rayleigh wave ellipticity

At this point the program will determine the eigenfunctions corresponding to the T,C pair and then return to Point C to read a new T,C,U,AR,ELLIP card.

The program is terminated by reading a MMAX card with MMAX.
LT.0. The last data card is followed by the system end of file card.

Case 2 - Output written on tape

A.	1-4	MMAX	I4	LE 0 End program
				GT 0 Number of layers including halfspace of the earth model

Card Sequence	Column	Name	Format	Explanation
	11-14	ITAPE	I4	LE 0 No tape output
				EQ 1 Earth model not written on FILE 01; eigen- functions are.
				GT 1 Earth model written as file header followed by eigenfunctions.
	21-30	FACT	F10.0	LE 0.0 use FACT = 4 GT 0.0 vertical wave lengths for layer dropping. See SURFACE.

B. Earth model (MMAX cards)

1-10	D(I)	F10.5	Layer thickness in km
11-20	A(I)	F10.5	P velocity in km/sec
21-30	B(I)	F10.5	S velocity in km/sec
31-40	RHO(I)	F10.5	Density of layer in gm/cm ³ .

C. Tape update

1-10	TMAX	F10.5	Last eigenfunction to be read when adding information to previously written eigenfunc- tion FILE.
11-15	KMODMX	I5	Number of earth models written on FILE to skip. KMODMX = 0 when writing new tape.

Card Sequence	Column	Name	Format	Explanation
D. Period Sequence				
a. Header				
	1-4	IFUNC	I4	LT 0 Write IFUNC on tape to denote end of eigenfunction list. Then terminate program.
				EQ 2 for all other cases IFUNC must equal 2. This signifies the FILE as containing the Rayleigh wave eigenfunctions.
	11-14	MODE	I4	Number of modes to be placed on tape for a given period.
	22-25	AMODE(I)	A4	Identification code for individual modes.
	27-20			
	.			
	67-70			
b. Period data (MODE cards)				
	1-14	T	E14.7	Period in seconds
	16-29	C	E14.7	Phase velocity in km/sec
	31-44	U	E14.7	Group velocity in km/sec
	46-59	AR	E14.7	Rayleigh wave amplitude factor
	61-74	RATIO	E14.7	Rayleigh wave ellipticity
c. End period data				
	1-80	blank card		

At this point the program returns to Point D to read a new IFUNC, MODE card. When the program finds an IFUNC.LT.0 , the program terminates.

NOTE: When the tape created here is to be used by other programs such as WIGGLE, SRFWVPLT, QUESTION, or RADPAT, the order of data is as is in the accompanying sample data set. The data are arranged in order of increasing period. Within each period set, the data are arranged in order of increasing mode, i.e., FUND first, 1 ST second, etc.

When using QUESTION and some of the other programs, both a LOVE and RAYLEIGH eigenfunction tape is used. Both tapes must have the same earth model for these other programs to work.

CREIGEN

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C      PROGRAM REIGEN
C      THE PROGRAM WILL ACCEPT UP TO ONE LIQUID LAYER AT THE SURFACE IN
C      WHICH CASE THE ELLIPTICITIES ARE THOSE AT THE TOP OF THE SOLID
C      LAYER ARRAY.
C      THIS PROGRAM COMPUTES VALUES OF DISPLACEMENT, STRESS AT MIDPOINT
C      OF EACH LAYER TOGETHER WITH VALUES OF DC/DB FOR THAT LAYER .
C      THE METHOD IS BASED ON A RUNGE-KUTTA INTEGRATION WITHIN EACH LAYER
C      THE INTEGRATION INTERVAL IS DDZ=D(I)/12., GOOD RESULTS WILL BE
C      OBTAINED WHEN DDZ IS LESS THAN C*T/2 PI, THUS, D(I)= 5 KM WILL
C      GIVE RESULTS TO 1 SEC FOR SEDIMENTARY LAYERS.
C      THE PROGRAM ALSO YIELDS VALUES OF VARIATIONAL INTEGRALS.
C      NUMERICAL ACCURACY IMPROVEMENT BY CHIUNG-CHUAN CHENG AUG 78
C      DOUBLE PRECISION YY1,YY2,YY3,YY4,YZ1,YZ2,YZ3,YZ4,AUR1,AUZ1,
1 ATZ1,ATR1,AUR2,AUZ2,ATZ2,ATR2,AA*BB,DAUR1,DAUZ1,DATZ1,DATR1,
1 DAUR2,DAUZ2,DATZ2,DATR2,EAUR1,EAUZ1,EATZ1,EATR1,EAUR2,EAUZ2,
1 EATZ2,EATR2,SAUR1,SAUZ1,SATZ1,SATR1,SAUR2,SAUZ2,SATZ2,SATR2,XNORM
COMMON D(100),A(100),B(100),RHO(100),DEPTH(100),AMPUR(100)
COMMON AMPUZ(100),STRESR(100),STRESSZ(100),DCDB(100)
COMMON C,U,T,FLAGR,SUMI0,SUMI1,SUMI2
COMMON SUMI3,UGR,WVNUM,DMR(5),DMZ(5),SMR(5),SMZ(5)
COMMON DMRSMZ,RATIO,DMZSMR,XLAMB(100),XMU(100)
COMMON AR,ARE,DCDA(100),DCDR(100)
COMMON WWT(4),WT(4)
COMMON YY1,YY2
COMMON AMODE(10)
DIMENSION YY1(100,5),YY2(100,5),YY3(100,5),YY4(100,5)
DIMENSION YZ1(100,5),YZ2(100,5),YZ3(100,5),YZ4(100,5)
DIMENSION DMRSMZ(5),DMZSMR(5)
11 FORMAT (2(14.6X),10(1X,A4))
20 FORMAT (1H1)
21 FORMAT (1H0,13X,2HM ,8X,7HDEPTH ,10X,1HD,12X,5HALPHA,11X,4HBETA,
1 11X,3HRHO,13X,2HMU,11X,5HLAMDA/1H )
23 FORMAT (1H ,10X,I5,7(5XF10.4))
25 FORMAT (1H ,10X,I5,30X,5(5X,F10.4))
44 FORMAT (4F10.5)
5000 FORMAT (1H1,13X,10H T(SEC) = ,F7.2,13H C(KM/SEC) = ,F7.4)
5001 FORMAT (1H0,13X,12H U(DG/DT) = ,F7.4,11H U(VAR) = ,F7.4,3X,
1 6H E = ,E11.4)
5002 FORMAT (1H0,13X,6H I0 = ,E11.4,3X,6H I1 = ,E11.4,3X,6H I2 = ,
1 E11.4)
5003 FORMAT (1H ,13X,6H I3 = ,E11.4,3X,6H L = ,E11.4,3X,6H K = ,
1 E11.4)
5004 FORMAT (1H0,3H M,5X,5HDEPTH,5X,2HUR,10X,2HUZ,10X,2HTZ,10X,2HTR,
1 9X,4HDCDA,8X,4HDCDB,8X,4HDCDR,/1H0)
5005 FORMAT (1H ,13,F10.2, 7(1X,E11.4))
5007 FORMAT (1H ,13X,6H AR = ,E11.4,3X,6HARE = ,E11.4)
6010 FORMAT (1H ,A4,3F10.5,5E11.4)
6011 FORMAT (1H ,7(4X,E11.4))
6012 FORMAT (1H ,2(14.6X))

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6013 FORMAT (1H ,I4,5X,7F10.4)
XXMIN = 1.0E-15
WWT(1)=0.0
WWT(2)=0.5
WWT(3)=0.5
WWT(4)=1.0
WT(1)=1./6.
WT(2)=1./3.
WT(3)=1./3.
WT(4)=1./6.
801 READ (60,12) MMAX,ITAPE,FACT
12 FORMAT (2(I4,6X),F10.0)
IF (FACT.LE.0.) FACT=4.
NMAX=MMAX
IF (MMAX) 777,777,778
778 WRITE (61,20)
WRITE (61,21)
BASE = 0.0
DO 22 I=1,MMAX
READ (60,44) D(I),A(I),B(I),RHO(I)
BASE = BASE + D(I)
DEPTH(I) = BASE - D(I)*0.5
XMU(I) = RHO(I)*B(I)*B(I)
XLAMB(I) = RHO(I)*(A(I)*A(I)-2.*B(I)*B(I))
IF (I = MMAX) 26,24,24
26 WRITE (61,23) I,DEPTH(I),D(I),A(I),B(I),RHO(I),XMU(I),XLAMB(I)
22 CONTINUE
24 WRITE (61,25) I,A(I),B(I),RHO(I),XMU(I),XLAMB(I)
DEPTH(MMAX) = BASE - D(MMAX)
IF (ITAPE.GT.0) CALL TAPEIN
IF (ITAPE.LE.1) GO TO 504
LMAX= MMAX+1
IF (B(1).LE.0.0) LMAX = LMAX - 1
WRITE (01,6012) LMAX
DPTH = 0.0
IF (B(1).LE.0.0) DPTH = D(1)
DO = 0.0
K = 1
IF (B(1).LE.0.0) K = 2
I = K - 1
WRITE (01,6013) I,DPTH,DO,A(K),B(K),RHO(K),XMU(K),XLAMB(K)
DO 200 I=1,MMAX
IF (B(I).LE.0.0) GO TO 200
WRITE (01,6013) I,DEPTH(I),D(I),A(I),B(I),RHO(I),XMU(I),XLAMB(I)
200 CONTINUE
504 IF (ITAPE.LE.0) GO TO 503
READ (60,11) IFUNC,MODE,(AMODE(I)*I=1,10)
WRITE (01,6012) IFUNC,MODE
IF (IFUNC.LT.0) GO TO 777
IIK = 0

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```

503 CONTINUE
  DO 500 J=1,MMAX
    DCDA(J)=0.0
    DCDB(J)=0.0
    DCDR(J)=0.0
    AMPUR(J)=0.0
    AMPUZ(J)=0.0
    STRESR(J)=0.0
    STRESZ(J)=0.0
500 CONTINUE
  READ (60,5) T,C,U,AR,RATIO
  5 FORMAT (5(E14.7,1X))
  IF (T) 801,504,499
499 CONTINUE
  MMAX = NMAX
C
C      LAYER DROPPING PROCEDURE--HALFSPACE DEFINED AS FACT WAVELENGTHS
C      BELOW REGION WHERE C IS LESS THAN B.
C
  DMAX = FACT*T*C
899 SUM = 0.0
  DO 900 II=1,MMAX
    MAX = II
    IF (C-B(II)) 901,900,900
901 SUM= SUM + D(II)
    IF (SUM.LE.DMAX) GO TO 900
    IF (II.EQ.MMAX) GO TO 902
    IF (A(II+1)-A(II)) 902,903,902
903 IF (B(II+1)-B(II)) 902,900,902
900 CONTINUE
902 MMAX = MAX
  SUMI0 = 0.0
  SUMI1 = 0.0
  SUMI2 = 0.0
  SUMI3 = 0.0
  WVNO = 6.2831853072/(C*T)
  WVNOSG=WVNO*WVNO
  WVNUM = WVNO
  OMEGA = 6.2831853072/T
  OMEGSG = OMEGA*OMEGA
  IF (B(1).GT.0.0) GO TO 4000
  RA = C/A(1)
  RA = WVNO * SQRT(RA*RA - 1.)
  IF (RA) 4001,4001,4002
4001 SUMI0 = RHO(1) * D(1)
  SUMI1 = 0.0
  SUMI2 = 0.0
  SUMI3 = 0.0
  TZZ= 0.0
  GO TO 4000

```

```

4002 SIN2RA = SIN(2.*RA*D(1))/(4.*RA)
COSRA = COS(RA*D(1))
COS2RM = 1./(COSRA*COSRA)
FAC1 = (0.5*D(1) + SIN2RA) * COS2RM
FAC3 = WVNO * ( 0.5 * D(1) - SIN2RA) * COS2RM
FAC2 = WVNO * FAC3 / (RA*RA)
FAC4 = RA * RA * FAC1
SUMI0 = RHO(1) * (FAC1+FAC2)
SUMI1 = XLAMB(1) * FAC2
SUMI2 = XLAMB(1) * FAC3
SUMI3 = XLAMB(1) * FAC4
TZZ = - RHO(1) * OMEGSQ * SIN(RA*D(1))/(RA*COSRA)
4000 CONTINUE
COVA = C/A(MMAX)
COVB = C/B(MMAX)
GAM = 2./COVB**2
GAMM1 = GAM - 1.
RA = WVNO*SQRT(ABS(COVA**2-1.))
RB = WVNO*SQRT(ABS(COVB**2-1.))
DET = WVNOSQ - RA * RB
H = RHO(MMAX) * OMEGSQ
BRKT = - GAMM1*WVNO + GAM*RA*RB/WVNO
C THE SUBSCRIPTS .1. AND .2. REFER TO TWO INDEPENDENT SOLUTIONS
ITER = 0
AUR1 = 1.D 00
AUZ1 = 0.D 00
ATZ1 = -H*BRKT/DET
ATR1 = -H*RA/DET
MMM1 = MMAX - 1
DO 1346 MM=1,MMM1
M = MMAX - MM
IF (B(M),LE.0.0) GO TO 1346
C CHOICE OF EFFICIENT VALUE OF DDZ FOR R-K INTEGRATION. WVNO*DDZ
C SHOULD GIVE INTEGRATION ERROR LESS THAN 0.01 PERCENT FOR ALL
C PERIODS.
XDIV=1.
NDIV = 1
DDZ = - D(M) / (4.*XDIV)
A12 = 1./(XLAMB(M)+2.*XMU(M))
A13 = WVNO * XLAMB(M) * A12
A21 = - OMEGSQ * RHO(M)
A24 = WVNO
A31 = -WVNO
A34 = 1./ XMU(M)
A42 = -A13
A43 = A21 + 4.*WVNOSQ*XMU(M)*(XLAMB(M)+XMU(M))*A12
YY3(M,5) = AUR1
YY1(M,5) = AUZ1
YY2(M,5) = ATZ1
YY4(M,5) = ATR1

```

```

DO 1338 KK=2,5
K = 6 - KK
C RUNGE-KUTTA INTEGRATION
DO 1400 JJ=1,NDIV
EAUR1 = AUR1
EAUZ1 = AUZ1
EATZ1 = ATZ1
EATR1 = ATR1
DAUR1 = 0.
DAUZ1 = 0.
DATZ1 = 0.
DATTR1 = 0.
DO 1401 LL = 1,4
C CLASSIC RUNGE KUTTA
SAUR1 = AUR1 + WWT(LL)*DDZ*DAUR1
SAUZ1 = AUZ1 + WWT(LL)*DDZ*DAUZ1
SATZ1 = ATZ1 + WWT(LL)*DDZ*DATZ1
SATR1 = ATR1 + WWT(LL)*DDZ*DATTR1
DAUR1 = A31*SAUZ1 + A34*SATR1
DAUZ1 = A12*SATZ1 + A13*SAUR1
DATZ1 = A21*SAUZ1 + A24*SATR1
DATTR1 = A42*SATZ1 + A43*SAUR1
EAUR1 = EAUR1 + WT(LL)*DDZ*DAUR1
EAUZ1 = EAUZ1 + WT(LL)*DDZ*DAUZ1
EATZ1 = EATZ1 + WT(LL)*DDZ*DATZ1
EATR1 = EATR1 + WT(LL)*DDZ*DATTR1
1401 CONTINUE
AUR1 = EAUR1
AUZ1 = EAUZ1
ATZ1 = EATZ1
ATR1 = EATR1
1400 CONTINUE
YY1(M,K) = AUZ1
YY2(M,K) = ATZ1
YY3(M,K) = AUR1
YY4(M,K) = ATR1
1338 CONTINUE
1346 CONTINUE
IF (B(1),GT.0,0) GO TO $347
YY1(1,1) = YY1(2,1)
YY2(1,1) = YY2(2,1)
YY3(1,1) = YY3(2,1)
YY4(1,1) = YY4(2,1)
1347 CONTINUE
4003 AUR2 = 0.0 00
AUZ2 = 1.0 00
ATZ2 = -H*RB/DET
ATR2 = -H*BRKT/DET
IF (ITER,EQ,0) GO TO 4004
AUR1 = 1.0 00

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AUZ1 = 0.D 00
ATZ1 = -H*BRKT/DET
ATR1 = -H*RA/DET
AUR2 = XNORM*AUR1 + AUR2
AUZ2 = XNORM*AUZ1 + AUZ2
ATZ2 = XNORM*ATZ1 + ATZ2
ATR2 = XNORM*ATR1 + ATR2
4004 DO 2346 MM=1,MMM1
M = MMAX - MM
IF (B(M).LE.0.0) GO TO 2346
DDZ = - D(M) / (4.*XDIV)
A12 = 1./(XLAMB(M)+2.*XMU(M))
A13 = WVNO * XLAMB(M) * A12
A21 = - OMEGSQ * RHO(M)
A24 = WVNO
A31 = -WVNO
A34 = 1./ XMU(M)
A42 = -A13
A43 = A21 + 4.*WVNOSQ*XMU(M)*(XLAMB(M)+XMU(M))*A12
YZ3(M,5) = AUR2
YZ1(M,5) = AUZ2
YZ2(M,5) = ATZ2
YZ4(M,5) = ATR2
DO 2338 KK=2,5
K = 6 - KK
DO 2400 JJ=1,NDIV
EAUR2 = AUR2
EAUZ2 = AUZ2
EATZ2 = ATZ2
EATR2 = ATR2
DAUR2 = 0.
DAUZ2 = 0.
DATZ2 = 0.
DATR2 = 0.
DO 2401 LL = 1,4
SAUR2 = AUR2 + WWT(LL)*DDZ*DAUR2
SAUZ2 = AUZ2 + WWT(LL)*DDZ*DAUZ2
SATZ2 = ATZ2 + WWT(LL)*DDZ*DATZ2
SATR2 = ATR2 + WWT(LL)*DDZ*DATTR2
DAUR2 = A31*SAUZ2 + A34*SATR2
DAUZ2 = A12*SATZ2 + A13*SAUR2
DATZ2 = A21*SAUZ2 + A24*SATR2
DATTR2 = A42*SATZ2 + A43*SAUR2
EAUR2 = EAUR2 + WT(LL)*DDZ*DAUR2
EAUZ2 = EAUZ2 + WT(LL)*DDZ*DAUZ2
EATZ2 = EATZ2 + WT(LL)*DDZ*DATZ2
EATR2 = EATR2 + WT(LL)*DDZ*DATTR2
2401 CONTINUE
AUR2 = EAUR2
AUZ2 = EAUZ2

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        ATZ2 = EATZ2
        ATR2 = EAATR2
2400  CONTINUE
        YZ1(M,K) = AUZ2
        YZ2(M,K) = ATZ2
        YZ3(M,K) = AUR2
        YZ4(M,K) = ATR2
2338  CONTINUE
2346  CONTINUE
        IF (B(1).GT.0.0) GO TO 2347
        YZ1(1,1) = YZ1(2,1)
        YZ2(1,1) = YZ2(2,1)
        YZ3(1,1) = YZ3(2,1)
        YZ4(1,1) = YZ4(2,1)
2347  CONTINUE
        AA = (YZ3(1,1) - RATIO*YZ1(1,1))
        BB = (RATIO*YY1(1,1) - YY3(1,1))
        XNORM = AA/BB
        BB = XNORM*YY1(1,1) + YZ1(1,1)
C      THESE PARAMETERS ARE SAVED FOR USE AS A NUMERICAL ANALYSIS
C      DIAGNOSTIC
        AMPUR(100) = (XNORM*YY3(1,1) + YZ3(1,1))/BB
        AMPUZ(100) = (XNORM*YY1(1,1) + YZ1(1,1))/BB
        STRESZ(100) = (XNORM*YY2(1,1) + YZ2(1,1))/BB
        STRESR(100) = (XNORM*YY4(1,1) + YZ4(1,1))/BB
        ITER = ITER + 1
        IF (ITER.GT.1) GO TO 2011
C      IF ROUNDOFF ERROR IS SIGNIFICANT, MODIFY INITIAL CHOICE AT BASE
C      OF LAYERS OF ONE OF THE INDEPENDENT SOLUTIONS
        XTEST = ABS(AMPUR(100)/RATIO-1.)
        IF (XTEST.GE.0.00001) GO TO 4003
2011  DCDB(99) = ITER
        DCDA(100) = 0.0
        DCDB(100) = 0.0
        DCDR(100) = 0.0
        AMPUR(99) = YY3(1,1)
        AMPUR(98) = YZ3(1,1)
        AMPUZ(99) = YY1(1,1)
        AMPUZ(98) = YZ1(1,1)
        STRESZ(99) = YY2(1,1)
        STRESZ(98) = YZ2(1,1)
        STRESR(99) = YY4(1,1)
        STRESR(98) = YZ4(1,1)
        DCDB(98) = BB
        DCDA(98) = 0.0
        DCDA(99) = 0.0
        DCDR(98) = 0.0
        DCDR(99) = 0.
C      COMBINATION OF TWO INDEPENDENT SOLUTIONS SUBJECT TO SURFACE R. C.
C      THAT UR/UZ = RATIO (I)

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C
DO 7000 M=1,MMAX
IF (B(M),LE,0.0) GO TO 7000
IF (M-MMAX) 7001,7002,7002
7001 DZ = D(M)/4.
DO 1339 KK=1,5
AUR = (XNORM*YY3(M,KK)      + YZ3(M,KK)) /BB
AUZ = (XNORM*YY1(M,KK)      + YZ1(M,KK)) /BB
ATZ = (XNORM*YY2(M,KK)      + YZ2(M,KK)) /BB
ATR = (XNORM*YY4(M,KK)      + YZ4(M,KK)) /BB
DURDZ = ATR/XMU(M) - WVNO*AUZ
DUZDZ = (ATZ + WVNO*XLAMB(M)*AUR)/(XLAMB(M)+2.*XMU(M))
DMR(KK) = AUR*AUR
DMZ(KK) = AUZ*AUZ
SMR(KK) = DURDZ*DURDZ
SMZ(KK) = DUZDZ*DUZDZ
DMRSMZ(KK) = AUR*DUZDZ
DMZSMR(KK) = AUZ*DURDZ
IF (KK - 3) 1339,1301,1339
C EIGENFUNCTIONS AT MIDPOINT OF LAYER ARE SAVED
1301 AMPUR(M) = AUR
AMPUZ(M) = AUZ
STRESZ(M) = ATZ
STRESR(M) = ATR
1339 CONTINUE
C INTEGRATION BY BOODE'S RULE
DMMR = (DZ/22.5)*(7.*(DMR(1)+DMR(5))+32.*DMR(2)+DMR(4))+12,
1 *DMR(3))
DMMZ = (DZ/22.5)*(7.*(DMZ(1)+DMZ(5))+32.*DMZ(2)+DMZ(4))+12,
1 *DMZ(3))
SMMZ = (DZ/22.5)*(7.*(SMZ(1)+SMZ(5))+32.*SMZ(2)+SMZ(4))+12,
1 *SMZ(3))
SMMR = (DZ/22.5)*(7.*(SMR(1)+SMR(5))+32.*SMR(2)+SMR(4))+12,
1 *SMR(3))
DRSZ = (DZ/22.5)*(7.*(DMRSMZ(1)+DMRSMZ(5))+32.*DMRSMZ(2)
1 +DMRSMZ(4))+12.*DMRSMZ(3))
DZSR = (DZ/22.5)*(7.*(DMZSMR(1)+DMZSMR(5))+32.*DMZSMR(2)
1 +DMZSMR(4))+12.*DMZSMR(3))
SUMI0 = SUMI0 + RHO(M)*(DMMR+DMMZ)
SUMI1 = (XLAMB(M)+2.*XMU(M))*DMMR+XMU(M)*DMMZ + SUMI1
SUMI2 = XMU(M)*DZSR - XLAMB(M)*DRSZ + SUMI2
SUMI3 = (XLAMB(M)+2.*XMU(M))*SMMZ + XMU(M)*SMMR + SUMI3
DLDL = -WVNOSQ*DMMR + 2.*WVNO*DRSZ - SMMZ
DLDM = -WVNOSQ*(2.*DMMR*DMMZ) - 2.*WVNO*DZSR - (2.*SMMZ+SMMR)
DLDR = OMEGSQ*(DMMR+DMMZ)
DCDB(M) = 2.*RHO(M)*B(M)*C * (DLDM-2.*DLDL)/WVNO
DCDA(M) = 2.*RHO(M)*A(M)*C * DLDL/WVNO
DCDR(M) = (C/WVNO)*(DLDR+XLAMB(M)*DLDE/RHO(M)*XMU(M)*DLDM/RHO(M))
IF (ABS(AUZ)+ABS(AUR)=XXMIN) 7002*7002,7000
7000 CONTINUE

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C      EXACT HALFSPACE CONTRIBUTION
7002 AMPUR(M) = AUR
      AMPUZ(M) = AUZ
      STRESR(M) = ATR
      STRESZ(M) = ATZ
      AP = -RHO(M)*(WVNO*AUR + RB*AUZ)/DET
      BP = -RHO(M)*(-RA*AUR/WVNO - AUZ)/DET
      A1 = -WVNO*AP/RHO(M)
      A2 = -WVNO*RB*BP/RHO(M)
      A3 = RA*AP/RHO(M)
      A4 = WVNO*BP/RHO(M)
      IF (RB) 7005,7006,7005
7005 DMMR = A1*A1/(2.*RA)+2.*A1*A2/(RA+RB)+A2*A2/(2.*RB)
      DMMZ = A3*A3/(2.*RA)+2.*A3*A4/(RA+RB)+A4*A4/(2.*RB)
      SMMZ = RA*A3*A3/2.+2.*RA*RB*A3*A4/(RA+RB)+RB*A4*A4/2.
      SMMR = RA*A1*A1/2.+2.*RA*RB*A1*A2/(RA+RB)+RB*A2*A2/2.
      DRSZ = -A1*A3/2,-(A1*A4*RB+A2*A3*RA)/(RA+RB)-A2*A4/2.
      DZSR = -A1*A3/2,-(A1*A4*RA+A2*A3*RB)/(RA+RB)-A2*A4/2.
      GO TO 7010
7006 UGR = B(M)
      FLAGR = 0,0
      SUMI0 = RHO(M)*10.**25
      SUMI1 = XMU(M)*10.**25
      SUMI2 = 0.0
      SUMI3 = 0.0
      ARE = 0.0
      DCDB(M) = -2.*WVNO*10.**25
      GO TO 531
7010 CONTINUE
      SUMI0 = SUMI0 + RHO(M)*(DMMR+DMMZ)
      SUMI1 = (XLAMB(M)+2.*XMD(M))*DMMR+XMU(M)*DMMZ + SUMI1
      SUMI2 = XMU(M)*DZSR - XLAMB(M)*DRSZ + SUMI2
      SUMI3 = (XLAMB(M)+2.*XMD(M))*SMMZ + XMU(M)*SMMR + SUMI3
      DLDR = OMEGSQ*(DMMR+DMMZ)
      DLDM = -WVNO*SQ*(2.*DMMR+DMMZ)-2.*WVNO*DZSR-(2.*SMMZ+SMMR)
      DLDL = -WVNO*SQ*DMMR + 2.*WVNO*DRSZ - SMMZ
      DCDA(M) = 2.*RHO(M)*A(M)*C*DLDL/WVNO
      DCDB(M) = 2.*RHO(M)*B(M)*C*(DLDM-2.*DLDL)/WVNO
      DCDR(M) = (C/WVNO)*(DLDR+XLAMB(M)*DLDL/RHO(M)+XMU(M)*DLDM/RHO(M))
7011 CONTINUE
      UGR = (WVNO*SUMI1 + SUMI2)/(OMEGA*SUMI0)
      FLAGR = OMEGSQ*SUMI0 - WVNO*SQ*SUMI1 - 2.*WVNO*SUMI2 - SUMI3
      ARE = 1./((2.*C*UGR*SUMI0))
531 CONTINUE
501 CONTINUE
      WRITE (61,5000) T,C
      WRITE (61,5001) U,UGR,RATIO
      WRITE (61,5002) SUMI0,SUMI1,SUMI2
      WRITE (61,5003) SUMI3,FLAGR,WVNUM
      WRITE (61,5007) AR,ARE

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        WRITE (61,5004)
AU = RATIO
AZ = 1.0
TR = 0.0
TZ = 0.0
IF (B(1).LE.0.0) TZ=TZZ
DEPTH(100)=0.0
IF (B(1).LE.0.0) DEPTH(100)=D(1)
M = 0
DUMM1 = 0.0
DUMM2 = 0.0
DUMM3 = 0.0
WRITE (61,5005) M,DEPTH(100),AU,AZ,TZ,TR,DUMM1,DUMM2,DUMM3
N = 1
IF (B(1).LE.0.0) N = 2
DO 5010 M = N,MMAX
DLDK = -2.* (WVNUM * SUMI1 + SUMI2)
DCDR(M) = DCDR(M)/DLDK
DCDA(M) = DCDA(M)/DLDK
5009 DCDB(M) = DCDB(M)/DLDK
5010 WRITE (61,5005) M,DEPTH(M),AMPUR(M),AMPUZ(M),STRESZ(M),STRESR(M),
1 DCDA(M),DCDB(M),DCDR(M)
WRITE (61,19)
19 FORMAT (1H0)
M = 100
WRITE (61,5005) M,DEPTH(M),AMPUR(M),AMPUZ(M),STRESZ(M),STRESR(M),
1 DCDA(M),DCDB(M),DCDR(M)
WRITE (61,19)
M = 98
WRITE (61,5005) M,DEPTH(M),AMPUR(M),AMPUZ(M),STRESZ(M),STRESR(M),
1 DCDA(M),DCDB(M),DCDR(M)
M = 99
WRITE (61,5005) M,DEPTH(M),AMPUR(M),AMPUZ(M),STRESZ(M),STRESR(M),
1 DCDA(M),DCDB(M),DCDR(M)
IF (ITAPE.LE.0) GO TO 503
MMAX = NMAX
K = IIK +1
WRITE (01,6010) AMODE(K),T,C,U,SUMI0,WVNUM,AR,RATIO,DUMM1
WRITE (01,6011) AU,AZ,TZ,TR,DUMM1,DUMM2,DUMM3
N = 1
IF (B(1).LE.0.0) N = 2
DO 6002 M= N,MMAX
WRITE (01,6011) AMPUR(M),AMPUZ(M),STRESZ(M),STRESR(M),DCDA(M),
1 DCDB(M),DCDR(M)
6002 CONTINUE
IIK = IIK+ 1
GO TO 503
777 CONTINUE
STOP
END

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```
SUBROUTINE TAPEIN
READ (60,1) TMAX,KMODMX
1 FORMAT(F10.5,I5)
KMODEL = 0
3002 CONTINUE
IF (KMODEL.EQ.KMODMX) RETURN
READ (01,6012) MMAX
6012 FORMAT (1H ,2(I4,6X))
DO 200 I = 1,MMAX
200 READ (01,6013) IDUM
6013 FORMAT (1H ,I4,5X,7F10.4)
2000 READ (01,6012) IFUNC,MODE
IF (IFUNC.LT.0) GO TO 3000
DO 2999 IM = 1,MODE
READ (01,6010) IDUM,T
6010 FORMAT (1H ,A4,3F10.5,5E11.4)
DO 2100 IN = 1,MMAX
2100 READ (01,6011) XDUM
6011 FORMAT (1H ,7(4X,E11.4))
2999 CONTINUE
IF (T.GE.TMAX) RETURN
GO TO 2000
3000 KMODEL = KMODEL + 1
GO TO 3002
END
```

31	2		
1.0	5.00	2.89	2.5
1.0	6.1	3.52	2.7
1.0	6.1	3.52	2.7
1.0	6.1	3.52	2.7
1.0	6.1	3.52	2.7
1.0	6.1	3.52	2.7
1.0	6.1	3.52	2.7
1.0	6.1	3.52	2.7
1.0	6.1	3.52	2.7
1.0	6.1	3.52	2.7
1.0	6.1	3.52	2.7
1.0	6.1	3.52	2.7
1.0	6.1	3.52	2.7
1.0	6.4	3.7	2.9
1.0	6.4	3.7	2.9
1.0	6.4	3.7	2.9
1.0	6.4	3.7	2.9
1.0	6.4	3.7	2.9
1.0	6.4	3.7	2.9
1.0	6.4	3.7	2.9
1.0	6.4	3.7	2.9
1.0	6.4	3.7	2.9
1.0	6.4	3.7	2.9
1.0	6.4	3.7	2.9
2.0	6.7	3.87	3.0
2.0	6.7	3.87	3.0
2.0	6.7	3.87	3.0
2.0	6.7	3.87	3.0
2.0	6.7	3.87	3.0
2.0	6.7	3.87	3.0
2.0	6.7	3.87	3.0
2.0	6.7	3.87	3.0
2.0	6.7	3.87	3.0
2.0	6.7	3.87	3.0
2.0	6.7	3.87	3.0
2.0	6.7	3.87	3.0
2.0	6.7	3.87	3.0
2.0	6.7	3.87	3.0
2.0	8.15	4.7	3.4

2	3	FUND 1 ST 2 ND		
2.0000000E 00	3.1142651E 00	3.0052558E 00	6.6519809E-03	7.7533051E-01
2.0000000E 00	3.7021885E 00	3.5109269E 00	2.0487942E-04	6.2238460E-01
2.0000000E 00	3.8652486E 00	3.6090615E 00	1.4074941E-04	5.6639548E-01

2	3	FUND 1 ST 2 ND		
2.5000000E 00	3.1358773E 00	3.0523695E 00	4.9933391E-03	7.9398409E-01
2.5000000E 00	3.7487547E 00	3.5414947E 00	1.6791463E-04	6.2641879E-01
2.5000000E 00	3.9224432E 00	3.6695645E 00	8.6686953E-05	5.6405590E-01

2	3	FUND 1 ST 2 ND		
3.0000000E 00	3.1511676E 00	3.0684696E 00	4.0669670E-03	8.0034681E-01
3.0000000E 00	3.7910240E 00	3.5610245E 00	1.5263218E-04	6.1663947E-01
3.0000000E 00	3.9830012E 00	3.6056533E 00	9.9571947E-05	5.4424362E-01

2	3	FUND 1 ST 2 ND		
3.5000000E 00	3.1650985E 00	3.0703514E 00	3.4696805E-03	8.0017772E-01

3.5000000E 00	3.8306909E 00	3.5773676E 00	1.4386220E-04	6.0192247E-01
3.5000000E 00	4.0642072E 00	3.4999735E 00	1.3432320E-04	5.0977601E-01
2	3	FUND 1 ST 2 ND		
4.0000000E 00	3.1793620E 00	3.0665456E 00	3.0437626E-03	7.9666199E-01
4.0000000E 00	3.8692162E 00	3.5826478E 00	1.4036134E-04	5.8613508E-01
4.0000000E 00	4.1702411E 00	3.4017723E 00	1.7067032E-04	4.6229274E-01
2	3	FUND 1 ST 2 ND		
5.0000000E 00	3.2100163E 00	3.0587665E 00	2.4481107E-03	7.8600404E-01
5.0000000E 00	3.9516559E 00	3.5479480E 00	1.5024156E-04	5.5422457E-01
5.0000000E 00	4.4235663E 00	3.4890519E 00	1.6888031E-04	3.4403332E-01
2	3	FUND 1 ST 2 ND		
6.0000000E 00	3.2420740E 00	3.0607984E 00	2.0292056E-03	7.7537596E-01
6.0000000E 00	4.0523164E 00	3.4753120E 00	1.7317689E-04	5.1928874E-01
6.0000000E 00	4.5971306E 00	3.9840677E 00	8.6218279E-05	2.6498797E-01
2	3	FUND 1 ST 2 ND		
7.0000000E 00	3.2734813E 00	3.0721896E 00	1.7125091E-03	7.6686421E-01
7.0000000E 00	4.1752123E 00	3.4192479E 00	1.9103228E-04	4.7781309E-01
7.0000000E 00	4.6848377E 00	4.3021607E 00	4.7082233E-05	2.3372602E-01
2	2	FUND 1 ST 2 ND		
8.0000000E 00	3.3031600E 00	3.0884357E 00	1.4674579E-03	7.4071229E-01
8.0000000E 00	4.3090794E 00	3.4664272E 00	1.8146382E-04	4.3075433E-01
2	2	FUND 1 ST 2 ND		
9.0000000E 00	3.3309471E 00	3.1049636E 00	1.2764350E-03	7.5653140E-01
9.0000000E 00	4.4275274E 00	3.6620983E 00	1.4106098E-04	3.8617375E-01
2	2	FUND 1 ST 2 ND		
1.0000000E 01	3.3572350E 00	3.1180076E 00	1.1269633E-03	7.5379098E-01
1.0000000E 01	4.5134921E 00	3.9041931E 00	9.8747813E-05	3.5185762E-01
2	2	FUND 1 ST 2 ND		
1.1000000E 01	3.3826617E 00	3.1261866E 00	1.0092459E-03	7.5200669E-01
1.1000000E 01	4.5722072E 00	4.0912724E 00	7.0955873E-05	3.2831591E-01
2	2	FUND 1 ST 2 ND		
1.2000000E 01	3.4078975E 00	3.1287238E 00	9.1556588E-04	7.5080240E-01
1.2000000E 01	4.6141496E 00	4.2182069E 00	5.4338796E-05	3.1348770E-01
2	2	FUND 1 ST 2 ND		
1.3000000E 01	3.4335251E 00	3.1259829E 00	8.3977392E-04	7.4991863E-01
1.3000000E 01	4.6458775E 00	4.3101667E 00	4.3387340E-05	3.0606461E-01
2	2	FUND 1 ST 2 ND		
1.4000000E 01	3.4599863E 00	3.1193709E 00	7.7717032E-04	7.4919921E-01
1.4000000E 01	4.6703708E 00	4.3906449E 00	3.4092357E-05	3.0547285E-01

		FUND 1 ST 2 ND		
2	2	3.4875658E 00	3.1104207E 00	7.2408858E-04 7.4897095E-01
1.5000000E 01	4.6883101E 00	4.4814298E 00	2.3859875E-05 3.1153613E-01	
2	2	FUND 1 ST 2 ND		
1.6000000E 01	3.5163950E 00	3.1007019E 00	6.7787876E-04 7.4802391E-01	
1.6000000E 01	4.6985934E 00	4.6169773E 00	8.8747948E-06 3.2445508E-01	
2	1	FUND 1 ST 2 ND		
1.7000000E 01	3.5464619E 00	3.0924973E 00	6.3644259E-04 7.4759430E-01	
2	1	FUND 1 ST 2 ND		
1.8000000E 01	3.5776256E 00	3.0865256E 00	5.9829302E-04 7.4734917E-01	
2	1	FUND 1 ST 2 ND		
1.9000000E 01	3.6096355E 00	3.0847227E 00	5.6235467E-04 7.4737325E-01	
2	1	FUND 1 ST 2 ND		
2.0000000E 01	3.6421553E 00	3.0883168E 00	5.2798860E-04 7.4775627E-01	
2	1	FUND 1 ST 2 ND		
2.2000000E 01	3.7071341E 00	3.1146846E 00	4.6254308E-04 7.4991790E-01	
2	1	FUND 1 ST 2 ND		
2.4000000E 01	3.7693573E 00	3.1656269E 00	4.0148728E-04 7.5427168E-01	
2	1	FUND 1 ST 2 ND		
2.6000000E 01	3.8262323E 00	3.2363947E 00	3.4616044E-04 7.6084514E-01	
2	1	FUND 1 ST 2 ND		
2.8000000E 01	3.8763158E 00	3.3190566E 00	2.9794579E-04 7.6928231E-01	
2	1	FUND 1 ST 2 ND		
3.0000000E 01	3.9192900E 00	3.4042079E 00	2.5727347E-04 7.7900857E-01	
2	1	FUND 1 ST 2 ND		
3.5000000E 01	3.9994231E 00	3.5978714E 00	1.8459469E-04 8.0515936E-01	
2	1	FUND 1 ST 2 ND		
4.0000000E 01	4.0513421E 00	3.7414352E 00	1.4075543E-04 8.2907604E-01	
2	1	FUND 1 ST 2 ND		
4.5000000E 01	4.0861130E 00	3.8427755E 00	1.1322909E-04 8.4821653E-01	
2	1	FUND 1 ST 2 ND		
5.0000000E 01	4.1105629E 00	3.9130282E 00	9.4880703E-05 8.6238892E-01	
2	1	FUND 1 ST 2 ND		
6.0000000E 01	4.1425549E 00	3.9985009E 00	7.2337643E-05 8.7866213E-01	

2	1	FUND 1 ST 2 ND		
7.0000000E 01	4.1629857E 00	4.0475750E 00	5,9053952E-05	8,8405416E-01
2	1	FUND 1 ST 2 ND		
8.0000000E 01	4.1776894E 00	4.0773227E 00	5.0216392E-05	8,8319131E-01
2	1	FUND 1 ST 2 ND		
9.0000000E 01	4.1891447E 00	4.0990825E 00	4,3848633E-05	8,7890215E-01
2	1	FUND 1 ST 2 ND		
1.0000000E 02	4.1985384E 00	4.1142765E 00	3,8996830E-05	8,7283818E-01
2	1	FUND 1 ST 2 ND		
1.5000000E 02	4.2297525E 00	4.1626364E 00	2,5310945E-05	8,3814402E-01
2	1	FUND 1 ST 2 ND		
2.0000000E 02	4.2482172E 00	4.1896077E 00	1,8762922E-05	8,1026544E-01
-2	1	FUND		

II-25

U(DC/DT) = 3.0053 U(VAR) = 3.0052 E = 0.7753E 00
 10 = 0.8030E 01 11 = 0.1003E 03 12 = -0.2537E 02
 13 = 0.12837E 02 L = 0.1017E-03 K = 0.1009E 01
 AR = 0.66525E-02 ARE = 0.6653E-02

M	DEPTH	UR	UZ	TZ	TR	DCDA	DCDB	DCDR
0.	0.7753E 00	0.1000E 01	0.	0.1012E 02	0.9842E 01	0.9434E-01	0.9364E-01	-0.1055E 00
0.50	0.3791E 00	0.1050E 01	0.	0.1012E 02	0.2170E 02	0.2507E-01	0.3621E-01	-0.1569E 00
1.50	-0.1022E-01	0.9107E 00	0.	0.1012E 02	0.4424E-02	0.4170E 00	0.4341E-01	-0.4170E 00
2.50	-0.1482E 00	0.6907E 00	0.	0.1012E 02	0.7808E-03	0.1451E 00	0.1451E 00	-0.1451E 00
3.50	-0.1522E 00	0.4825E 00	0.	0.1012E 02	0.1380E-03	0.7710E-01	0.2602E 01	-0.2602E 01
4.50	-0.1203E 00	0.3225E 00	0.	0.1012E 02	0.1208E 02	0.1380E-03	0.7710E-01	-0.2602E 01
5.50	-0.8599E-01	0.2098E 00	0.	0.1012E 02	0.9478E 01	0.1208E 02	0.1208E 02	-0.1208E 02
6.50	-0.5852E-01	0.1337E 00	0.	0.1012E 02	0.6376E 01	0.6376E 01	0.2460E-04	-0.3614E-01
7.50	-0.3882E-01	0.8348E-01	0.	0.1012E 02	0.4198E 01	0.5296E 01	0.4600E-05	-0.1577E 01
8.50	-0.2517E-01	0.5059E-01	0.	0.1012E 02	0.2743E 01	0.3363E 01	0.103E-05	-0.6564E-02
9.50	-0.1507E-01	0.2932E-01	0.	0.1012E 02	0.1795E 01	0.2089E 01	0.6648E-06	-0.2634E-02
10.50	-0.8234E-02	0.1707E-01	0.	0.1012E 02	0.1165E 01	0.1301E 01	0.1604E-05	-0.1051E-02
11.50	-0.5064E-02	0.1017E-01	0.	0.1012E 02	0.4290E 00	0.5069E 00	0.1685E-06	-0.1377E 01
12.50	-0.3076E-02	0.6009E-02	0.	0.1012E 02	0.2554E 00	0.3029E 00	0.2892E-07	-0.1238E-03
13.50	-0.1837E-02	0.3531E-02	0.	0.1012E 02	0.1509E 00	0.1794E 00	0.516E-09	-0.1572E-04
14.50	-0.1087E-02	0.2067E-02	0.	0.1012E 02	0.8665E-01	0.1056E 00	0.1462E-09	-0.5465E-05
15.50	-0.6398E-03	0.1207E-02	0.	0.1012E 02	0.5189E-01	0.6187E-01	0.2508E-10	-0.1882E-05
16.50	-0.3748E-03	0.7034E-03	0.	0.1012E 02	0.3029E-01	0.3615E-01	0.4304E-11	-0.6435E-06
17.50	-0.2190E-03	0.4093E-03	0.	0.1012E 02	0.1765E-01	0.2108E-01	0.7382E-12	-0.9429E-07
18.50	-0.1276E-03	0.2380E-03	0.	0.1012E 02	0.1027E-01	0.1227E-01	0.1267E-12	-0.3202E-07
19.50	-0.9740E-04	0.1814E-03	0.	0.1012E 02	0.7831E-02	0.9358E-02	0.2616E-13	-0.3796E-07
20.	0.	0.7753E 00	0.1000E 01	0.	0.9053E-04	0.9526E-04	0.	0.
100	0.	0.7753E 00	0.1000E 01	0.	0.9053E-04	0.9526E-04	0.	0.
98	0.	0.2207E 08	-0.1802E 08	0.	0.5748E 09	-0.7375E 09	0.	0.1245E 10
99	0.	0.4110E 08	-0.3358E 08	0.	0.1065E 10	-0.1374E 10	0.	0.1000E 01

$U(DC/DT) = 3.5109$ $U(VAR) = 3.5108$ $E = 0.6224E 00$

$10 = 0.1876E 03$ $11 = 0.2653E 04$ $12 = 0.1816E 03$
 $13 = 0.2449E 03$ $L = 0.3397E-03$ $K = 0.8486E 00$
 $AR = 0.2049E-03$ $ARE = 0.2050E-03$

DEPTH	UR	UZ	TZ	TR	DCDA	DCDB	DCDR
0	0.6224E 00	0.1000E 01	0.	0.4102E 01	0.	0.2837E-02	0.7371E-03
0.50	0.2420E 00	0.1012E 01	-0.1165E 02	0.	0.7696E 01	0.8030E-02	0.7651E-03
1.50	0.2633E 00	0.7461E 00	-0.2878E 02	0.	0.7100E-03	0.7643E-02	0.9118E-02
2.50	-0.5523E 00	0.2860E 00	-0.3835E 02	0.	0.1572E 01	0.2012E-03	0.1366E-01
3.50	-0.6773E 00	-0.2952E 00	-0.4190E 02	-0.	0.1048E 02	0.4427E-04	0.3000E-01
4.50	-0.6907E 00	-0.9060E 00	-0.4074E 02	-0.	0.2469E 02	0.8925E-05	0.3922E-01
5.50	-0.6250E 00	-0.1478E 01	-0.3581E 02	-0.	0.3872E 02	0.6987E-06	0.4958E-01
6.50	-0.5061E 00	-0.1959E 01	-0.2802E 02	-0.	0.5103E 02	0.2213E-05	0.6116E-01
7.50	-0.3613E 00	-0.2314E 01	-0.1641E 02	-0.	0.6077E 02	0.1667E-04	0.7312E-01
8.50	-0.2257E 00	-0.2529E 01	-0.8340E 01	-0.	0.6793E 02	0.7460E-04	0.8490E-01
9.50	-0.1512E 00	-0.2618E 01	-0.3841E 00	-0.	0.7366E 02	0.2984E-05	0.9778E-01
10.50	-0.3557E-01	-0.2619E 01	-0.7734E 01	-0.	0.7987E 02	0.3078E-03	0.9080E-01
11.50	-0.1470E 00	-0.2515E 01	0.1296E 02	-0.	0.8052E 02	0.7355E-04	0.9393E-01
12.50	0.1942E 00	-0.2349E 01	0.1564E 02	-0.	0.7699E 02	0.1962E-04	0.8765E-01
13.50	0.2334E 00	-0.2150E 01	0.1704E 02	-0.	0.7135E 02	0.5168E-05	0.7674E-01
14.50	0.2529E 00	-0.1934E 01	0.1777E 02	-0.	0.6465E 02	0.1565E-05	0.6427E-01
15.50	0.2613E 00	-0.1711E 01	0.1811E 02	-0.	0.5745E 02	0.7364E-06	0.51953E-02
16.50	0.2621E 00	-0.1484E 01	0.1819E 02	-0.	0.5010E 02	0.8213E-06	0.4065E-01
17.50	0.2548E 00	-0.1258E 01	0.1799E 02	-0.	0.4288E 02	0.2011E-05	0.3083E-01
18.50	0.2351E 00	-0.1040E 01	0.1737E 02	-0.	0.3617E 02	0.6972E-05	0.2268E-01
19.50	0.1926E 00	-0.8378E 00	0.1597E 02	-0.	0.3060E 02	0.2683E-04	0.1631E-01
20.50	0.2110E 00	-0.6038E 00	0.1250E 02	-0.	0.2381E 02	0.1241E-04	0.1720E-01
21.50	0.1457E 00	-0.4350E 00	0.8232E 01	-0.	0.1545E 02	0.7332E-06	0.7566E-02
22.50	0.6673E-01	-0.2349E 00	0.5171E 01	-0.	0.9649E 01	0.4333E-07	0.2999E-02
23.50	0.4161E-01	-0.1439E 00	0.3193E 01	-0.	0.6935E 01	0.2560E-08	0.1143E-02
24.50	0.2566E-01	-0.8771E-01	0.1961E 01	-0.	0.3624E 01	0.1528E-09	0.4281E-03
25.50	0.1581E-01	-0.5316E-01	0.1207E 01	-0.	0.2198E 01	0.1045E-10	0.1587E-03
26.50	0.9824E-02	-0.3182E-01	0.7500E 00	-0.	0.1317E 01	0.4354E-11	0.5801E-04
27.50	0.6235E-02	-0.1644E-01	0.4780E 00	-0.	0.7645E 00	0.3972E-10	0.2056E-04
28.50	0.4050E-02	-0.9764E-02	0.3189E 00	-0.	0.4122E 00	0.6451E-09	0.6816E-05
29.50	0.2268E-02	-0.4033E-02	0.2142E 00	-0.	0.2012E 00	0.1089E-07	0.1943E-05
30.50	0.9364E-03	-0.2182E-02	0.1551E 00	-0.	0.1600E 00	0.9673E-09	0.3133E-06
31.50	0.4000E 00	-0.1000E 01	-0.1659E-03	-0.	0.6581E-04	0.	0.1514E-06
100	0.	0.6224E 00	0.1000E 01	-0.1659E-03	-0.6581E-04	0.	0.
98	0.	0.1294E 13	-0.9593E 12	0.2059E 14	-0.3308E 14	0.	0.7308E 13
99	0.	0.3015E 13	-0.2235E 13	0.4797E 14	-0.7707E 14	0.	0.1000E 01

$U(DC/DT) = 3.6091$ $U(VAR) = 3.6090$ $E = 0.5664E 00$

10 = 0.2546E 03 11 = 0.3938E 04 12 = -0.3146E 03
 13 = 0.4223E 03 L = 0.2283E-02 K = 0.8128E 00
 AR = 0.1407E-03 ARE = 0.1408E-03

DEPTH	UR	UZ	TR	DCDA	DCDB	DCDR
H						
0						
1						
2						
3						
4						
5						
6						
7						
8						
9						
10						
11						
12						
13						
14						
15						
16						
17						
18						
19						
20						
21						
22						
23						
24						
25						
26						
27						
28						
29						
30						
31						
40						
100						

0. 0.5664E 00 0.1000E 01 -0.2713E-03 -0.1792E-03 0. 0.
 98 0. 0.1648E 12 -0.1183E 12 0.2213E 13 -0.3907E 13 0.
 99 0. 0.4633E 12 -0.3326E 12 0.6222E 13 -0.1099E 14 0.
 0. -0.3729E 11 0. 0.1000E 01 0.

III. LEIGEN

PROGRAMMER: R.B. HERRMANN / DEC 1971

PURPOSE:

This program accepts the card output of the program SURFACE. This output consists of the period, phase velocity pairs of the Love wave dispersion in plane layered solid media. The program LEIGEN determines the Love wave eigenfunctions, partial of phase velocity with respect to shear velocity and density as a function of depth within the medium. One liquid layer may be placed at the surface.

INPUT/OUTPUT

Input is from card. The eigenfunctions and phase velocity partial derivatives are listed on the printer. The eigenfunctions and phase velocity partial derivatives may also be listed on FILE 01. This file can then be used by other programs to generate synthetic surface wave seismograms and amplitude spectra.

THEORY

The differential equation to be solved is

$$\frac{d}{dz} \left\{ \mu \frac{dV}{dz} \right\} + (\rho \omega^2 - k^2 \mu) V = 0$$

with the boundary conditions $V(0) = 0$, $\frac{dV}{dz} = 0$ at $z = 0$.

$k = 2\pi/(CT)$ where C is the phase velocity and T is the period, ρ is the density and μ is the Lame constant $\mu = \rho \beta^2$, where β is the shear velocity in the layer. The other restriction on the solution is that $V(z) \rightarrow 0$ as $z \rightarrow \infty$.

The equation of motion can be derived from the following expression for the Lagrangian and variational methods. The Lagrangian of the system is

$$K = \omega^2 I_0 - k^2 I_1 - I_2,$$

where

$$I_0 = \int_0^\infty \rho V^2 dz, I_1 = \int_0^\infty \mu V^2 dz, \text{ and } I_2 = \int_0^\infty \mu \frac{dV}{dz}^2 dz.$$

The group velocity U can be defined in terms of the energy integrals as

$$U = \frac{kI_1}{\omega I_0}$$

The amplitude factor of Harkrider is defined as

$$A_L = \frac{1}{2 C U I_0} .$$

For a system of constant velocity layers, the partial of the phase velocity with respect to changes in the shear velocity or density in the m 'th layer are determined from the following relations:

$$\left(\frac{\partial c}{\partial \beta_m} \right)_{\rho, \omega} = 2 \rho_m \beta_m \left(\frac{\partial c}{\partial \mu_m} \right)_{\rho, \omega}$$

$$\left(\frac{\partial c}{\partial \rho_m} \right)_{\beta, \omega} = \left(\frac{\partial c}{\partial \rho_m} \right)_{\mu, \omega} + \beta_m^2 \left(\frac{\partial c}{\partial \mu_m} \right)_{\rho, \omega}$$

where

$$\left(\frac{\partial c}{\partial \rho_m} \right)_{\mu, \omega} = - \frac{c^3}{2I_1} \int_{z_{m-1}}^{z_m} v^2(z) dz$$

$$\left(\frac{\partial c}{\partial \mu_m} \right)_{\rho, \omega} = \frac{c}{2k^2 I_1} \int_{z_{m-1}}^{z_m} \{ k^2 v^2(z) + (dv/dz)^2 \} dz$$

where z_{m-1} and z_m are the lower and upper depth limits of the m 'th layer (note that z increases downward).

The numerical integration performed in LEIGEN uses the Haskell formulation of the Love wave problem. The integration is conducted from the halfspace to the free surface.

Several parameters are calculated which aid in judging the numerical error of the integration process. From the program SURFACE, the following parameters are available T , C , U , and A_L . The group velocity U is determined by implicit differentiation in SURFACE. Experience has shown that the U and A_L parameters determined by SURFACE are very reliable.

LEIGEN lists several parameters which can be used as goodness of fit parameters. The group velocity $U = U(ENER)$, which has been determined by using the energy integrals, can be compared to $U(DC/DT)$, which was determined by SURFACE. The $A_L = ALE$, the amplitude factor determined using the energy integrals, can be compared to the $A_L = AL$, which was determined by SURFACE using differentiation of the period equation. The Lagrangian = L should be zero if there are no numerical errors; this has been found to be a poor guide to goodness of fit. At the bottom of each printer listing, there is a line headed by $M = 0$. This contains the values of $UT (V)$ and $TT (\mu dV/dz)$ as computed at the free surface. If there is no numerical error, $TT = 0$.

The partial derivatives for the last row in the listing of the quantities as a function of depth are for changes in the parameters of the halfspace.

PROGRAM DESCRIPTION

PROGRAM LEIGEN : This program performs the desired numerical integration and lists the results on the printer.

SUBROUTINE TAPEIN : This subroutine makes it possible to add more eigenfunctions at different periods or for different models on the same tape. This allows expansion of the tape.

INPUT DATA

There are two slightly different arrangements to the data for the cases when the output is and is not written on FILE 01.

Case 1 - No output written on FILE 01.

Card

<u>Sequence</u>	<u>Column</u>	<u>Name</u>	<u>Format</u>	<u>Explanation</u>
A.	1-4	MMAX	I4	LE 0 END program
				GT 0 Number of layers in- cluding half- space in the model.

Card
Sequence

A.	Column	Name	Format	Explanation
	1-4	MMAX	I4	LE 0 End program
				GT 0 Number of layers including halfspace in the model
	11-14	ITAPE	I4	LE 0 No tape output (In this case must be LE 0)

B. Earth model (MMAX cards)

1-10	D(I)	F10.5	Layer thickness in km.
11-20	A(I)	F10.5	P velocity of layer in km/sec
21-30	B(I)	F10.5	S velocity of layer in km/sec
31-40	RHO(I)	F10.5	Density of layer in gm/cm ³ .

C. Period data

1-14	T	E14.7	LT 0 read new MMAX, ITAPE card (Point A)
			EQ 0 does nothing in this case
			GT 0 Period in seconds
16-29	C	E14.7	Phase velocity in km/sec
31-44	U	E14.7	Group velocity in km/sec
46-59	AL	E14.7	Amplitude factor

At this point the program will determine the appropriate eigenfunctions and return to Point C to read a new T, C, U, AL data card.

The program is terminated by reading a new MMAX card with MMAX.LE. 0. The last system card is an end of file.

Case 2 - Output written on FILE 01

Card	Sequence	Column	Name	Format	Explanation
A.		1-4	MMAX	I4	LE 0 End program
					GT 0 Number of layers including halfspace contained in the model
	11-14		ITAPE	I4	LE 0 No tape output
					EQ 1 Eigenfunctions written on FILE 01 but velocity-density model is not written again
					GT 1 Velocity-density model written on FILE as header as well as eigenfunctions

B. Earth model (MMAX cards)

1-10	D(I)	F10.5	Layer thickness in km
11-20	A(I)	F10.5	P velocity of layer in km/sec
21-30	B(I)	F10.5	S velocity of layer in km/sec
31-40	RHO(I)	F10.5	Density of layer in gm/cm ³

Card	Sequence	Column	Name	Format	Explanation
C. Tape update					
	1-10	IMAX		F10.5	Skip model already on tape and read through the eigenfunctions with T = TMAX. Tape now ready to write after T = TMAX.
	11-15	KMODMX	I5		Number of earth models to skip. KMODMX = 0 if tape is new and has nothing written on it.

Card Sequence	Column	Name	Format	Explanation
------------------	--------	------	--------	-------------

D. Period sequence

a. Header

1-4	IFUNC	I4	LT 0 Signifies end of eigenfunction written on FILE. This is essential when reading the FILE.
			EQ 1 Otherwise, IFUNC=1. This identifies the FILE as containing Love wave eigenfunctions. See program REIGEN where IFUNC=2 indicates Rayleigh wave eigenfunctions.
11-14	MODE	I4	Number of modes to be placed on tape for given period.
22-25	AMODE(I)	A4	Identification code for mode
26-30			
.			
.			
66-70			

b. Period data (MODE cards)

1-14	T	E14.7	Period in seconds
16-29	C	E14.7	Phase velocity in km/sec
31-34	U	E14.7	Group velocity in km/sec
46-59	AL	E14.7	Amplitude factor

c. End period data

1-80 blank card

At this point, program returns to Point D to read in a new IFUNC, MODE card. When the program finds an IFUNC LT 0, the program terminates.

Note, when the FILE created here is to be used by other programs such as WIGGLE, SRFWVPLT, QUESTION, or RADPAT, the order of data is as in the example, i.e., short periods first; within each period set, the modes are listed in increasing order, FUND first, 1'ST second, etc.

C THIS PROGRAM COMPUTES VALUES OF DISPLACEMENTS, STRESS AT THE
C MIDPOINT OF EACH LAYER TOGETHER WITH THE VALUE OF DC/DB FOR THAT
C LAYER. THE METHOD IS BASED ON THE HASKELL MATRIX WITH SPECIAL
C CONSIDERATION GIVEN TO THE HIGH FREQUENCIES. THE PROGRAM ALSO
C GIVES THE VARIOUS VALUES OF THE INTEGRALS MAKING UP THE
C LAGRANGIAN.
C MMAX = MINUS OR ZERO = END PROGRAM, = PLUS = NUMBER OF LAYERS
C D = LAYER THICKNESS, A = P VEL, B = S VEL, RHO = DENSITY
C DEPTH = DEPTH TO MIDPOINT OF LAYER XMU, XLAMB = LAME'S CONSTANTS
C T = PERIOD C = PHASE VELOCITY U = GROUP VELOCITY (DC/DT DERIVED)
C THE PROGRAM READS IN ONE SET OF T,C,U,AL OBSERVATIONS AT A TIME
C T = NEG READ IN NEW MODEL
C T = ZERO CHECK TO SEE WHETHER TO ENTER NEW PARAMETERS FOR
C TAPE WRITE

C A ONE LAYER OCEANIC MODEL WILL BE ACCEPTED, IN WHICH CASE THE
C LIQUID LAYER IS IGNORED

```

DIMENSION D(100),A(100),B(100),DEPTH(100),AMP(100),STRESS(100)
DIMENSION DCDR(100),DCDR(100),DMM(5),SMM(5),RHO(100)
DIMENSION APODE(10)
DIMENSION XMU(100),XLAMB(100)
XXMIN=1.0E-20
5 FORMAT(4(E14.7,1X))
11 FORMAT(2(I4,6X),10(1X,A4))
20 FORMAT(1H1)
21 FORMAT(1H0,13X,2HM ,8X,7HDEPTH ,10X,1HD,12X,5HALPHA,11X,4HBETA,
1 11X,3HRHO,13X,2HMU,11X,6HLAMBDA /1H )
23 FORMAT(1H ,10X,15,7(5X,F10.4))
25 FORMAT(1H ,10X,15,30X,5(5X,F10.4))
44 FORMAT(4F10.5)
5000 FORMAT(1H1,16X,1(10H T(SEC) = ,F7.2,13H C(KM/SEC) = ,F7.5,1X))
5001 FORMAT(1H0,16X,1(12H U(DC/DT) = ,F7.4,11H U(ENER) = ,F7.4,2X))
5002 FORMAT(1H0,16X,1(8H IC = ,E11.4,7H I1 = ,E11.4,2X))
5003 FORMAT(1H ,16X,1(8H I2 = ,E11.4,7H L = ,E11.4,2X))
5004 FORMAT(1H0,3H M,8X,5HDEPTH,1(6X,4HDISP,8X,6HSTRESS,8X,5HDC/DB,8X,
15HDC/DR )/1H0)
5005 FORMAT(1H ,13, F13.2,4(2X,E11.4))
5006 FORMAT(1H ,16X,1(8H WVNC = ,E11.4,20X)/1H )
5007 FORMAT(1H ,16X,1(8H AL = ,E11.4,7H ALE = ,E11.4,2X))
6010 FORMAT(1H ,A4,3F10.5,5E11.4)
6011 FORMAT(1H ,7(4X,E11.4))
6012 FORMAT(1H ,2(I4,6X))
6013 FORMAT(1H ,I4,5X,7F10.4)
12 FORMAT(I4,6X,I4,6X,F10.4)
801 READ(60,12) MMAX,ITAPE,ANY
NMAX = MMAX
IF(MMAX)777,777,778
778 WRITE(61,20)
WRITE(61,21)

```

```

BASE=0.0
DO 22 I=1,MMAX
READ(60,44) D(I),A(I),B(I),RHO(I)
BASE=BASE+D(I)
DEPTH(I)=BASE - D(I)*0.5
XMU(I) = RHO(I)*B(I)*B(I)
XLAMB(I) = RHO(I)*(A(I)*A(I)-2.*B(I)*B(I))
IF(I-MMAX)26,24,24
26 WRITE(61,23)I,DEPTH(I),D(I),A(I),B(I),RHO(I),XMU(I),XLAMB(I)
22 CONTINUE
24 WRITE(61,25)I,A(I),B(I),RHO(I),XMU(I),XLAMB(I)
DEPTH(MMAX) = BASE - D(MMAX)
IF(ITAPE.GT.0) CALL TAPEIN
IF(ITAPE.LE.1) GO TO 504
LMAX = MMAX + 1
IF(B(1).EQ.0.0) LMAX = LMAX - 1
WRITE(01,6012) LMAX
DPTH = 0.0
IF(B(1).EQ.0.0) DPTH = D(1)
DD = 0.0
K = 1
IF(B(1).EQ.0.0) K = 2
I = K - 1
WRITE(01,6013) I,DPTH,DD,A(K),B(K),RHO(K),XMU(K),XLAMB(K)
DO 200 I=1,MMAX
IF(B(I).LE.0.0) GO TO 200
WRITE(01,6013) I,DEPTH(I),D(I),A(I),B(I),RHO(I),XMU(I),XLAMB(I)
200 CONTINUE
504 IF(ITAPE.LE.0) GO TO 503
READ(60,11) IFUNC,MCODE,(AMODE(I),I=1,10)
WRITE(01,6012) IFUNC,MODE
IF(IFUNC.LT.0) GO TO 777
IIK = 0
503 CONTINUE
DO 500 J = 1,100
AMP(J) = 0.0
STRESS(J) = 0.0
DCDR(J) = 0.0
500 DCDB(J) = 0.0
READ(60,5) T,C,U,AL
IF(T) 801,504,499
499 CONTINUE
MMAX = NMAX
C
C      LAYER DROPPING PROCEDURE--HALFSPACE DEFINED AS ANY WAVELENGTHS
C      BELOW REGION WHERE C IS LESS THAN B
C      UNLESS OTHERWISE SPECIFIED, ANY=7.0
C
IF (ANY.LE.0) ANY=7.0
DMAX=ANY*C*T

```

```

899 SUM = 0
DO 900 II=1,MMAX
MAX = II
IF(C - B(II)) 901,900,900
901 SUM = SUM + D(II)
IF(SUM - DMAX) 900,900,902
900 CONTINUE
902 MMAX = MAX
WVNO = 5.2831853/(C*T)
WVNUM = WVNO
WVNOSQ=WVNO*WVNO
OMEGSQ =(6.2831853/T)**2
COVB = C/B(MMAX)
H = RHO(MMAX)*B(MMAX)*B(MMAX)
RB = WVNO*SQRT(ABS(COVB**2 - 1.))
UT = 1
TT = - H * RB
AMP(MMAX) = UT
STRESS(MMAX) = TT
IF(RB)1001,1000,1001
1000 DM = 1.0E25
SM = 0
GO TO 1002
1001 DM = 0.5/RB
SM = 0.5*RP
1002 CONTINUE
DLDR = OMEGSQ * DM
DLDM = -(WVNOSQ*DM + SM)
DCDE(MMAX) = 2.*RHO(MMAX)*B(MMAX)*C*DLDM/WVNO
DCDR(MMAX) = (C/WVNO)*(DLDR + B(MMAX)*B(MMAX)*DLDM)
MMI1 = MMAX - 1
SUMI0 = RHO(MMAX) * DM
SUMI1 = H * DM
SUMI2 = H * SM
DO 1340 K=1,MMI1
M = MMAX - K
IF(B(M).EQ.0.0) GO TO 1340
COVB = C/B(M)
RB = WVNO*SQRT(ABS(COVB**2 - 1.))
H=RHO(M)*B(M)*B(M)
DZ = D(M)/4.
DMM(1) = UT * UT
SMM(1) = (TT/H)**2
DO 1339 KK=2,5
XKK=KK-1
Q=RB*DZ*XKK
IF(C-B(M))1207,1221,1231
1231 SINQ = SIN(Q)
Y=SINQ/RB
Z=-RB*SINQ

```

```

COSQ = COS(Q)
GO TO 1242
1221 Y=DZ*XKK
Z=0.0
COSQ=1.0
GO TO 1242
1207 EXQP = EXP(Q)
EXQM=1./EXP
Y=(EXQP-EXQM)/(2.*RB)
Z=RB*RB*Y
COSQ=(EXQP+EXQM)/2.
1242 EUT = COSQ*LT - Y*TT/H
ETT = -H*Z*LT + COSQ*TT
DMM(KK) = EUT*EUT
SMM(KK) = (ETT*ETT)/(H*H)
IF(KK=3)1339,1301,1339
1301 AMP(M) = EUT
STRESS(M) = ETT
1339 CONTINUE
UT = EUT
TT = ETT
DM=(DZ/22.5)*(7.*(DMM(1)+DMM(5))+32.*(DMM(2)+DMM(4))+12.*DMM(3))
SM=(DZ/22.5)*(7.*(SMM(1)+SMM(5))+32.*(SMM(2)+SMM(4))+12.*SMM(3))
DLDM = -(WVNO*SQ*DM + SM)
DLDR = OMEGSO * DM
DCDB(M) = 2.*RHO(M)*B(M)*C*DLDM/WVNO
DCDR(M) = (C/WVNO)*(DLDR+B(M)*B(M)*DLDM)
SUMI0 = SUMI0 + RHO(M)*DM
SUMI1 = SUMI1 + H*DM
SUMI2 = SUMI2 + H*SM
1340 CONTINUE
AMP(100) = 1.0
STRESS(100) = TT/UT
DCDR(100) = 0.0
DCDR(100) = 0.0
DEPTH(100) = 0.0
IF(B(1),EQ,0.0) DEPTH(100) = D(1)
DLDK = -2.*WVNO*SUMI1
DO 1500 L=1,MMAX
IF(B(L),EQ,0.0) GO TO 1500
AMP(L) = AMP(L)/UT
STRESS(L) = STRESS(L)/UT
DCDB(L) = DCDB(L)/DLDK
DCDR(L) = DCDR(L)/DLDK
C EXCLUSION OF LOW AMPLITUDES FROM FINAL OUTPUT
IF(B(L) = B(MMAX))1506,1507,1507
1507 CONTINUE
IF(ABS(AMP(L))-XXMIN) 1501,1502,1502
1501 AMP(L) = 0.0
STRESS(L) = 0.0

```

```

DCDB(L) = 0.0
DCDR(L) = 0.0
GO TO 1500
1502 CONTINUE
1506 CONTINUE
1500 CONTINUE
  SUMI0 = SUMI0/(UT*UT)
  SUMI1 = SUMI1/(UT*UT)
  SUMI2 = SUMI2 /(UT*UT)
  UGR = SUMI1/(C*SUMI0)
  FLAGR=0MEGSQ*SUMI0-KVNO*SQ*SUMI1-SUMI2
  ALE = 1. / (2.*C*UGR*SUMI0)
501 CONTINUE
  WRITE(61,5000) T,C
  WRITE(61,5001) U,UGR
  WRITE(61,5002) SUMI0,SUMI1
  WRITE(61,5003) SUMI2,FLAGR
  WRITE(61,5007) AL,ALE
  WRITE(61,5006) WNUM
  WRITE(61,5004)
  M = 0
  DPH = DEPTH(100)
  UT = 1.0
  TT = 0.0
  DCB = 0.0
  DCR = 0.0
  WRITE(61,5005) M,DPH,UT,TT,DCB,DCR
  N = 1
  IF(B(1).EQ.0.0) N = 2
  DO 5010 M = N,MMAX
5010 WRITE(61,5005)M,DEPTH(M),AMP(M),STRESS(M),DCDB(M),DCDR(M)
  WRITE(61,45)
45 FORMAT(1H0)
  M = 0
  WRITE(61,5005) M,DEPTH(100),AMP(100),STRESS(100),DCDB(100),DCDR(10
10)
  IF(ITAPE.LE.0) GO TO 503
  MMAX = NMAX
  DUMM1=0.0
  DUMM2=0.0
  DUMM3=0.0
  K = IIK + 1
  WRITE(01,6010) AMODE(K),T,C,U,SUMI0,WNUM,AL,DUMM1,DUMM2
  WRITE(01,6011)UT,TT,DUMM1,DUMM2,DUMM3,DCB,DCR
  DO 6002 M = N,MMAX
  WRITE(01,6011)AMP(M),STRESS(M),DUMM1,DUMM2,DUMM3,DCDB(M),DCDR(M)
6002 CONTINUE
  IIK = IIK + 1
  GO TO 503
777 CONTINUE

```

END

```
SUBROUTINE TAPEIN
READ(60,1) TMAX,KMODEMX
1 FORMAT(F10.5,15)
KMODEL = 0
3002 CONTINUE
IF(KMODEL.EQ.KMODEMX) RETURN
3001 READ(01,6012) MMAX
6012 FORMAT(1H ,2(I4,6X))
DO 200 I = 1,MMAX
200 READ(01,6013) IDUM
6013 FORMAT(1H ,I4,5X,7F10.4)
2000 READ(01,6012) JFUNC,MODE
IF(JFUNC.LT.0) GO TO 3000
DO 2999 IM = 1,MODE
READ(01,6010) IDUM,T
6010 FORMAT(1H ,A4,3F10.5,5E11.4)
6011 FORMAT(1H ,7(4X,E11.4))
DO 2100 IN = 1,MMAX
2100 READ(01,6011) XDUM
2999 CONTINUE
IF(T.EQ.TMAX) RETURN
GO TO 2000
3000 KMODEL = KMODEL + 1
GO TO 3002
END
```

DATA PAGE 1

31	2		
1.0	5.00	2.89	2.5
1.0	6.1	3.52	2.7
1.0	6.1	3.52	2.7
1.0	6.1	3.52	2.7
1.0	6.1	3.52	2.7
1.0	6.1	3.52	2.7
1.0	6.1	3.52	2.7
1.0	6.1	3.52	2.7
1.0	6.1	3.52	2.7
1.0	6.1	3.52	2.7
1.0	6.1	3.52	2.7
1.0	6.1	3.52	2.7
1.0	6.1	3.52	2.7
1.0	6.1	3.52	2.7
1.0	6.1	3.52	2.7
1.0	6.1	3.52	2.7
1.0	6.1	3.52	2.7
1.0	6.4	3.7	2.9
1.0	6.4	3.7	2.9
1.0	6.4	3.7	2.9
1.0	6.4	3.7	2.9
1.0	6.4	3.7	2.9
1.0	6.4	3.7	2.9
1.0	6.4	3.7	2.9
1.0	6.4	3.7	2.9
1.0	6.4	3.7	2.9
1.0	6.4	3.7	2.9
2.0	6.7	3.87	3.0
2.0	6.7	3.87	3.0
2.0	6.7	3.87	3.0
2.0	6.7	3.87	3.0
2.0	6.7	3.87	3.0
2.0	6.7	3.87	3.0
2.0	6.7	3.87	3.0
2.0	6.7	3.87	3.0
2.0	6.7	3.87	3.0
2.0	6.7	3.87	3.0
2.0	6.7	3.87	3.0
2.0	6.7	3.87	3.0
	6.15	4.7	3.4

1	3	FUND 1 ST 2 ND	
2.000000E 00	3.4074768E 00	3.2084552E 00	7,3882933E-03
2.000000E 00	3.7060527E 00	3.4703516E 00	9,9816568E-04
2.000000E 00	3.8574009E 00	3.5638103E 00	5,3670618E-04

1	3	FUND 1 ST 2 ND	
2.500000E 00	3.4501979E 00	3.2832408E 00	5,1014991E-03
2.500000E 00	3.7627710E 00	3.5214867E 00	7,9472811E-04
2.500000E 00	3.9238278E 00	3.6557350E 00	3,6684530E-04

1	3	FUND 1 ST 2 ND	
3.000000E 00	3.4807968E 00	3.3257429E 00	3,8951203E-03
3.000000E 00	3.8108110E 00	3.5587629E 00	6,4451305E-04
3.000000E 00	3.9836080E 00	3.6276105E 00	3,4717436E-04

1	3	FUND 1 ST 2 ND	
3.500000E 00	3.5055192E 00	3.3529931E 00	3,1705755E-03

DATA PAGE 2

3.500000E 00	3.8535174E 00	3.5874673E 00	5,4247159E-04
3.500000E 00	4.0572187E 00	3.5507559E 00	4,0134764E-04

1	3	FUND 1 ST 2 ND	
4.000000E 00	3.5268961E 00	3.3727236E 00	2,6871940E-03
4.000000E 00	3.8933347E 00	3.6028673E 00	4,7919983E-04
4.000000E 00	4.1492589E 00	3.4714303E 00	4,5483497E-04

1	3	FUND 1 ST 2 ND	
5.000000E 00	3.5637321E 00	3.4015544E 00	2,0759147E-03
5.000000E 00	3.9736222E 00	3.5934040E 00	4,3275715E-04
5.000000E 00	4.3836053E 00	3.3874595E 00	4,8021588E-04

1	3	FUND 1 ST 2 ND	
6.000000E 00	3.5956625E 00	3.4234237E 00	1,6996357E-03
6.000000E 00	4.0638349E 00	3.5553564E 00	4,3267716E-04
6.000000E 00	4.6303347E 00	3.6631601E 00	3,2393162E-04

1	2	FUND 1 ST 2 ND	
7.000000E 00	3.6244602E 00	3.4409496E 00	1,4433825E-03
7.000000E 00	4.1675513E 00	3.5162360E 00	4,3739184E-04

1	2	FUND 1 ST 2 ND	
8.000000E 00	3.6511838E 00	3.4547122E 00	1,2588173E-03
6.000000E 00	4.2836896E 00	3.5001658E 00	4,2720734E-04

1	2	FUND 1 ST 2 ND	
9.000000E 00	3.6765957E 00	3.4650219E 00	1,1209179E-03
9.000000E 00	4.4041342E 00	3.5340773E 00	3,9193020E-04

1	2	FUND 1 ST 2 ND	
1.000000E 01	3.7012129E 00	3.4724826E 00	1,0147778E-03
1.000000E 01	4.5195807E 00	3.6479562E 00	3,2751996E-04

1	2	FUND 1 ST 2 ND	
1.100000E 01	3.7254500E 00	3.4776014E 00	9,3083309E-04
1.100000E 01	4.6152358E 00	3.8675161E 00	2,3690678E-04

1	2	FUND 1 ST 2 ND	
1.200000E 01	3.7495535E 00	3.4810791E 00	8,6269161E-04
1.200000E 01	4.6781983E 00	4.2090951E 00	1,2724813E-04

1	1	FUND 1 ST 2 ND	
1.300000E 01	3.7736849E 00	3.4834946E 00	8,0597303E-04

1	1	FUND 1 ST 2 ND	
1.400000E 01	3.7979346E 00	3.4853888E 00	7,5768043E-04

1	1	FUND 1 ST 2 ND	
1.500000E 01	3.8223424E 00	3.4872092E 00	7,1565338E-04

1	1	FUND 1 ST 2 ND	
1.6000000E 01	3.8469107E 00	3.4893358E 00	6,7838894E-04
1	1	FUND 1 ST 2 ND	
1.7000000E 01	3.8716151E 00	3.4919915E 00	6,4476916E-04
1	1	FUND 1 ST 2 ND	
1.8000000E 01	3.8964126E 00	3.4955367E 00	6,1398530E-04
1	1	FUND 1 ST 2 ND	
1.9000000E 01	3.9212464E 00	3.5001182E 00	5,8547207E-04
1	1	FUND 1 ST 2 ND	
2.0000000E 01	3.9460508E 00	3.5059884E 00	5,5878564E-04
1	1	FUND 1 ST 2 ND	
2.2000000E 01	3.9952857E 00	3.5217619E 00	5,0968754E-04
1	1	FUND 1 ST 2 ND	
2.4000000E 01	4.0435336E 00	3.5433432E 00	4,6502086E-04
1	1	FUND 1 ST 2 ND	
2.6000000E 01	4.0902402E 00	3.5708577E 00	4,2393653E-04
1	1	FUND 1 ST 2 ND	
2.8000000E 01	4.1349285E 00	3.6036696E 00	3,8603595E-04
1	1	FUND 1 ST 2 ND	
3.0000000E 01	4.1772286E 00	3.6409709E 00	3,5111261E-04
1	1	FUND 1 ST 2 ND	
3.5000000E 01	4.2711602E 00	3.7480045E 00	2,7625661E-04
1	1	FUND 1 ST 2 ND	
4.0000000E 01	4.3478094E 00	3.8616224E 00	2,1787531E-04
1	1	FUND 1 ST 2 ND	
4.5000000E 01	4.4088597E 00	3.9699129E 00	1,7342323E-04
1	1	FUND 1 ST 2 ND	
5.0000000E 01	4.4570596E 00	4.0673087E 00	1,3987656E-04
1	1	FUND 1 ST 2 ND	
6.0000000E 01	4.5254651E 00	4.2225180E 00	9,5084958E-05
1	1	FUND 1 ST 2 ND	
7.0000000E 01	4.5695616E 00	4.3327979E 00	6,8241050E-05
1	1	FUND 1 ST 2 ND	

DATA PAGE 4

8.000000E 01 4.5991852E 00 4.4111944E 00 5.1209016E-05

1 1 FUND 1 ST 2 ND
9.000000E 01 4.6198920E 00 4.4680420E 00 3.9813745E-05

1 1 FUND 1 ST 2 ND
1.000000E 02 4.6348776E 00 4.5099175E 00 3.1839386E-05

1 1 FUND 1 ST 2 ND
1.500000E 02 4.6708521E 00 4.6136025E 00 1.3670287E-05

1 1 FUND 1 ST 2 ND
2.000000E 02 4.6835724E 00 4.6510571E 00 7.5841452E-06

-1 1 FUND

$$T(\text{SEC}) = 2.00 C(\text{KM/SEC}) = 3,40748$$

$U(DC/DT) = 3.2085$ $U(E\delta ER) = 3.2084$

```

I0 = 0.6190E 01 I1 = 0.6767E 02
I2 = 0.3568E 01 L = -0.4387E-05
AL = 0.7388E-02 ALE = 0.7389E-02
WVNO = 0.9220E 00

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LAMBDA

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BETA

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2

X
I
I
I

T(SEC) = 2.00 C(KM/SEC) = 3.70605

U(DC/DT) = 3.4704 U(ENER) = 3.4703

I0 = 0.3894E 02 I1 = 0.5008E 03
 I2 = 0.2445E 02 L = -0.3946E-04
 AL = 0.9982E-03 ALE = 0.9984E-03
 WVNO = 0.8477E 00

M	DEPTH	DISP	STRESS	DG/DB	DC/DR
0	0.	0.1000E 01	0.	0.	0.
1	0.50	0.9427E 00	-0.4742E 01	0.5079E-01	-0.1431E-01
2	1.50	0.6365E 00	-0.9864E 01	0.3729E-01	0.3486E-02
3	2.50	0.3208E 00	-0.1112E 02	0.1859E-01	0.6454E-02
4	3.50	-0.1974E-01	-0.1152E 02	0.1224E-01	0.7460E-02
5	4.50	-0.3587E 00	-0.1102E 02	0.2018E-01	0.6200E-02
6	5.50	-0.6700E 00	-0.9669E 01	0.4080E-01	0.3056E-02
7	6.50	-0.9293E 00	-0.7569E 01	0.6567E-01	-0.1017E-02
8	7.50	-0.1117E 01	-0.4883E 01	0.8939E-01	-0.4780E-02
9	8.50	-0.1217E 01	-0.1819E 01	0.1040E 00	-0.7092E-02
10	9.50	-0.1224E 01	0.1386E 01	0.1049E 00	-0.7248E-02
11	10.50	-0.1154E 01	0.3018E 01	0.1064E 00	0.1865E-03
12	11.50	-0.1076E 01	0.3122E 01	0.9271E-01	0.2437E-03
13	12.50	-0.9965E 00	0.3219E 01	0.7961E-01	0.2986E-03
14	13.50	-0.9143E 00	0.3308E 01	0.6718E-01	0.3507E-03
15	14.50	-0.8299E 00	0.3390E 01	0.5554E-01	0.3995E-03
16	15.50	-0.7436E 00	0.3463E 01	0.4479E-01	0.4446E-03
17	16.50	-0.6555E 00	0.3529E 01	0.3504E-01	0.4854E-03
18	17.50	-0.5659E 00	0.3586E 01	0.2638E-01	0.5217E-03
19	18.50	-0.4750E 00	0.3634E 01	0.1889E-01	0.5531E-03
20	19.50	-0.3829E 00	0.3674E 01	0.1264E-01	0.5793E-03
21	21.00	-0.2636E 00	0.2892E 01	0.1345E-01	0.1329E-02
22	23.00	-0.1618E 00	0.1775E 01	0.5066E-02	0.5006E-03
23	25.00	-0.9926E-01	0.1089E 01	0.1907E-02	0.1885E-03
24	27.00	-0.6088E-01	0.6690E 00	0.7176E-03	0.7101E-04
25	29.00	-0.3730E-01	0.4113E 00	0.2696E-03	0.2675E-04
26	31.00	-0.2279E-01	0.2535E 00	0.1008E-03	0.1007E-04
27	33.00	-0.1383E-01	0.1573E 00	0.3726E-04	0.3795E-05
28	35.00	-0.8221E-02	0.9946E-01	0.1335E-04	0.1432E-05
29	37.00	-0.4615E-02	0.6575E-01	0.4387E-05	0.5473E-06
30	39.00	-0.2132E-02	0.4803E-01	0.1143E-05	0.2268E-06
31	40.00	-0.1116E-02	0.4371E-01	0.1948E-06	0.7389E-07
0	0.	0.1000E 01	-0.2034E-04	0.	0.

T(SEC) = 2.00 C(KM/SEC) = 3.85740

U(DC/DT) = 3.5638 U(ENER) = 3.5638

I0 = 0.6776E 02 I1 = 0.9315E 03
 I2 = 0.5091E 02 L = -0.2338E-04
 AL = 0.5367E-03 ALE = 0.5368E-03
 WENO = 0.8144E 00

M	DEPTH	DISP	STRESS	DG/DB	DC/DR
0	0.	0.1000E 01	0.	0.	0.
1	0.50	0.9359E 00	-0.5296E 01	0.2890E-01	-0.9306E-02
2	1.50	0.5920E 00	-0.1141E 02	0.2088E-01	0.2637E-02
3	2.50	0.2193E 00	-0.1324E 02	0.1158E-01	0.5685E-02
4	3.50	-0.1823E 00	-0.1333E 02	0.1113E-01	0.5835E-02
5	4.50	-0.5598E 00	-0.1165E 02	0.1974E-01	0.3010E-02
6	5.50	-0.8636E 00	-0.1844E 01	0.3304E-01	-0.1349E-02
7	6.50	-0.1054E 01	-0.4124E 01	0.4425E-01	-0.5020E-02
8	7.50	-0.1105E 01	0.7408E 00	0.4764E-01	-0.6133E-02
9	8.50	-0.1010E 01	0.1550E 01	0.4150E-01	-0.4119E-02
10	9.50	-0.7827E 00	0.9550E 01	0.2895E-01	-0.5789E-05
11	10.50	-0.4835E 00	0.1176E 02	0.1652E-01	0.3145E-02
12	11.50	-0.1762E 00	0.1252E 02	0.8371E-02	0.4134E-02
13	12.50	0.1412E 00	0.1256E 02	0.7925E-02	0.4188E-02
14	13.50	0.4505E 00	0.1188E 02	0.1528E-01	0.3295E-02
15	14.50	0.7340E 00	0.1052E 02	0.2877E-01	0.1657E-02
16	15.50	0.9754E 00	0.8554E 01	0.4535E-01	-0.3559E-03
17	16.50	0.1161E 01	0.6098E 01	0.6127E-01	-0.2288E-02
18	17.50	0.1280E 01	0.3292E 01	0.7292E-01	-0.3703E-02
19	18.50	0.1325E 01	0.2981E 00	0.7768E-01	-0.4281E-02
20	19.50	0.1294E 01	-0.2713E 01	0.7446E-01	-0.3890E-02
21	21.00	0.1161E 01	-0.3938E 01	0.1311E 00	0.1265E-02
22	23.00	0.9951E 00	-0.3521E 01	0.9645E-01	0.9761E-03
23	25.00	0.8465E 00	-0.3164E 01	0.6989E-01	0.7549E-03
24	27.00	0.7126E 00	-0.2863E 01	0.4962E-01	0.5860E-03
25	29.00	0.5910E 00	-0.2611E 01	0.3423E-01	0.4578E-03
26	31.00	0.4795E 00	-0.2403E 01	0.2265E-01	0.3613E-03
27	33.00	0.3764E 00	-0.2238E 01	0.1498E-01	0.2899E-03
28	35.00	0.2797E 00	-0.2111E 01	0.7926E-02	0.2386E-03
29	37.00	0.1879E 00	-0.2020E 01	0.3759E-02	0.2039E-03
30	39.00	0.9934E-01	-0.1965E 01	0.1289E-02	0.1833E-03
31	40.00	0.5579E-01	-0.1950E 01	0.2936E-03	0.9988E-04
0	0.	0.1000E 01	0.1078E+04	0.	0.

IV. SUBROUTINE EIGEN

PROGRAMMER : R. B. HERRMANN / MAY 1973

PURPOSE:

This subroutine reads the eigenfunctions from the tape on Unit ITAPE and stores the eigenfunctions for a set of up to 15 focal depths labeled COMMON. This subroutine requires the subroutines MODEL prior to use and it uses the subroutine NONPRP for linear interpolation to obtain the eigenfunctions at various focal depths.

INPUT/OUTPUT

The eigenfunction tape is loaded onto the tape Unit ITAPE. Input and output parameters are stored in COMMON.

```

SUBROUTINE SIGER(ITAPE,TT,T)
COMMON/ONE/AY(100,4),CY(100,4),N,DEPTH(100),D(100)
COMMON/TWO/LAYER(100),A(100),B(100),RHO(100),XMU(100),XLAMB(100)
COMMON/THREE/ANODE(10),WVNC(10),ELIP(10),ABSORP(10),DENOM(10)
COMMON/FORE/NUMDPT,AAA(15),PYY1(15,10),PYY2(15,10),PYY3(15,10),
LPYY4(15,10),IFUNC
DIMENSION PYR(4)
2000 READ(ITAPE,1) IFUNC,MODE
 1 FORMAT(1H ,2(14.6X))
MAXY = 2 * IFUNC
IF(IFUNC.LT.0) TT = - TT
IF(IFUNC.LT.0) GO TO 9999
DO 2999 IM = 1,100
READ(ITAPE,2) ANODE(IM),T,C,U,SUMI0,WVNC(IM),AL,ELIP(IM),
ABSORP(IM)
5 FORMAT(1H ,A4,3F10.5,5E+1.4)
C THE SPECTRA ARE COMPUTED AT A DISTANCE OF 1000 KILOMETERS
DENOM(IM) = SQRT(WVAC(IM))*0.3989423*1.0E+20
DENOM(IM) = DENOM(IM)/AL
DENOM(IM) = DENOM(IM) * 6.2831853 * 31.622776
DO 2100 IN=1,N
2100 READ(ITAPE,3)AY(IN,1),AY(IN,2),AY(IN,3),AY(IN,4),DCDA,DCDB,DCDR
 6 FORMAT(1H ,2(4X,E11.4))
IF(TT.NE.T) GO TO 2999
WVN02 = WVNC(IM)**2
GO TO (2400,2200),IFUNC
C RAYLEIGH WAVES
INITIALLY AY(IN,1) = UR, AY(IN,2) = UZ, AY(IN,3) = TZ,
AY(IN,4) = TR
NOW AY(IN,1) = UR, AY(IN,2) = UZ, AY(IN,3) = DURDZ,
AY(IN,4) = DUZDZ
AND CY(IN,1) = DURDZ, CY(IN,2) = DUZDZ, CY(IN,3) = D2URDZ2,
CY(IN,4) = D2UZDZ2.
2200 DO 2300 IN = 1,1
  H = RHO(IN)*C*U*WVN02
  XL2M = XLAMB(IN) + 2.*XMU(IN)
  CY(IN,1) = -WVNC(IM)*AY(IN,2) + AY(IN,4)/XMU(IN)
  CY(IN,2) = (WVNC(IM)*XLAMB(IN)*AY(IN,1)+AY(IN,3))/XL2M
  CY(IN,3) = (-H/XMU(IN) + WVN02*(3.+XL2M-2.*XMU(IN))/XL2M)*AY(IN,1)
1 -WVNC(IM)*(1.+XLAMB(IN)/XMU(IN))*AY(IN,3)/XL2M
  CY(IN,4) = (-(H+WVN02*XLAMB(IN))*AY(IN,2)+WVNC(IM)*(1.+XLAMB(IN)/
2 XMU(IN))*AY(IN,4))/XL2M
  AY(IN,3) = CY(IN,1)
  AY(IN,4) = CY(IN,2)
2300 CONTINUE
  GO TO 2501
C LOVE WAVES
INITIALLY AY(IN,1) = UT, AY(IN,2) = TT
NOW AY(IN,1) = UT, AY(IN,2) = DUTDZ

```

```
C      AND CY(IN,1) = LUTDZ, CY(IN,2) = D2UTDZ2
2400 DO 2500 IN=1,N
      CY(IN,1) = AY(IN,2)/XMU(IN)
      CY(IN,2) = .WVNC*(XMU(IN)- C*C*RHO(IN))*AY(IN,1)/XMU(IN)
      AY(IN,2) = CY(IN,1)
2500 CONTINUE
2501 CONTINUE
DO 2555 ID = 1,NUMDFT
AA = AAA(ID)
NPER = ID
IF(AA.LT.0) AA = 0.0
CALL NONPRP(AA,PAXY,PYR)
GO TO(2551,2552),IFLNC
2551 CONTINUE
PYY1(NPER,1)= .WVNC(IM)*PYR(1)/DENOM(IM)
PYY2(NPER,1)= PYR(2)/DENOM(IM)
GO TO 2555
2552 CONTINUE
PYY1(NPER,1)= .WVNC(IM)*PYR(1)/DENOM(IM)
PYY2(NPER,1)= .WVNC(IM)*PYR(2)/DENOM(IM)
PYY3(NPER,1)= PYR(3)/DENOM(IM)
PYY4(NPER,1)= PYR(4)/DENOM(IM)
2555 CONTINUE
2999 CONTINUE
IF(TT.NE.T1) GO TO 2000
9999 CONTINUE
RETURN
END
```


V. SUBROUTINE MODEL

PROGRAMMER : R. B. HERRMANN / MAY 1973

PURPOSE:

This subroutine reads the earth model which is stored on the eigenfunction tape immediately in front of the eigenfunctions. The earth model is placed in labeled COMMON.

INPUT/OUTPUT

The earth model is read off the eigenfunction tape which is stored on tape Unit ITAPE.

```
SUBROUTINE MODEL(ITAPE)
COMMON/ONE/AY(100,4),CY(100,4),N,DEPTH(100),D(100)
COMMON/TWO/LAYER(100),A(100),B(100),RHO(100),XMU(100),XLAMB(100)
READ(ITAPE,1) N
1 FORMAT(1H ,2(14.6X))
C READ IN MODEL
DO 100 I=1,N
100 READ(ITAPE,2) LAYER(I),DEPTH(I),D(I),A(I),B(I),RHO(I),XMU(I),
1 XLAMB(I)
2 FORMAT(1H ,14.5A,7F10.4)
RETURN
END
```

VI. SUBROUTINE NONPRP

PROGRAMMER : R. B. HERRMANN / MAY 1972

PURPOSE:

This subroutine is called by the subroutine EIGEN in order to obtain the eigenfunctions at a specified depth by linear interpolation of known eigenfunctions at depths above and below the desired value.

INPUT/OUTPUT

All values needed are in labeled COMMON or are in the subroutine argument.

NONPRP PAGE 1

```
SUBROUTINE NONPRP(A,MAXY,PYR)
COMMON/DNE/AY(100,4),CY(100,4),N,DEPTH(100),D(100)
DIMENSION PYR(4)
DO 100 I=1,N
  IFICA = I
  D1=DEPTH(I)-0.5*D(I)
  D2=DEPTH(I)+0.5*D(I)
  IF(A.GE.D1.AND.A.LE.D2) GO TO 101
100 CONTINUE
101 DO 110 JJ = 1,MAXY
    AA = AY(IFICA,JJ)
    CC = CY(IFICA*JJ)
    AA = AA - CC * DEPTH(IFICA)
110 PYR(JJ) = AA + CC*A
RETURN
END
```

VII. SUBROUTINE EXCIT

PROGRAMMER: R. B. HERRMANN / MAY 1972

PURPOSE:

This subroutine calculates the theoretical surface-wave amplitude spectra at a given period from the eigenfunctions of that period for a given set of source orientation factors. The source can be either a dislocation source specified by the strike, dip, and slip, or it can be an explosive source.

INPUT/OUTPUT

All information needed for the functioning of the subroutine is taken from the argument of the subroutine.

EXCIT PAGE 1

```
SUBROUTINE EXCIT(FAC1,FAC2,FAC3,FAC4,FAC5,FAC6,FAC7,FAC8,FAC9,
1 PY1,PY2,PY3,PY4,IFLNC,SAZ,CAZ,AMP,IEQEX)
GO TO (100,200),IFQEX
100 CONTINUE
S2AZ = 2.*SAZ*CAZ
C2AZ = CAZ*CAZ - SAZ*SAZ
GO TO (1,2),IFUNC
1 CONTINUE
XRL = PY1*((FAC1-FAC4)*C2AZ + (FAC2+FAC3)*S2AZ )
XIM = - PY2*((FAC5+FAC8)*CAZ + (FAC6-FAC7)*SAZ)
AMP = SQRT(XRL * XRL + XIM * XIM )
RETURN
2 CONTINUE
XRL = PY1*((FAC3+FAC2)*C2AZ + (FAC4-FAC1)*S2AZ) + FAC9 * PY4
1 + J*5*FAC9*PY1
XIM = (PY3+PY2)*((FAC7-FAC6)*CAZ+(FAC8+FAC5)*SAZ)
AMP = SQRT(XRL * XRL + XIM * XIM )
RETURN
200 CONTINUE
GO TO (10,20), IFUNC
10 AMP = 0
RETURN
20 AMP = PY4 - PY1
RETURN
END
```

VIII. SRFWVPLT

PROGRAMMER: R. B. HERRMANN / Jun 72

PURPOSE:

This program calculates the theoretical surface-wave amplitude spectra for each individual mode for a choice of two source spectrums of the dislocation source. The two source spectrums are $S(\omega) = 1$, which corresponds to the delta-response of the system, and $S(\omega) = 1/i\omega$, which corresponds to the response of the system to a step dislocation. Depending on the surface-wave eigenfunction tape used, the program will plot the Love or Rayleigh wave amplitude spectra. The theoretical amplitude spectra are generated at a reference distance of 1000 km or 9° from the source. To specify the double-couple source, the orientations of the pressure and tension axes are used. As an aid in identifying the plots, the focal mechanism specified by the pressure and tension axes is plotted. The spectra are plotted on a 3×3 log-log scale, of approximately 2.5 inches per cycle.

INPUT/OUTPUT

The specification of the focal depth, focal mechanism orientation, and plotting options is taken from card on FILE 60. Printer output is on FILE 61. The eigenfunctions are read from the tape on FILE 01 (this is the tape written by REIGEN or LEIGEN using exactly the same order of input as shown in the sample input for these programs). FILE 10 is used for the CALCOMP plotting.

THEORY

The theory of this program is taken from Levshin and Yanson (1971). The formulation in terms of pressure and tension axes is used.

PROGRAM DESCRIPTION

PROGRAM SRFWVPLT : This program reads in the data and controls the plotting.

SUBROUTINE NONPRP: This subroutine performs a linear interpolation on the eigenfunctions to arrive at the values for a specified depth.

SUBROUTINE ALOGAXES: This subroutine establishes the 3×3 cycle log-log scale for the plot.

SUBROUTINE LIND: This is a subroutine for plotting the amplitude spectra for the various modes. Any amplitude information which is less than three orders of magnitude of the largest amplitude plotted, is not plotted. This allows a neater plot by eliminating plotting below the y-axis or along the x-axis.

SUBROUTINE DIPDD: This is a subroutine used by the subroutine NODPL. It takes a vector and redefines the vector in terms of two angles.

SUBROUTINE NODPL; This subroutine plots a focal mechanism on an equal area projection for a visual representation of the focal mechanism specified by the given orientations of the pressure and tension axes.

Card Sequence	Column	Name	Format	Explanation
A.	1-10	SIZE	F10.5	CALCOMP reduction factor between 0.0 and 1.0 .
B.	1-10	DEL	F10.5	Azimuth of the T axis measured from North
	11-20	BET	F10.5	The angle the T axis makes with the downward vertical axis
	21-30	ALP	F10.5	Azumuth of the P axis measured clockwise from North
	31-40	GAM	F10.5	The angle the P axis makes with the downward vertical axis
C.	1-10	AA	F10.5	.LT. - 10 read in new mechanism parameters .GT. - 10 and .LT. 0 end program .GT. 0 focal depth in km

Card Sequence	Column	Name	Format	Explanation
D.	1-10	PHI	F10.5	.LT.O read new AA .GE.O receiver azimuth
	11-20	RAD	F10.5	.GT.O plot radial spectra of Rayleigh wave .LT.O no radial plot
	21-30	VER	F10.5	.GT.O plot vertical spectra of Rayleigh wave .LE.O no vertical plot
	31-40	SRC	F10.5	.GT.O spectra for 1/iw source-step response .LE.O spectra for 1.0 source-impulse response (when using the LOVE tape VER and RAD have no meaning)

After performing the operations specified by the PHI, VER, RAD, SRC card, the program will read another PHI, VER, RAD, SRC card.

The example given on the following page is for the case DEL = 45, BET = 90., ALP = 135, and GAM = 90. This corresponds to a left lateral vertical strike slip fault striking north. Since the plots given here are only of amplitude spectra and not the phase spectra, the plots will also correspond to a right lateral strike slip fault striking north. Basically, an interchange of the P and T axes will not affect the amplitude spectra.

C PROGRAM SRFWVPLT
C THIS PROGRAM ACCEPTS ANY ARBITRARY ORIENTATION FOR A DOUBLE COUPLE
C FOCAL MECHANISM. THE INPUT IS THE ORIENTATION OF THE PRESSURE AND
C TENSION AXES. THE PROGRAM PLOTS THEORETICAL LOVE OR RAYL WAVE
C SPECTRA AT A DISTANCE OF 9 DEGREES OR 1000 KILOMETERS ON A
C STANDARD 3 X 3 LOG-LOG SCALE. THERE IS AN OPTION FOR RADIAL AND
C OR VERTICAL DISPLACEMENT SPECTRUM FOR RAYLEIGH WAVES. USING THE
C EIGENFUNCTION TAPE FOR LOVE WAVES WILL GIVE THE LOVE WAVE SPECTRA.
COMMON/ONE/AY(100,4),CY(100,4),N,DEPTH(100),D(100),
1 LAYER(100),A(100),B(100),RHO(100),XMU(100),XLAMB(100),
1 IBUF(1026), ELIP
DIMENSION AMODE(10),WVNC(10),ELIP(60,10),ABSORP(10),DENOM(10)
DIMENSION LABEL(3),LABLE(4),XX(60),YY1(60,10),NPERM(10),YY2(60,10)
DIMENSION LABE(7),LAB(3)
DIMENSION X(62),Y(62),PYR(4),PYI(4)
DIMENSION PYI(60,10),PYY2(60,10),PYY3(60,10),PYY4(60,10)
DIMENSION TT(60)
CHARACTER LAB,LABE,LABEL,LABLE
CALL PLOTS(IBUF,1026,10)
CALL PLCT(0.0,-11.0,-3)
CALL PLCT(0.0,2.0,-3)
READ(60,4) SIZE
C SIZE = REDUCTION FACTOR USED IN CALL FACTOR OF CALCOMP LIBRARY
C IF SIZE = 1.0 PLOT IS 3 CYCLE LOG-LOG + 2.5 INCHES PER CYCLE
C HERE SIZE MUST BE LE. 1.0
IF(SIZE.LE.0.0,SIZE.GT.1.0) SIZE = 1.0
CALL FACTOP(SIZE)
LAB(1) = 6H U2
LAB(2) = 6H UR
LAB(3) = 6H UT
LABE(1) = 6HDELT
LABE(2) = 6H BETA
LABE(3) = 6H AL
LABE(4) = 6HPH
LABE(5) = 6HGAM
LABEL(1) = 6HPERICD
LABEL(2) = 6H (SEC)
LABEL(1) = 6H
LABEL(2) = 6HAMP (C
LABEL(3) = 6HM-SEC)
1 FORMAT(1H ,2(I4,6X))
2 FORMAT(1H ,14,5X,7F10.4)
4 FORMAT(8F10.5)
5 FORMAT(1H ,A4,3F10.5,5E11.4)
6 FORMAT(1H ,7(4X,E11.4))
DEGRAD = 0.017452329
4321 READ(60,4) DEL,BET,ALP,GAM
C DEL IS THE AZIMUTH OF THE T AXIS
C BET IS THE ANGLE THE T AXIS MAKES WITH THE DOWNWARD VERTICAL
C ALP IS THE AZIMUTH OF THE P AXIS

```

C      GAM IS THE ANGLE THE P AXIS MAKES WITH THE DOWNWARD VERTICAL
      WRITE(61,7) DEL,BET,ALP,GAM
7 FORMAT(1HD, 4F10.3)
      CALL SYMBOL(1.35,5.5,0.14,1HT,0.0,1)
      CALL SYMBOL(3.59,5.5,0.14,1HP,0.0,1)
      CALL SYMBOL(0.56,5.0,0.14,LABE,0.0,28)
      CALL NUMBER(0.56,4.7,0.14,DEL,0.0,-1)
      CALL NUMBER(1.68,4.7,0.14,BET,0.0,-1)
      CALL NUMBER(2.80,4.7,0.14,ALP,0.0,-1)
      CALL NUMBER(3.92,4.7,0.14,GAM,0.0,-1)
      DEL = DEGRAD * DEL
      BET = DEGRAD * BET
      ALP = DEGRAD * ALP
      GAM = DEGRAD * GAM
      FACT1 = (COS(BET)**2) - (COS(GAM)**2)
      FACT2 = (SIN(BET))**2
      FACT3 = (SIN(GAM))**2
      FACT4 = SIN(2.*BET)
      FACT5 = SIN(2.*GAM)
      CALL NODPL(DEL,BET,ALP,GAM)
      CALL PLOT(9.0*0.0,-3)
2001 READ(60,4) AA
      IF(AA.LE.-10.) GO TO 4321
      IF(AA.LT.0) GO TO 9998
      DO 37 I=1,50
      DO 37 J=1,10
      PYY1(I,J) = 0
      PYY2(I,J) = 0
      PYY3(I,J) = 0
37 PYY4(I,J) = 0
      READ(01,1)N
C      READ IN MODEL
      DO 100 I=1,N
100 READ(01,2) LAYER(I),DEPTH(I),D(I),A(I),B(I),RHO(I),XMU(I),XLAMB(I
1)
      NPER = 0
      DO 1313 I = 1,10
1313 NPERM(I) = 0
2000 READ(01,1) IFUNC, MODE
      MAXY = 2 * IFUNC
      IF(IFUNC.LT.0) GO TO 9999
      IF(NPER.EQ.0) MODEMX = MODE
      IF(NPER.EQ.0) IFUN = IFLNC
      DO 1314 IZ = 1,MODE
1314 NPERM(IZ) = NPERM(IZ) + 1
      NPER = NPER + 1
      DO 2999 IM = 1,MODE
      READ(01,5) AMODE(IM),T,C,U,SUM10,WVNO(IM),AL,ELIP(NPER,IM),ABSORP(
1IM)
      XX(NPER) = T

```

```

      TT(NPER) = T
C   THE SPECTRA ARE COMPUTED AT A DISTANCE OF 1000 KILOMETERS
C   DENOM(IM) = SQRT (WVNC(IM))*0.3989423*1.0E+20
C   DENOM(IM) = DENOM(IM)/AL
C   DENOM(IM) = DENOM(IM) * 5.2831853 * 31.622776
DO 2100 IN=1,N
2100 READ(01,6) AY(IN,1),AY(IN,2),AY(IN,3),AY(IN,4),DCDA,DCDB,DCDR
WVN02 = WVNC(IM)**2
GO TO (2400,2200),IFUNC
C   RAYLEIGH WAVES
C   INITIALLY AY(IN,1) = LR, AY(IN,2) = UZ, AY(IN,3) = TZ
C   AY(IN,4) = TR
C   NOW AY(IN,1) = UR, AY(IN,2) = UZ, AY(IN,3) = DURDZ,
C   AY(IN,4) = DUZDZ
C   AND CY(IN,1) = DURDZ, CY(IN,2) = DUZDZ, CY(IN,3) = D2URDZ2,
C   CY(IN,4) = D2UZDZ2.
2200 DO 2300 IN = 1,N
      H = RHO(IN)*C*C*WVN02
      XL2M = XLAMB(IN) + 2.*XMU(IN)
      CY(IN,1) = -WVNC(IM)*AY(IN,2) + AY(IN,4)/XMU(IN)
      CY(IN,2) = (WVNC(IM)*XLAMB(IN)*AY(IN,1)+AY(IN,3))/XL2M
      CY(IN,3) = (-H/XMU(IN) + WVN02*(3.*XL2M-2.*XMU(IN))/XL2M)*AY(IN,1)
      1 -WVNC(IM)*(1.+XLAMB(IN)/XMU(IN))*AY(IN,3)/XL2M
      CY(IN,4) = (-H+WVN02*XLAMB(IN))*AY(IN,2)+WVNC(IM)*(1.+XLAMB(IN)/
      1 XMU(IN))*AY(IN,4))/XL2M
      AY(IN,3) = CY(IN,1)
      AY(IN,4) = CY(IN,2)
2300 CONTINUE
GO TO 2501
C   LOVE WAVES
C   INITIALLY AY(IN,1) = LT, AY(IN,2) = TT
C   NOW AY(IN,1) = LT, AY(IN,2) = DUTDZ
C   AND CY(IN,1) = DUTOZ, CY(IN,2) = D2UTDZ2
2400 DO 2500 IN=1,N
      CY(IN,1) = AY(IN,2)/XMU(IN)
      CY(IN,2) = WVN02*(XMU(IN)- C*C*RHO(IN))*AY(IN,1)/XMU(IN)
      AY(IN,2) = CY(IN,1)
2500 CONTINUE
2501 CONTINUE
CALL NONPRP(AA,MAXY,PYR)
GO TO (2551,2552),IFUNC
2551 PYY1(NPER,IM) = WVNC(IM)*PYR(1)/DENOM(IM)
PYY2(NPER,IM) = PYR(2) / DENOM(IM)
GO TO 2552
2552 CONTINUE
PYY1(NPER,IM) = WVNC(IM) * PYR(1) / DENOM(IM)
PYY2(NPER,IM) = WVNC(IM) * PYR(2) / DENOM(IM)
PYY3(NPER,IM) = PYR(3)/DENOM(IM)
PYY4(NPER,IM) = PYR(4)/DENOM(IM)
2555 CONTINUE

```

```

2999 CONTINUE
GO TO 2000
9999 CONTINUE
REWIND 1
DO 3160 I = 1,NPER
3160 XX(I) = ALOG10(XX(I))
XMAX = XX(NPER)
LMAX = XMAX
IF(XMAX - LMAX) 115,115,116
116 LMAX = LMAX + 1
115 CONTINUE
LMIN = LMAX - 3
X1 = LMIN
9997 CONTINUE
READ(60,4) PHI,RAD,VER,SRC
IF(PHI-LT,0) GO TO 2001
C PHI = AZIMUTH AT WHICH SPECTRA IS DESIRED
C RAD -- GT 0 PLOT RADIAL SPECTRA OF RAYLEIGH WAVE
C -- LE 0 NO RADIAL SPECTRA PLOT
C VER -- GT 0 PLOT VERTICAL SPECTRA OF RAYLEIGH WAVE
C -- LE 0 NO VERTICAL SPECTRA PLOT
C (WHEN USING LOVE TAPE, RAD AND VER HAVE NO MEANING)
C SRC -- GT 0 SPECTRA MULTIPLIED BY SOURCE SPECTRUM S(CMEGA) =
C 1.0/CMEGA
C E.G., STEP RESPONSE OF SYSTEM SEISMIC MOMENT = 1.0
C DYN-E-CM
C -- LE 0 SPECTRA MULTIPLIED BY SOURCE SPECTRUM S(CMEGA) = 1.0
C E.G., IMPULSE RESPONSE OF SYSTEM
PHIR = DEGRAD * PHI
SINA = SIN(ALP - PHIR)
COSA = COS(ALP - PHIR)
SIND = SIN(DEL - PHIR)
COSD = COS(DEL - PHIR)
SINDP = SIN(2.*(DEL-PHIR))
SINAP = SIN(2.*(ALP-PHIR))
IFUNC = IFUN
YMAXR = 1.0E-38
YMAXV = 1.0E-38
DO 9995 J=1,MODEMX
NP = NPERM(J)
DO 9995 I=1,NP
PYR(1) = PYY1(I,J)
PYR(2) = PYY2(I,J)
PYR(3) = PYY3(I,J)
PYR(4) = PYY4(I,J)
GO TO (2710,2720),IFUNC
2710 CONTINUE
XREAL = 0.5 * PYR(1)*(FACT2*SINDP - FACT3*SINAP)
XIMAG = -0.5*PYR(2)*(FACT4*SIND - FACT5*SINA)
GO TO 2730

```

```

2720 CONTINUE
  XREAL = FACT1*PYR(4) - PYR(1)*(FACT2*COSD*COSD - FACT3*COSA*COSA)
  XIMAG = 0.5*(PYR(2)+PYR(3))*(FACT4*COSD-FACT5*COSA)
2730 CONTINUE
  XREAL = 1.0E 25 * XREAL .
  XIMAG = 1.0E 25 * XIMAG
  DUM = SQRT(XREAL**2 + XIMAG**2)
  DUM = 1.0E-25 * DUM
  IF(SRC.GT.0) DUM = DUM * TT(I) / 6.2831853
  IF(DUM.EQ.0) DUM = 1.0E-38
  IF(IFUNC.EQ.1) GO TO 9988
  IF(RAD.LE.0) GO TO 9989
  YY2(I,J) = DUM * ABS(ELIP(I,J))
  IF(YY2(I,J) .GT. YMAXR) YMAXR = YY2(I,J)
  YY2(I,J) = ALOG10(YY2(I,J))
9989 IF(VER.LE.0) GO TO 9995
9988 YY1(I,J) = DUM
  IF(YY1(I,J) .GT. YMAXV) YMAXV = YY1(I,J)
  YY1(I,J) = ALOG10(YY1(I,J))
9995 CONTINUE
  YMAXV = ALOG10(YMAXV)
  YMAXR = ALOG10(YMAXR)
  DO 9996 IC = 1,IFUNC
  IF(IFUNC.EQ.1) GO TO 9993
  IF(IC.EQ.1.AND.VER.LE.0) GO TO 9996
  IF(IC.EQ.2.AND.RAD.LE.0) GO TO 9996
9993 CONTINUE
  IF(IC.EQ.1) YMAX = YMAXV
  IF(IC.EQ.2) YMAX = YMAXR
  IF(IFUNC.EQ.1) LABLE(1) = LAB(3)
  IF(IFUNC.EQ.2.AND.IC.EQ.1) LABLE(1) = LAB(1)
  IF(IFUNC.EQ.2.AND.IC.EQ.2) LABLE(1) = LAB(2)
  XAXLEN = 7.535
  YAXLEN = 7.507
  DELTAX = 3./7.535
  DELTAY = 3./7.507
  SCP = 0.03 * YAXLEN
  LMAX = YMAX
  IF(YMAX = LMAX) 110,110,111
111 LMAX = LMAX + 1
110 CONTINUE
  LMIN = LMAX - 3
  Y1 = LMIN
  CALL ALOGAXES(XAXLEN,YAXLEN,3,3,LABEL,LABLE,12,18,X1,Y1,DELTAX,
+ DELTAY)
  DO 3171 IM = 1,MODEMX
  NP = NPERM(IM)
  DO 3172 IZ = 1,NP
  X(IZ) = XX(IZ)
  GO TO (3168,3169),IC

```

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```
3168 Y(IZ) = YY1(IZ,IM)
      GO TO 3170
3169 Y(IZ) = YY2(IZ,IM)
3170 CONTINUE
      IF(Y(IZ).LT.Y1) Y(IZ) = Y1
3172 CONTINUE
      X(NP+1)=X1
      X(NP+2) = DELTAX
      Y(NP+1) = Y1
      Y(NP+2) = DELTAY
3181 CALL LIND(X,Y,NP,IM)
3171 CONTINUE
      CALL SYMBOL( 2.6 ,YAXLEN+0.1,SCP,4HH = ,0,0,4)
      CALL NUMBER( 4.1 ,YAXLEN+0.1,SCP,AA,0,0,1)
      CALL SYMBOL( 2.1,YAXLEN + 0.6,SCP,4HPHI=,0.0,4)
      CALL NUMBER(4.35,YAXLEN+0.6,SCP,PHI,0,0,1)
      MMM = MODEMX - 1
      DO 3300 IM= 1,MMM
      FPN = IM
      IT = IM + 1
      XQ = XAXLEN - 1,
      YY = YAXLEN - 1, - (FPN - 1,) * 0.3
      CALL SYMBOL(XQ,YY+0.1,0.2,IT,0,0,-1)
      CALL SYMBOL(XQ+0.3,YY,0.2,1H=,0,0,1)
3300 CALL NUMBER(XQ+0.6,YY,0.2,FPN,0,0,-1)
      CALL PLOT(11.0,0.0,-3)
      CALL FACTOR(1,0)
      CALL PLOT(0.0,-11.0,-3)
      CALL PLCT(0.0,2,0,-3)
      CALL FACTOR(SIZE)
9996 CONTINUE
      GO TO 9997
9998 CONTINUE
      CALL PLOT(10.5,0,0,999)
      STOP
      END

SUBROUTINE NONPRP(A,MAXY,PYR)
COMMON/ONE/AY(100,4),CY(100,4),N,DEPTH(100),D(100),
1      LAYER(100),AJ(100),B(100),RHO(100),XMU(100),XLAMB(100),
1 IBUF(1026), ELIP(60,10)
DIMENSION PYR(4)
DO 100 I=1,N
IFOCA = I
D1 = DEPTH(I) - 0.5*D(I)                               D2 = DEPTH(I) + 0.5*D(I)
IF(A.GE.D1.AND.A.LE.D2) GO TO 101
100 CONTINUE
101 DO 110 JJ = 1,MAXY
      AA = AY(IFOCA,JJ)
```

```

CC = CY(IFCCA,JJ)
AA = AA - CC * DEPTH(IFCCA)
110 PYR(JJ) = AA + CC*A
RETURN
END

SUBROUTINE ALOGAXES(XAXLEN,YAXLEN,NOCX,NOCY,TTLX,TTLY,MTX,MTY,X1,
1Y1,DELTAX,DELTAY)
CHARACTER TTLX, TTLY
DIMENSION TTLX(10),TTLY(10)
SLT=.02*YAXLEN
SST=.01*YAXLEN
SP=-.06*YAXLEN
SS=.035*YAXLEN
SSP=SP+SS+.06
TTLP=-.11*YAXLEN-.1
STTL=.035*YAXLEN
XNUM = 1
YL = Y1
YU = Y1+NOCY
IF(ABS(YL),GE,10.,OR,ABS(YU),GE,10.) XNUM = XNUM + 1
IF(ABS(YL),GE,100.,OR,ABS(YU),GE,100.) XNUM = XNUM + 1.
IF( Y1 ,LT, 0 ) XNUM = XNUM + 1.0
CALL PLOT(-SLT,0.0,2)
CALL PLOT(0.0,-SLT,3)
CALL PLOT(0.0,0.0,2)
XP0=X1
YP0=Y1
IF(NOCX,EG.0)GO TO 4
ANOCX=NOCX
FACTX=XAXLEN/ANOCX
CALL SYMBOL(-,6*SS,SP,SS,2H10,0,0,2)
CALL NUMBER(999,, SSP,0.5*SS,X1,0.0,-1)
CALL PLOT(0.0,0.0,3)
DO 3 J=1,NOCX
DO 2 I=1,10
X=I
X=ALOG10(X)*FACTX+(J-1)*FACTX
IF(I.EQ.1)GO TO 2
CALL PLOT(X,0.0,2)
CALL PLOT(X,-SST,2)
2 CALL PLOT(X,0.0,3)
CALL PLOT(X,-SLT,2)
CALL SYMBOL(X-,6*SS,SP,SS,2H10,0,0,2)
XP0=XP0+1.0
CALL NUMBER(999,, SSP,0.5*SS,XP0,0.0,-1)
3 CALL PLOT(X,0.0,3)
XTL=MTX
XTL=(XAXLEN-XTL*STTL)/2.0
CALL SYMBOL(XTL,TTLP,STTL,TTLX,0.0,MTX)

```

```

          GO TO 6
1 CALL AXIS(0.0,0.0,TTLX,-MTX,XAXLEN,0.0,X1,DELTAX)
6 CALL PLOT(0.0,0.0,3)
      IF(NOCY.EQ.0)GO TO 10
                           ANOCY=NOCY
      SP = SP + (XNUM - 1.5) * 0.5 * SS
      TTLP = TTLP - (XNUM - 1)*0.5*SS
      FACTY=YAXLEN/ANOCY
      CALL SYMBOL(SP-,4,-.5*SS,SS,2H10,0,0,2)
      CALL NUMBER(999,,.5*SS-.06,,5*SS,Y1,0.0,-1)
                           CALL PLOT(0.0,0.0,3)
      DO 9 J=1,NOCY
          DO 8 I=1,10
              Y=I
              Y=ALOG10(Y)*FACTY+(J-1)*FACTY
              IF(I.EQ.1)GO TO 8
              CALL PLOT(0.0,Y,2)
                           CALL PLOT(-SST,Y,2)
8 CALL PLOT(0.0,Y,3)
              CALL PLOT(-SLT,Y,2)
              CALL SYMBOL(SP-,4,Y-.5*SS,SS,2H10,0,0,2)
              YPO=YPO+1.0
              CALL NUMBER(999,,Y+.5*SS-.06,,5*SS,YPO,0.0,-1)
9 CALL PLOT(0.0,Y,3)
              YTL=MTY
              YTL=(YAXLEN-YTL*STTL)/2.0
              CALL SYMBOL(TTLP-.2,YTL,STTL,TTLY,90.,MTY)
                           RETURN
10 CALL AXIS(0.0,0.0,TTLY,MTY,YAXLEN,90.,Y1,DELTAY)
      RETURN
                           END

SUBROUTINE LIND(X,Y,NP,IM)
DIMENSION X(1),Y(1)
X1 = X(NP+1)
Y1 = Y(NP+1)
DELTAX = X(NP+2)
DELTAY= Y(NP+2)
DO 100 I=1,NP
    XX = (X(I) - X1 )/DELTAX
    YY = (Y(I) - Y1)/DELTAY
    IF( I.EQ.1 ) GO TO 96
    IF( YY .LE. 0 ) GO TO 95
    IF( IM .EQ. 1 ) CALL PLOT(XX,YY,2)
    IF( IM .EQ. 1 ) GO TO 100
    CALL SYMBOL(XX,YY,0,14,IM,0.0,-2)
    GO TO 100
95 IF(Y(I-1) .LE. Y1 ) GO TO 105
    XXX = X(I) - (X(I) - X(I-1))/(Y(I)-Y(I-1))*Y(I)
    IF(XXX .LT. X(I-1) ,OR. XXX .GT. X(I)) XXX = 0.5*(X(I)+X(I-1))

```

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```
    XXX = (XXX-X1)/DELTAX
    CALL PLOT(XXX,0,0,2)
105 CONTINUE
    CALL PLOT(XX,0,2)
    IF( I.EQ.NP ) GO TO 100
102 IF( Y(I+1) .LE. Y1 ) GO TO 100
    XX = X(I) - (X(I)-X(I+1))/(Y(I)-Y(I+1))*Y(I)
    IF( XX .LT. X(I) ,OR, XX.GT. X(I+1)) XX = 0.5*(X(I)+X(I+1))
    XX = (XX-X1)/DELTAX
    CALL PLOT(XX,0,0,2)
    GO TO 100
96 IF( YY .LE. u ) GO TO 97
    IF( IM.EQ. 1 ) CALL PLOT(XX,YY,3)
    IF( IM .EQ. 1 ) GO TO 100
    CALL SYMBOL(XX,YY,0.14,IM,0.0,-1)
    GO TO 100
97 CALL PLOT(XX,0,3)
    GO TO 102
100 CONTINUE
    RETURN
    END
```

```
SUBROUTINE DIPDD(X,Y,Z,PHI,DELTA)
    TWPI = 6.2831853
    CCON = 57.29578
    VNORM = SQRT(X*X+Y*Y+Z*Z)
    IF(Z.GT.0) GO TO 10
    X = -X
    Y = -Y
    Z = -Z
10 IF(ABS(X).LE.0.0001.AND.ABS(Y).LT.0.001) GO TO 20
    DELTA = ARSIN (Z/VNORM)
    IF(ABS(X).LT.0.0001) GO TO 21
    PHI = ATAN(Y/X)
    IF(X.LT.0) PHI = PHI + TWPI/2,
    IF(PHI.LT.0) PHI = PHI + TWPI
    GO TO 30
21 IF(Y.LT.0) PHI = -TWPI/4, + TWPI
    IF(Y.GT.0) PHI = TWPI/4
    GO TO 30
20 DELTA = TWPI/4.
    PHI = 0
30 PHI = PHI * CCON
    DELTA = DELTA*CCON
    DELTA = 90. - DELTA
    RETURN
    END
```

SUBROUTINE NODPL(DEL,BET,ALP,GAM)

```

DIMENSION X(365),Y(365),DELTA(2),GAMMA(2)
CALL SYMBOL(2.5,3,25,0.1,1HN,0.0,1)
CALL PLOT(2.5,3,13,3)
CALL PLOT(2.5,2,93,2)
CALL PLOT(2.5,1,13,3)
CALL PLOT(2.5,0,93,2)
CALL PLOT(1,4,2,03,3)
CALL PLOT(1,6,2,03,2)
CALL PLOT(3,4,2,03,3)
CALL PLOT(3,6,2,03,2)
CALL SYMBOL(2,5,2,03,0,1,3,0,0,-1)
DEGRAD = 0.017452329
DO 111 I=1,361
ANGI = DEGRAD * ( I - 1 )
X(I) = 2.5 + COS(ANGI)
111 Y(I) = 2.03 + SIN(ANGI)
X(362) = 0.0
Y(362) = 0.0
X(363) = 1.0
Y(363) = 1.0
CALL LINE(X,Y,361,1,0,0)
PLOT P, T SYMBOLS
SQTWR = 1.414214
R = SQTWR * SIN(GAM/2.)
X1= 2.5 + R * SIN(ALP) - 0.03
Y1= 2.03 + R * COS(ALP) - 0.05
CALL SYMBOL(X1,Y1,0.1,1TP,0.0,1)
R = SQTWR * SIN(BET/2.)
X1= 2.5 + R * SIN(DEL) - 0.03
Y1= 2.03 + R * COS(DEL) - 0.05
CALL SYMBOL(X1,Y1,0.1,1TP,0.0,1)
TWPI = 6.2831853
DINC = 1.
NP = 179
X1 = SIN(BET)*COS(DEL)
Y1 = SIN(BET) * SIN(DEL)
Z1 = COS(BET)
X2 = SIN(GAM)*COS(ALP)
Y2 = SIN(GAM)*SIN(ALP)
Z2 = COS(GAM)
CALL DIPDD(X1+X2,Y1+Y2,Z1+Z2,GAMMA(1),DELTA(1))
CALL DIPDD(X1-X2,Y1-Y2,Z1-Z2,GAMMA(2),DELTA(2))
WRITE(6,1) ((GAMMA(L),DELTA(L)),L=1,2)
1 FORMAT(1H ,4F12.3)
DO 2 I = 1,2
GAMMA(I) = GAMMA(I) - 180,
IF(GAMMA(I).LT.0) GAMMA(I) = GAMMA(I) + 360,
2 CONTINUE
DO 100 K = 1,2
IF(DELTA(K).GE.1.) GO TO 50

```

```

NPOINT = NP+2
DO 40 I=1,NPOINT
  THETA =(GAMMA(K)-90.+(I-1)*DINC)*DEGRAD
  IF(THETA.LT.0) THETA = THETA + TWPI
  IF(THETA.GT.TWPI) THETA = THETA-TWPI
  X(I) = SIN(THETA) + 2.5
  Y(I) = COS(THETA) + 2.03
40 CONTINUE
GO TO 90
50 IF(ABS(DELTA(K)-90.) .GE. 1.) GO TO 70
  NPOINT = 2
  DO 60 I=1,2
    THETA =(GAMMA(K)-90.+(I-1)*180.)*DEGRAD
    IF(THETA.LT.0) THETA = THETA + TWPI
    IF(THETA.GT.TWPI) THETA = THETA - TWPI
    X(I) = SIN(THETA) + 2.5
    Y(I) = COS(THETA) + 2.03
60 CONTINUE
GO TO 90
70 THETA = (GAMMA(K)-90.) * DEGRAD
  IF(THETA.LT.0) THETA = THETA + TWPI
  X(1) = SIN(THETA) + 2.5
  Y(1) = COS(THETA) + 2.03
  THETA=(GAMMA(K)+90.) * DEGRAD
  IF(THETA.GT.TWPI) THETA = THETA - TWPI
  X(NP+2) = SIN(THETA) + 2.5
  Y(NP+2) = COS(THETA) + 2.03
  NPOINT = NP + 2
  ANG = (90.-DELTA(K)) * DEGRAD
  TANI = SIN(ANG) / COS(ANG)
  DO 80 I=1,NP
    THETA = GAMMA(K)-90. + I*DINC
    IF(THETA.LT.0) THETA = THETA + 360.
    IF(THETA.GT.360.) THETA = THETA - 360.
    ALPHA = (THETA-GAMMA(K)) * DEGRAD
    XJ = ATAN(TANI/COS(ALPHA))
    R = SQTWR * SIN(XJ/2.)
    X(I+1) = R*SIN(THETA*DEGRAD) + 2.5
    Y(I+1) = R*COS(THETA*DEGRAD) + 2.03
80 CONTINUE
90 CONTINUE
  X(NPOINT+1) = 0.0
  Y(NPOINT+1) = 0.0
  X(NPOINT+2) = 1.0
  Y(NPOINT+2) = 1.0
  CALL LINE(X,Y,NPOINT,1,0,0)
100 CONTINUE
RETURN
END

```

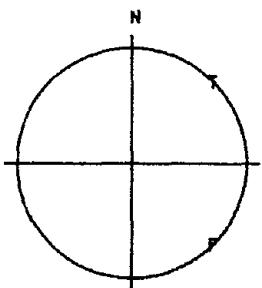
DATA PAGE 1

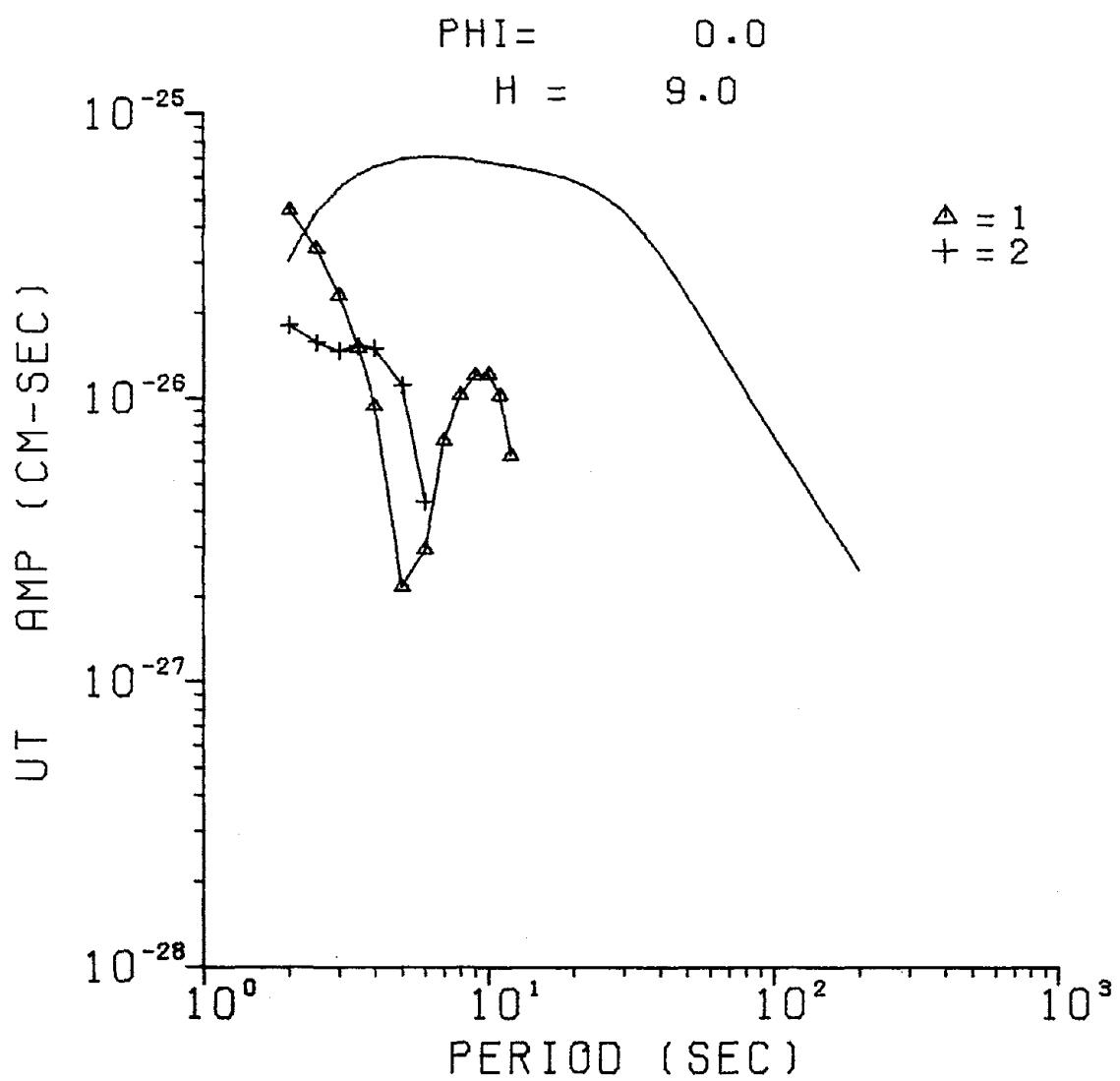
0.6			
45.	90.	135.	90.
9,			
0.0	1.0	1.0	1.0
-1.0			
-1.0			

45,000	90,000	135,000	90,000
89,995	89.993	179.995	90.000

T P

DELT	BETA	ALPH	GAM
45	90	135	90





IX. SRFWVSR

PROGRAMMER: R. B. HERRMANN / Jan 1973

PURPOSE:

This program accepts a focal mechanism and focal depth and compares theoretical amplitude spectra for an impulsive source with observed spectral amplitudes. The source spectrum at each period as well as the anelastic attenuation coefficients are determined.

INPUT/OUTPUT

The surface wave amplitude spectra are read from card on FILE 60. One eigenfunction tape is used on FILE 01. Output is on the printer, FILE 61.

THEORY

The source spectrum S and attenuation coefficient g are determined by a least squares process in order to fit the observations to the model

$$y = S A \exp(-gr),$$

where r is the distance from the source and A is the theoretical amplitude spectra at a particular station for the focal mechanism and focal depth assumed.

Herrmann (1974) and Herrmann and Mitchell (1975) established a criteria for accepting the quality of data for determining S and g . A correlation coefficient R is determined for each fit to the data. Also output is the critical R value -- RCRIT. The g and S determinations are accepted if R is greater than RCRIT. 95% confidence limits are given to G and S .

PROGRAM DESCRIPTION

PROGRAM SRFWVSR : This is the main input routine.

SUBROUTINE LINLSQ: This subroutine determines g and S together with confidence levels.

SUBROUTINE RCOEF: This determines the correlation coefficient between the Y vs A values.

SUBROUTINE NONPRP: This performs linear interpolation of the eigenfunctions with respect to depth.

SUBROUTINE FISHER; This determines the values of the Student "t" coefficient at a 95% confidence level.

REFERENCES

Herrmann, R. B. (1974). Surface-wave generation by central United States earthquakes, Ph. D. dissertation, Saint Louis University.

Herrmann, R. B. and B. J. Mitchell (1975). Statistical analysis and interpretation of surface-wave anelastic attenuation data for the stable interior of North America, Bull. Seism. Soc. Am. 65, 1115-1128.

INPUT DATA

Card Sequence	Column	Name	Format	Explanation
A.	1-10	DISTMX	F10.5	Only data at distances less than DISTMX are used to determine S and g
	11-20	EQEX	F10.5	.EQ.0 Earthquake data .NE.0 Explosion data
B.	2-5	NUMSRC	I4	Number of focal mechanisms to be entered (LE.10)
	12-15	NUMDPT	I4	Number of focal depths to be searched through (LE.10)

C. Focal mechanism parameters (NUMSRC cards)

1-10	XDEL	F10.5	Azimuth of T axis measured from north
11-20	XBET	F10.5	Angle T axis makes with respect to downward vertical axis
21-30	XALP	F10.5	AZ of P axis measured from north
31-40	XGAM	F10.5	Angle P axis makes with respect to downward vertical axis

D. Focal depth (NUMDPT cards)

1-10	AAA	F10.5	Focal depth
------	-----	-------	-------------

E. Data header

1-10	TT	F10.5	Period of data
------	----	-------	----------------

Card Sequence	Column	Name	Format	Explanation
Data header (cont'd)				
	11-15	MODAL	I5	Mode of data Fund = 1, 1'st = 2 etc.
	16-20	NUMSTA	I5	dummy set to 0 or blank
	21-25	IURUZ	I5	1 UZ data 2 UR data (has no meaning for Love wave data)
	30-40	ATTEN	E11.4	LT 0 search through admissible region then do least squares EQ 0 Do least squares only GT 0 Obtain source spectrum for given known attenuation value

F. Spectral amplitude data

2-5	ISTA	A4	Station identifier
11-20	AZ	F10.5	Station azimuth from source
21-30	DIST	F10.5	Station distance from source in km
31-40	WEIGHT	F10.5	Observational weight to data
41-51	AMPL	E11.4	Amplitude spectrum value normalized for geometrical spreading to 1000 km
61-70	GPV		Group velocity of amplitude arrival

The program now returns to Point F to read another amplitude observation. When a card is found for which ISTA = 'HALT' , then the program stops reading in data and proceeds with the least squares analysis.

When the least squares analysis is completed, the program returns to Point E to read in new period data. The program terminates when a value of TT is found which is negative.

The data must be arranged in order of increasing period which is the way the eigenfunctions are placed on tape.

For an explosion the XDEL, XBET, XALP, XGAM quantities have no meaning, so these values can be read to zero.

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C PROGRAM SRFWVSRC
C THIS PROGRAM COMPUTES THEORETICAL SURFACE WAVE TRANSMISSION
C FACTORS FOR RAYLEIGH OR LOVE WAVES AT A DISTANCE OF 1000
C KILOMETERS. OBSERVED SPECTRA NORMALIZED TO A DISTANCE OF 1000
C KILOMETERS ARE READ IN. A LINEAR LEAST SQUARES FIT IS MADE TO THE
C DATA TO OBTAIN THE SOURCE SPECTRUM AND ATTENUATION FACTORS
C SIMULTANEOUSLY. COMPUTATIONS CAN ALSO BE PERFORMED FOR SOURCE
C SPECTRUM ONLY WHEN ATTENUATION IS INPUTTED.
C RAYLEIGH EIGENFUNCTION TAPE ENABLES USE OF UZ OR UR
C LOVE EIGENFUNCTION TAPE ENABLES USE OF UT
C THE PROGRAM WILL ALSO DETERMINE THE SOURCE SPECTRUM AND THE
C GAMMA COEFFICIENTS FOR EXPLOSION DATA IN WHICH CASE THE DEL,BET,
C ALP,GAM PARAMETERS HAVE NO MEANING AND MAY ALL BE SET TO ZERO
COMMON AY(100,4),CY(100,4),N,DEPTH(100),D(100)
COMMON LAYER(100),A(100),B(100),RHO(100),XMU(100),XLAMB(100)
COMMON ELIP
DIMENSION AMODE(10),WVNC(10),ELIP(10), ABSORP(10),DENOM(10)
DIMENSION XDEL(10),XBET(10),XALP(10),XGAM(10),FACT1(10),FACT2(10)
DIMENSION FACT3(10),FACT4(10),FACT5(10),AAA(10),ISTA(50),AZ(50)
DIMENSION DIST(50),AMPL(50),WEIGHT(50),YY(50)
DIMENSION PYR(4),PYI(4)
DIMENSION PYY1(10,10),PYY2(10,10),PYY3(10,10),PYY4(10,10)
DIMENSION GPV(50)
DIMENSION UU(10)
IHALT = 4HHALT
1 FORMAT(1H ,2(14,6X))
2 FORMAT(1H ,14,5X,7F10.4)
4 FORMAT(8F10.5)
5 FORMAT(1H ,A4,3F10.5,5E11.4)
6 FORMAT(1H ,7(4X,E11.4))
7 FORMAT(F10.5,31S,4X,E11.4)
8 FORMAT(1H1,20X,5HF0CAL,11X,4HDEL ,6X,4HBET ,6X,4HALP ,6X,4HGAM )
9 FORMAT(1H ,20X,15,5X,4(4X,F6.1))
10 FORMAT(1X,A4,5X,3F10.5,E11.4,9X,F10.5)
DEGRAD = 0.017452329
      READ(60,4) DISTMX, EGEX
C EQEX .EQ. 0.0 EARTHQUAKE DATA
C EQEX .NE. 0.0 EXPLOSION DATA
READ(60,1) NUMSRC,NUMDPT
DO 25 I = 1,NUMSRC
25 READ(60,4) XDEL(I),XBET(I),XALP(I),XGAM(I)
      WRITE(61,8)
DO 26 I = 1,NUMSRC
26 WRITE(61,9) I,XDEL(I),XBET(I),XALP(I),XGAM(I)
DO 27 I = 1,NUMSRC
XDEL(I) = DEGRAD * XDEL(I)
XBET(I) = DEGRAD * XBET(I)
XALP(I) = DEGRAD * XALP(I)
XGAM(I) = DEGRAD * XGAM(I)
FACT1(I) = (COS(XBET(I))**2) - (COS(XGAM(I))**2)

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SREFWVSR0 PAGE 2

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FACT2(I) = (SIN(XBET(I)))**2
FACT3(I) = (SIN(XGAM(I)))**2
FACT4(I) = SIN(2.*XBET(I))
27 FACT5(I) = SIN(2.*XGAM(I))
DO 28 I = 1,NUMDPT
28 READ(60,4) AAA(I)
C   TT = PERIOD OF DATA
C   MODAL = MODE OF DATA F = 1, 1ST = 2, ETC
C   NUMSTA = DUMMY
C   IURUZ 1 = UZ DATA 2 = UR DATA
C   ATTEN LT 0 SEARCH
C       EQ 0 LINEAR LEAST SQUARES ONLY
C       GT 0 OBTAIN SOURCE FOR GIVEN ATTEN
XXMAX = 0.0
READ(60,7) TT,MODAL,NUMSTA,IURUZ,ATTEN
IF(TT.LE.0) GO TO 9999
DO 29 I = 1,100
NUMSTA = I
READ(60,10) ISTA(I),AZ(I),DIST(I),WEIGHT(I),AMPL(I),GPV(I)
IF(ISTA(I).EQ.IHALT) GO TO 291
IF(AMPL(I).GT.XXMAX) XXMAX = AMPL(I)
29 CONTINUE
GO TO 292
291 NUMSTA = NUMSTA - 1
292 CONTINUE
C   DELETION OF AMPLITUDES NEAR NODES OF RADIATION PATTERN
XLIM = 0.2 * XXMAX
J = 1
DO 4000 I = 1,NUMSTA
IF(AMPL(I).LT.XLIM) GO TO 4000
C   IF(DIST(I).GT.DISTMX) GO TO 4000
C   ISTA(J) = ISTA(I)
C   AZ(J) = AZ(I)
C   DIST(J) = DIST(I)
C   WEIGHT(J) = WEIGHT(I)
C   AMPL(J) = AMPL(I)
C   GPV(J) = GPV(I)
J = J + 1
4000 CONTINUE
NUMSTA = J - 1
DO 37 I = 1,10
DO 37 J = 1,10
PYY1(I,J) = 0
PYY2(I,J) = 0
PYY3(I,J) = 0
37 PYY4(I,J) = 0
READ(01,1)N
C   READ IN MODEL
DO 100 I=1,N
100 READ(01,2) LAYER(I),DEPTH(I),D(I),A(I),B(I),RHO(I),XMU(I),XLAMB(I)
C
CHANGE MADE HERE
```

```

1)
2000 READ(01,1) IFUNC, MODE
      MAXY = 2 * IFUNC
      IF(IFUNC.LT.0) GO TO 9999
      DO 2999 IM = 1,MODE
      READ(01,5) AMODE(IM),T,C,U,SUMIO,WVNO(IM),AL,ELIP(IM),ABSORP(IM)
      UU(IM) = 0
C      THE SPECTRA ARE COMPUTED AT A DISTANCE OF 1000 KILOMETERS
      DENOM(IM) = SQRT (WVNO(IM))*0.3989423
      DENOM(IM) = DENOM(IM)/AL
      DENOM(IM) = DENOM(IM) * 6.2831853 * 31.622776
      DO 2100 IN=1,N
2100 READ(01,6) AY(IN,1),AY(IN,2),AY(IN,3),AY(IN,4),DCDA,DCDB,DCDR
      IF(TT.NE.T) GO TO 2999
      WVNO2 = WVNO(IM)*#2
      IF(IFUNC-2)2400,2200,2200
C      RAYLEIGH WAVES
C      INITIALLY AY(IN,1) = UR, AY(IN,2) = UZ, AY(IN,3) = TZ,
C      AY(IN,4) = TR
C      NOW AY(IN,1) = UR, AY(IN,2) = UZ, AY(IN,3) = DURDZ,
C      AY(IN,4) = DUZDZ
C      AND CY(IN,1) = DURDZ, CY(IN,2) = DUZDZ, CY(IN,3) = D2URDZ2,
C      CY(IN,4) = D2UZDZ2.
C2200 DO 2300 IN = 1,N
      H = RHO(IN)*C*C*WVNO2
      XL2M = XLAMB(IN) + 2.*XMU(IN)
      CY(IN,1) = -WVNO(IM)*AY(IN,2) + AY(IN,4)/XMU(IN)
      CY(IN,2) = (WVNO(IM)*XLAMB(IN)*AY(IN,1)+AY(IN,3))/XL2M
      CY(IN,3) = (-H/XMU(IN) + WVNO2*(3.*XL2M-2.*XMU(IN))/XL2M)*AY(IN,1)
      CY(IN,4) = (-H+WVNO2*XLAMB(IN))/AY(IN,3)+WVNO(IM)*(1.+XLAMB(IN)/
      1.*XMU(IN))*AY(IN,4))/XL2M
      AY(IN,3) = CY(IN,1)
      AY(IN,4) = CY(IN,2)
2300 CONTINUE
      GO TO 2501
C      LOVE WAVES
C      INITIALLY AY(IN,1) = UT, AY(IN,2) = TT
C      NOW AY(IN,1) = UT, AY(IN,2) = DUTDZ
C      AND CY(IN,1) = DUTDZ, CY(IN,2) = D2UTDZ2
2400 DO 2500 IN=1,N
      CY(IN,1) = AY(IN,2)/XMU(IN)
      CY(IN,2) = WVNO2*(XMU(IN)- C*C*RHO(IN))*AY(IN,1)/XMU(IN)
      AY(IN,2) = CY(IN,1)
2500 CONTINUE
2501 CONTINUE
      DO 2555 ID = 1,NUMDPT
      AA = AAA(ID)
      NPER = ID
      IF(AA.LT.0) AA = 0.0

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      CALL NONPRP(AA,MAXY,PYR)
      GO TO(2551,2552),IFUNC
2551 PYY1(NPER,IM) = WVNC(IM)*PYR(1)/DENOM(IM)
      PYY2(NPER,IM) = PYR(2) / DENOM(IM)
      GO TO 2555
2552 CONTINUE
      PYY1(NPER,IM) = WVNC(IM) * PYR(1) / DENOM(IM)
      PYY2(NPER,IM) = WVNC(IM) * PYR(2) / DENOM(IM)
      PYY3(NPER,IM) = PYR(3)/DENOM(IM)
      PYY4(NPER,IM) = PYR(4)/DENOM(IM)
2555 CONTINUE
2999 CONTINUE
      IF(TT.NE.T) GO TO 2999
2998 CONTINUE
      IF(MODAL.LF.0.OR.MODAL.GT.MODE) GO TO 2997
      RAD = 0
      VER = 0
      IF(IFUNC.EQ.1) GO TO 2633
      IF(IURUZ-2)2631,2632,2632
2631 VER = 1
      GO TO 2633
2632 RAD = 1.
2633 CONTINUE
      DO 3000 I = 1,NUMDPT
      DO 3001 J = 1*NUMSRC
      WRITE(61,11) AAA(I),J,TT,UU(MODAL),AMODE(MODAL)
11 FORMAT(1H ,10X,8HDEPTH = ,F10.3,10X,9HSOURCE = , I3,10X,4HT = ,
     1 F10.2,10X,4HU = ,F10.5,10X,A4/)
      DEL = XDEL(J)
      ALP = XALP(J)
      BET = XBET(J)
      GAM = XGAM(J)
      DO 3002 K = 1*NUMSTA
      PYR(1) = PYY1(I,MODAL)
      PYR(2) = PYY2(I,MODAL)
      PYR(3) = PYY3(I,MODAL)
      PYR(4) = PYY4(I,MODAL)
      IF(EQEX.NE.0.0) GO TO 2725
      PHIR = AZ(K) * DEGRAD
      SINR = SIN(ALP - PHIR)
      COSR = COS(ALP - PHIR)
      SIND = SIN(DEL - PHIR)
      COSD = COS(DEL - PHIR)
      SINDP = SIN(2.*(DEL-PHIR))
      SINAP = SIN(2.*(ALP-PHIR))
      IF(IFUNC-2)2710,2720,2720
2710 CONTINUE
      XREAL = 0.5 * PYR(1) *(FACT2(J)*SINDP-FACT3(J)*SINAP)
      XIMAG = -0.5*PYR(2)*(FACT4(J)*SIND-FACT5(J)*SINA)
      GO TO 2730

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2720 CONTINUE
  XREAL=FACT1(J)*PYR(4)-PYR(1)*(FACT2(J)*COSD*COSD-FACT3(J)*COSA*COS
  1A)
  XIMAG = 0.5*(PYR(2)+PYR(3))*(FACT4(J)*COSD-FACT5(J)*COSA)
2730 CONTINUE
  GO TO 2731
2725 XREAL = PYR(4) - PYR(1)
  XIMAG = 0.0
  IF(IFUNC,EQ,1) XREAL = 0.0
2731 CONTINUE
  DUM = SQRT(XREAL**2 + XIMAG**2)
  IF(DUM.EQ.0) WEIGHT(K) = 0.0
  IF(IFUNC,EQ,1) GO TO 9988
  IF(RAD.LE,0) GO TO 9989
  YY(K) = DUM * ABS(ELIP(MODAL))
9989 IF(VER.LE,0) GO TO 9995
9988 YY(K) = DUM
9995 CONTINUE
3002 CONTINUE
C
  CALL LINLSQ(NUMSTA,WEIGHT,DIST,YY,AMPL,AZ,ISTA,ATTEN,GPV,TT)
3001 CONTINUE
3000 CONTINUE
  XXMAX = 0.0
2997 READ(60,7) TT,MODAL,NUMSTA,IURUZ,ATTEN
  IF(TT.LE,0) GO TO 9999
  DO 31 I = 1,100
    NUMSTA = I
    READ(60,10) ISTA(I),AZ(I),DIST(I),WEIGHT(I),AMPL(I),GPV(I)
    IF(ISTA(I).EQ,THALT) GO TO 311
    IF(AMPL(I).GT,XXMAX) XXMAX = AMPL(I)
  31 CONTINUE
  GO TO 312
  311 NUMSTA = NUMSTA - 1
  312 CONTINUE
C  DELETION OF AMPLITUDES NEAR NODES OF RADIATION PATTERN
  XLIM = 0.2 * XXMAX
  J = 1
  DO 5000 I = 1,NUMSTA
    C  IF(AMPL(I).LT.XLIM) GO TO 5000
    IF(DIST(I).GT,DISTMX) GO TO 5000
    ISTA(J) = ISTA(I)
    AZ(J) = AZ(I)
    DIST(J) = DIST(I)
    WEIGHT(J) = WEIGHT(I)
    AMPL(J) = AMPL(I)
    GPV(J) = GPV(I)
    J = J + 1
  5000 CONTINUE
  NUMSTA = J - 1

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```

IF(TT.EQ.T) GO TO 2998
GO TO 2000
9999 CONTINUE
REWIND 1
STOP
END

SUBROUTINE LINLSQ(N,W,XX,YY,A,AZ,ISTA,ATEN,GPV,TT)
C THIS SUBROUTINE SEARCHES FOR MINIMUM OF STATISTIC SHOWN AND GIVES
C 95 PERCENT CONFIDENCE LEVELS
DIMENSION W(1),XX(1),YY(1),A(1),AZ(1),ISTA(1),Y(52),X(52),ERR(52)
DIMENSION ISIGN(2)
DIMENSION IAST(20)
DIMENSION GPV(1),RAD(52)
DIMENSION D(52)
ISIGN(1) = 4H ++
ISIGN(2) = 4H X/
200 FORMAT(1H ,1HS= ,E11.4,A4,E11.4,1HX,4HK= ,E11.4,4H +- ,E11.4
1,10X,8HSIGMA = ,E11.4,10X,4HR = ,E11.4)
201 FORMAT(1H ,A4,5X,F10.2,5X,F10.2,4(4X,E11.4),2F10.4,4X,E11.4)
202 FORMAT(1H ,4H STA,11X,4HAZ ,9X,6HDIST ,4X,
1 11HRAD PATTERN,4X,11H OBSERVED ,4X,11H PREDICTED ,4X,2X,A4,
2 5HDEV ,4X,6HWEIGHT ,4X,6HGP VEL ,3X,12HCORR OBS RAD /)
203 FORMAT(1H0,8HNUMSTA =,I5,5X,18HESTIMATED MOMENT =,E11.4,5X,7HRCRIT
1 = ,E11.4/)
ATTEN = ATEN
IMP = 0
JUMP = 1
IF(ATEN.EQ.0) GO TO 5000
IF(ATEN.LT.0) JUMP = 2
IF(ATEN.LT.0) ATEN = -0.00004
5001 CONTINUE
GO TO (301,302,301),JUMP
302 IMP = IMP + 1
ATTEN = ATEN + 0.00004
IF(IMP.EQ.27) JUMP = 1
IF(IMP.EQ.27) GO TO 5000
301 CONTINUE
C W*(A - S*Y)**2 = W * (A - S*YY*EXP(-ATTEN*XX))**2 = MIN
SUM1 = 0
SUM2 = 0
ISIG = 1
DO 10 I = 1,N
Y(I) = YY(I) * EXP(-ATTEN*XX(I))
SUM1 = SUM1 + W(I) * A(I) * Y(I)
10 SUM2 = SUM2 + W(I) * Y(I) * Y(I)
S = SUM1/SUM2
SUM0 = 0
SUM1 = 0
SUM2 = 0

```

```

SUM0 = 0.0
DO 20 I = 1,N
  SUM0 = SUM0 + W(I)
  SUM1 = SUM1 + W(I) * Y(I)
  SUM2 = SUM2 + W(I)*Y(I)*Y(I)
  Y(I) = S * Y(I)
  ERR(I) = A(I) - Y(I)
20 SUM3 = SUM3 + W(I) * ERR(I) * ERR(I)
T = 1.960
C   TO PREVENT DIVISION BY ZERO
IF(N.EQ.1) N = 100
SIGMA = SQRT(SUM3/(N-1))
IF(N.EQ.100) N = 1
DS = T + SIGMA *(SQRT(SUM0/ABS(SUM0+SUM2-SUM1*SUM1)))
DS=DS*1.0E20
SLOPE = ATTEN
DSLOPE = 0.9
IF(JUMP.EQ.1) GO TO 6000
SS=S*1.0E20
WRITE(61,200)SS,ISIGN(ISIG),DS,SLOPE,DSLOPE,SIGMA
B1 = - SLOPE
B0 = ALOG(S)
5000 SUM1 = 0
SUM2 = 0
SUM3 = 0
SUM4 = 0
SUM5 = 0
SUM6 = 0
C   W * (ALOG(A/YY) - B0 - B1*XX) ** 2 = MIN
DO 100 I = 1,N
  Q(I) = ALOG(A(I)/YY(I))
  SUM1 = SUM1 + W(I) * XX(I) * XX(I)
  SUM2 = SUM2 + W(I) * XX(I)
  SUM3 = SUM3 + W(I)
  SUM4 = SUM4 + W(I) * Q(I)
  SUM5 = SUM5 + W(I) * Q(I) * XX(I)
100 SUM6 = SUM6 + W(I) * Q(I) * Q(I)
DET = SUM3 * SUM1 - SUM2 * SUM2
ISIG = 2
GO TO (401,402,402),JUMP
401 CONTINUE
  B0 = (SUM4 * SUM1 - SUM2 * SUM5)/DET
  B1 = (-SUM4*SUM2 + SUM3 * SUM5)/DET
402 CONTINUE
C   TO PREVENT DIVISION BY ZERO
IF(N.LE.2) N = 200
SIGMA2 = (SUM6 - B0*SUM4 - B1*SUM5)/(N-2)
IF(N.EQ.200) N = 2
T = 100.
IF(N.GT.2) CALL FISHER(N-2,T)

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```

T2 = T * T
XC = 100000,
IF(N.GT.2) XC = T2/(N-2)
RCRIT = SQRT(XC)/SQRT(1.+XC)
S = EXP(BU)
SIGMA = SQRT(ABS(SIGMA2))
SEBU = T * SQRT(ABS(SUM1*SIGMA2/DET))
DS = EXP(SEBU)
GO TO (501,502,502),JUMP
501 CONTINUE
DO 101 I = 1,N
ERR(I) = EXP(Q(I)-B0-B1*XX(I))
101 Y(I) = YY(I)*S*EXP(B1*XX(I))
502 CONTINUE
SEB1 = T * SQRT(ABS(SUM3*SIGMA2/DET))
SLOPE = - B1
DSLOPE = SEB1
6000 CONTINUE
CALL RCOEF(N,W,Y,A,RCOEFF)
SS=5*1.0E20
WRIT(E(61,200)SS , ISIGN(ISIG),DS,SLOPE,DSLOPE,SIGMA,RCOEFF
IF(JUMP.EQ.2) GO TO 5001
IF(JUMP.EQ.2.OR.JUMP.EQ.3) ISIG = 1
ESTMOM = SS * 6.2831853 / TT
WRITE(61,203) N,ESTMOM,RCRIT
WRITE(61,202) ISIGN(ISIG)
DO 5999 I = 1,N
5999 RAD(I) = A(I) * YY(I) / Y(I)
C BECAUSE OF THE NORMALIZATION USED YY(I) AND RAD(I) ARE
C 1.0E20 TOO LARGE
DO 6001 I = 1,N
6001 WRITE(61,201) ISTA(I),AZ(I),XX(I),YY(I),A(I),Y(I),ERR(I),W(I),
1 GPV(I),RAD(I)
IF(SLOPE.LT.0.OR.SLOPE.GT.0.01) ATTEN = 0,0
IF(SLOPE.LT.0.OR.SLOPE.GT.0.01) JUMP = 3
IF(SLOPE.LT.0.OR.SLOPE.GT.0.01) GO TO 5001
RETURN
END

SUBROUTINE RCOEF(N,W,Y,A,RCOEFF)
DIMENSION W(1),Y(1),A(1)
IF(N.EQ.1) RCOEFF = 1.0
IF(N.EQ.1) RETURN
SUM0 = 0
SUM1 = 0
SUM2 = 0
SUM3 = 0
SUM4 = 0
SUM5 = 0
DO 601 I = 1,N

```

```

SUM0 = SUM0 + W(I)
SUM1 = SUM1 + W(I) * Y(I)
SUM2 = SUM2 + W(I) * Y(I) * Y(I)
SUM3 = SUM3 + W(I) * A(I)
SUM4 = SUM4 + W(I) * A(I) * A(I)
601 SUM5 = SUM5 + W(I) * Y(I) * A(I)
      RCOEFF = (SUM0*SUM5 - SUM1*SUM3)/SORT((SUM0*SUM2 - SUM1*SUM1)*
      1 (SUM4 * SUM0 - SUM3 * SUM3))
      RETURN
      END

SUBROUTINE NONPRP(A,MAXY,PYR)
COMMON AY(100,4),CY(100,4),N,DEPTH(100),D(100)
COMMON LAYER(100),AJ(100),B(100),RHO(100),XMU(100),XLAMB(100)
COMMON IBUF(1026)
COMMON ELIP(50,10)
DIMENSION PYR(4)
DO 100 I=1,N
IFOCIA = I
D1 = DEPTH(I) - 0.5*D(I)
D2 = DEPTH(I) + 0.5*D(I)
IF(A.GE.D1.AND.A.LE.D2) GO TO 101
100 CONTINUE
101 DO 110 JJ = 1,MAXY
      AA = AY(IFOCIA,JJ)
      CC = CY(IFOCIA,JJ)
      AA = AA - CC * DEPTH(IFOCIA)
110 PYR(JJ) = AA + CC*A
      RETURN
      END

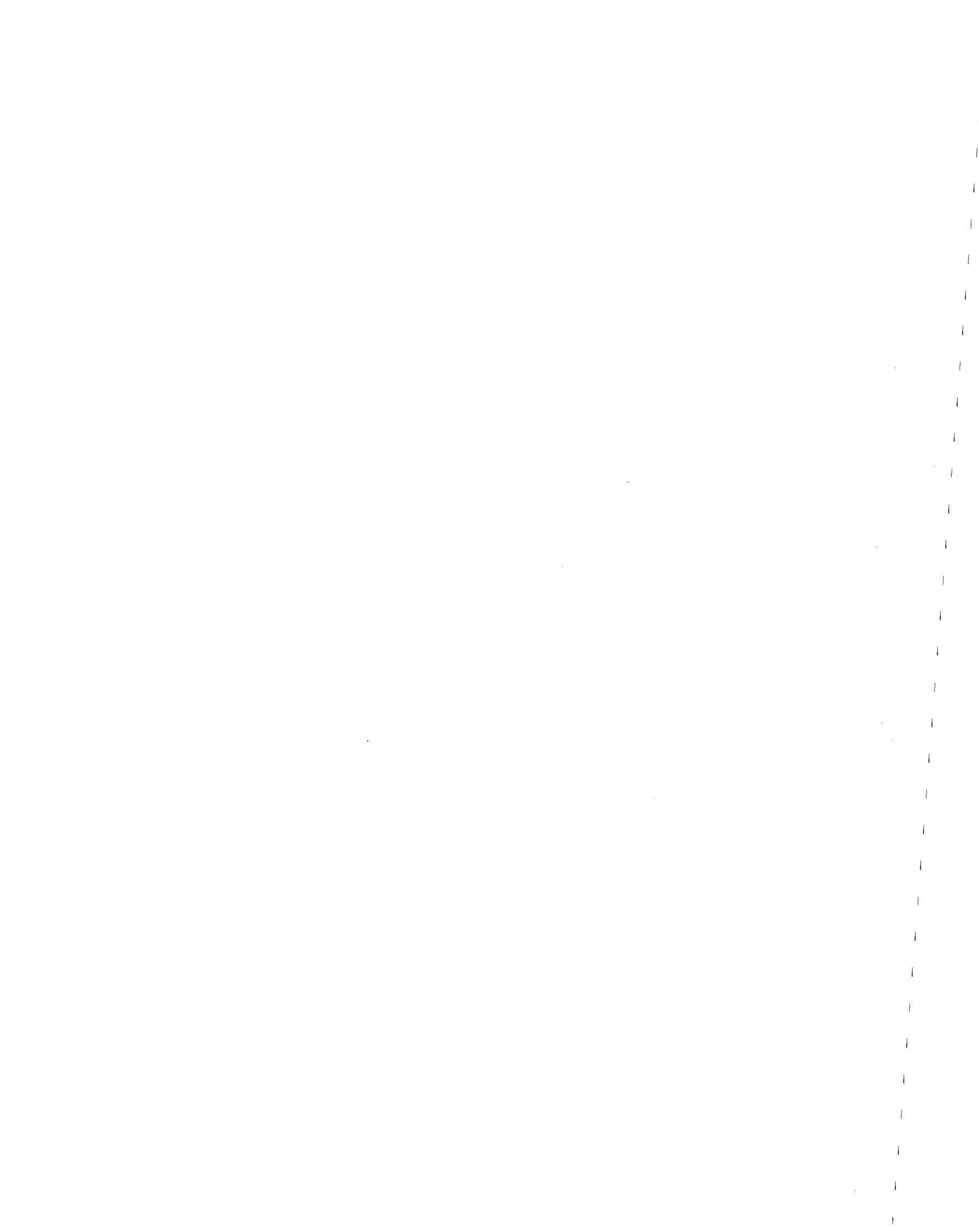
SUBROUTINE FISHER(N,T)
DIMENSION FT(34)
FT(1) = 12.706
FT(2) = 4.303
FT(3) = 3.182
FT(4) = 2.776
FT(5) = 2.571
FT(6) = 2.447
FT(7) = 2.365
FT(8) = 2.306
FT(9) = 2.262
FT(10) = 2.228
FT(11) = 2.201
FT(12) = 2.179
FT(13) = 2.160
FT(14) = 2.145
FT(15) = 2.131
FT(16) = 2.120
FT(17) = 2.110

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```
FT(18) = 2.101
FT(19) = 2.093
FT(20) = 2.086
FT(21) = 2.080
FT(22) = 2.074
FT(23) = 2.069
FT(24) = 2.064
FT(25) = 2.060
FT(26) = 2.056
FT(27) = 2.052
FT(28) = 2.048
FT(29) = 2.045
FT(30) = 2.042
FT(31) = 2.021
FT(32) = 2.000
FT(33) = 1.980
FT(34) = 1.960
I = N - 2
IF(I.LE.0) I = 1
IF(I.GT.30) GO TO 100
T = FT(I)
RETURN
100 IF(I.GT.30.AND.I.LE.40) T = (I-30)*(FT(31)-FT(30))/10,+FT(30)
    IF(I.GT.40.AND.I.LE.60) T = (I-40)*(FT(32)-FT(31))/20,+FT(31)
    IF(I.GT.60.AND.I.LE.120) T= (I-60)*(FT(33)-FT(32))/60,+FT(32)
    IF(I.GT.120) T = FT(34)
RETURN
END
```

3500.	0.0					
1	1					
181.	52.	272.	89.			
12.						
16.	1	0	1			
SLMT	3.2	335.3	1.0	6.4891E-04	3.3989	16.
LHCT	3.6	1425.30	1.0	6.8459E-04	3.4802	16.
AAMT	36.2	941.4	1.0	5.1925E-03	3.4276	16.
BLOT	40.3	526.4	1.0	4.4153E-03	3.4484	16.
OTTT	44.4	1650.	1.0	4.9422E-03	3.4836	16.
MNTT	47.2	1793.	1.0	4.5981E-03	3.4970	16.
CLET	47.4	1014.30	1.0	4.4010E-03	3.3878	16.
STJT	55.1	3574.1	1.0	4.2669E-03	3.4486	16.
WEST	59.8	1815.0	1.0	5.1962E-03	3.4810	16.
BLAT	76.0	917.1	1.0	2.3541E-03	3.5017	16.
ATLT	111.6	612.0	1.0	3.4088E-03	3.4991	16.
LUBT	261.1	1069.7	1.0	2.9095E-03	3.2037	16.
TUCT	264.6	1916.8	1.0	1.6224E-03	3.2604	16.
ALQT	271.7	1455.8	1.0	1.9990E-03	3.3494	16.
DUGT	291.2	2024.7	1.0	3.2214E-03	3.3583	16.
NEWT	303.7	2489.9	1.0	5.5236E-03	3.2151	16.
LONT	305.1	2881.8	1.0	3.3148E-03	3.3513	16.
MSOT	309.7	2316.3	1.0	5.3376E-03	3.3697	16.
PNTT	311.8	2613.9	1.0	5.7971E-03	3.3968	16.
EDMT	324.7	2651.9	1.0	7.7983E-03	3.3830	16.
FFCT	340.9	2307.3	1.0	2.4308E-03	3.4970	16.
MBCT	350.2	4758.5	1.0	1.7778E-03	3.4853	16.
HALT						
-1.						

S =	0.3902E 24	X/	0.1409E 01	DEPTH =	12.000	SOURCE =	1	T =	16.00	U =	3'48934
		K=	0.1168E-03	+	0.1610E-03	SIGMA =	0.3694E 06	FUND	R =	0.8267E 00	
FOCAL	181.6	BET	52.0	ALP	272.0	GAM	89.0				
NUMSTA =	22	ESTIMATED MOMENT =	0.1532E 24	RCRIT =	0.4252E 00	PREDICTED					
STA	AZ	DIST	RAD PATTERN	OBSERVED		X/Z DEV					
SLMT	3.20	335.30	0.1046E-06	0.8489E-03	0.3924E-03	0.2163E 01					
LHCT	3.60	1425.30	0.1302E-06	0.6846E-03	0.4302E-03	0.1591E 01					
AAMT	36.20	941.40	0.1732E-05	0.5193E-02	0.6054E-02	0.6633E-02	0.6653E 00				
BLOT	40.30	526.40	0.1809E-05	0.4415E-02	0.4942E-02	0.5950E-02	0.8307E 00				
OTT	44.40	1650.00	0.1849E-05	0.4598E-02	0.5870E-02	0.6428E-02	0.6833E 00				
MNTT	47.20	1793.00	0.1855E-05	0.4401E-02	0.4267E-02	0.4678E-02	0.5267E-02				
CLET	47.40	1014.30	0.1854E-05	0.4401E-02	0.4267E-02	0.4678E-02	0.9121E 00				
STJT	55.10	3374.10	0.1778E-05	0.5196E-02	0.2354E-02	0.3470E-02	0.6784E 00				
WEST	59.60	1815.00	0.1668E-05	0.9898E-06	0.1210E-05	0.3409E-02	0.4395E-02				
BLAT	76.00	917.10	0.1190E-05	0.1674E-05	0.2910E-02	0.2450E-02	0.2450E-02				
ATLT	111.60	612.00	0.7115E-06	0.5186E-06	0.1622E-02	0.1618E-02	0.1618E-02				
LUBT	261.10	1069.70	0.5186E-06	0.5186E-06	0.1622E-02	0.1618E-02	0.1618E-02				
TUCT	264.60	1916.80	0.5186E-06	0.5186E-06	0.1622E-02	0.1618E-02	0.1618E-02				
ALQT	271.70	1455.80	0.2627E-06	0.1999E-02	0.8647E-03	0.2312E 01	0.2312E 01				
DUGT	291.20	2024.70	0.1190E-05	0.3221E-02	0.3666E-02	0.8787E 00	0.8787E 00				
NEWT	303.70	2489.90	0.1674E-05	0.5524E-02	0.4885E-02	0.1131E 01	0.1131E 01				
LONT	305.10	2881.80	0.1711E-05	0.3315E-02	0.4768E-02	0.6952E 00	0.6952E 00				
MSOT	309.70	2316.30	0.1802E-05	0.5338E-02	0.5366E-02	0.9948E 00	0.9948E 00				
PNTT	311.61	2813.90	0.1829E-05	0.5797E-02	0.5138E-02	0.1128E 01	0.1128E 01				
EDMT	324.70	2651.90	0.1778E-05	0.7798E-02	0.5096E-02	0.1532E 01	0.1532E 01				
FFCT	340.90	2307.30	0.1224E-05	0.2431E-02	0.3647E-02	0.6665E 00	0.6665E 00				
MBCF	350.20	4758.50	0.7173E-06	0.1778E-02	0.1606E-02	0.1107E 01	0.1107E 01				



X. RADPAT

PROGRAMMER: R. B. HERRMANN / May 1973

PURPOSE:

This program accepts observed surface-wave amplitude spectra for a given period, corrects for amplitude decrease due to anelastic attenuation, and plots the amplitudes as a function of receiver azimuth. For a given value of the source spectrum at the same period, orientation of the pressure and tension axes, and focal depth, the theoretical surface wave radiation pattern is plotted.

The radiation patterns are scaled to fit in a square 2 3/4 inches on a side. Beneath each radiation pattern, the period is written, as well as a scale relating the plot size to units of dyne-cm.

The amplitude apectra used by the program are values at a reference distance of 9°.

The spectra can be aligned on the page by specifying an (II,JJ) coordinate for the plot location. The (II,JJ) specifications are

(1,1)	(1,2)	(1,3)
(2,1)	(2,2)	(2,3)
(3,1)	(3,2)	(3,3)

To plot more than nine patterns, or to plot another page, just make the next II or JJ in the (II,JJ) pair negative. The program recognizes the negative quantity, moves the origin of the plotting system 13 inches, and plots the radiation pattern in the location (ABS(II), ABS(JJ)).

INPUT / OUTPUT

The eigenfunction tape are read from FILE 01. The observed surface-wave amplitude spectra are read from card on FILE 60. The FILE 10 is used for CALCOMP plots. If the card data is for Love waves, the eigenfunction tape must be the one containing the Love wave eigenfunctions. The program also requires the subroutines NONPRP, EIGEN, MODEL. These subroutines are stored separately on disk. Printer output is on FILE 61.

INPUT DATA

Card Sequence	Column	Name	Format	Explanation
A.				
	13-72	(IFORM(I), 15A4 I=1,15)		Format for spec- tral amplitude data input
B.				
	1-10	XDEL	F10.5	Azimuth of T axis measured clock- wise from North
	11-20	XBET	F10.5	The angle the T axis makes with the downward vertical z-axis
	21-30	XALP	F10.5	Azimuth of P axis measured clock- wise from North
	31-40	XGAM	F10.5	The angle the P axis makes with the downward vertical z-axis
C.				
	1-10	AAA(1)	F10.5	Focal depth in km
D. Data header				
	1-10	TT	F10.5	.LE.0 end program .GT.0 period in seconds
	11-15	MODAL	I5	MODAL Mode of sur- face wave data
	16-20	IDUMP	I5	Dummy variable-not used-leave blank
	21-25	IURUZ	I5	.EQ. 1 UZ for Rayl .EQ. 2 IR fpr Rayl (has meaning only for Rayleigh wave data).

D. Data header (cont'd)

Card Sequence	Column	Name	Format	Explanation
	30-40	ATTEN	E11.4	Anelastic attenuation coefficient for surface wave in units of km ⁻¹
		SRC	E11.4	Source spectrum at period TT in units of dyne-cm
E.	2-5	II	I5	Coordinates of radiation pattern positioning
F.	12-15	JJ	I5	" "
	IFORM - Format read in from first data card above	ISTA		Station code
		AZ		Azimuth of station from source measured clockwise from North in degrees
		DIST		Distance from source to station in kilometers
		WEIGHT		Weighting factor on data, not used by this program
		AMPL		Spectral amplitude observed at the station at period TT in units of cm-sec
		GPV		Group velocity of this arrival at period TT-- informative only-- not used by this program.

The program returns to Point F to read in more ISTA, AZ, DIST, WEIGHT, AMPL, GPV data until it finds one with ISTA = HALT. The HALT indicates the end of a data set. The radiation pattern is now plotted. The program then returns to Point D to read in a new TT, MODAL, IDUMP, IURUZ, ATTEN, SRC card to see if the plotting should continue or whether the program should be terminated.

***** NOTE -- The data must be arranged in order of increasing period. Also no period interpolation is performed. Hence the periods used must agree exactly with those written on the tape.

```

C PROGRAM RADPAT
C THIS PROGRAM ACCEPTS OBSERVED AMPLITUDES OF LOVE OR RAYLEIGH WAVES
C AND USING SOURCE SPECTRUM AND ATTENUATION OF A GIVEN PERIOD YIELDS
C A PLOT OF THE OBSERVED AND THEORETICAL RADIATION PATTERN.
C THE PROGRAM YIELDS AMPLITUDES AT A DISTANCE OF 1000 KILOMETERS
COMMON/ONE/AY(100,4),CY(100,4),N,DEPTH(100),D(100)
COMMON/TWO/LAYER(100),A(100),B(100),RHO(100),XMU(100),XLAMB(100)
COMMON/THREE/AMODE(10),WNO(10),ELIP(10),ABSORP(10),DENOM(10)
COMMON/FOUR/NUMDPT,AAA(15),PYY1(15,10),PYY2(15,10),PYY3(15,10),
1PYY4(15,10),IFUNC
DIMENSION ISTA(100),AZ(100),DIST(100),WEIGHT(100),AMPL(100),GPV(10
10),IBUF(1000)
DIMENSION SINAZ(361),COSAZ(361)
DIMENSION XC(100),YC(100),XT(363),YT(363),SA(361),CA(361),SD(361)
DIMENSION CD(361),SCP(361),SAP(361)
DIMENSION PYR(4)
DIMENSION IFORM(15)
CALL PLCTS(IBUF,1000,10)
CALL PLOT(0.0,-11.0,-3)
CALL PLOT(0.0,1.0,-3)
DEGRAD = 0.01745329
IHALT = 4HHALT
1 FORMAT(1H ,2(I4,6X))
4 FORMAT(8F10.5)
15 FORMAT(1H ,2(E11.4,4X,E11.4,9X))
16 FORMAT(1H ,5(E11.4,4X),2F10.1,2I5)
17 FORMAT(F10.5,3I5,4X,E11.4,9X,E11.4)
18 FORMAT(12X,10A6)
19 FORMAT(1H )
20 FORMAT(1H ,A4,5X,7HDEPTH = , F10.2)
CALL MODEL(01)
READ(60,18) (IFRM(I),I=1,10)
WRITE(61,18)(IFRM(I),I=1,10)
READ(60,4) XDEL,XBET,XALP,XGAM
WRITE(61,4) XDEL,XBET,XALP,XGAM
XDEL = DEGRAD * XDEL
XBET = DEGRAD * XBET
XGAM = DEGRAD * XGAM
XALP = DEGRAD * XALP
FACT1 = (COS(XBET)**2) + (COS(XGAM)**2)
FACT2 = (SIN(XBET))**2
FACT3 = (SIN(XGAM))**2
FACT4 = SIN(2.*XBET)
FACT5 = SIN(2.*XGAM)
DO 4004 K = 1,361
XJ = K - 1
ANG = DEGRAD * XJ
SA(K) = SIN(XALP-ANG)
CA(K) = COS(XALP-ANG)
SD(K) = SIN(XDEL-ANG)

```

```

CD(K) = COS(XDEL-ANG)
SDF(K) = SIN(2.*(XDEL-ANG))
SAP(K) = SIN(2.*(XALP-ANG))
SINAZ(K) = SIN(ANG)
4004 COSAZ(K) = COS(ANG)
NUMDPT = 1
T = 0
READ(60,4) AAA(1)
C IF WE WISH TO TAKE AVERAGE OVER FAULT PLANE USE FOLLOWING.
C SMALL EFFECT
C NUMDPT = 9
C DEEP = AAA(1)
C DO 4010 I = 1,NUMDPT
C XJ = I - 5
C THIS HAS FAULT PLANE EXTENDING FROM AAA(1) +-2.6 KM
C AAA(1) = DEEP + 0.7 * XJ
C4010 CONTINUE
2999 CONTINUE
      WRITE(61,10)
      READ(60,17) TT,MODAL,IDLMP,IURUZ,ATTEN,SRC
      IF(TT.LE.0) GO TO 9999
      WRITE(61,17) TT,MODAL,IDLMP,IURUZ,ATTEN,SRC
      READ(60,1) II,JJ
      CALL POSIT(II,JJ,XX,YY)
      DBMAX = 0.0
      DO 29 I = 1,100
      NUMSTA = I
      READ(60,IFORM) ISTA(I),AZ(I),DIST(I),WEIGHT(I),AMPL(I),GPV(I)
      IF(ISTA(I).EQ.1HALT) GO TO 291
29 CONTINUE
      GO TO 292
291 NUMSTA = NUMSTA - 1
292 CONTINUE
C      CORRECTION FOR ANELASTIC ATTENUATION
      DO 300 I = 1,NUMSTA
      PHI = DEGRAD * AZ(I)
      DUM = AMPL(I) * EXP(ATTEN*DIST(I))
      AMPL(I) = DUM
      YO(I) = DUM * COS(PHI)
      XO(I) = DUM * SIN(PHI)
      IF(DUM.GT.DBMAX) DBMAX = DUM
      WRITE(61,IFORM) ISTA(I),AZ(I),DIST(I),WEIGHT(I),AMPL(I),GPV(I)
300 CONTINUE
      IF(T.EQ.TT) GO TO 2961
      CALL EIGEN(W1,TT,T)
2961 CONTINUE
      SUM1 = 0
      SUM2 = 0
      SUM3 = 0
      SUM4 = 0

```

```

DO 3000 I = 1,NUMDPT
SUM1 = SUM1 + PYY1(I,MODAL)
SUM2 = SUM2 + PYY2(I,MODAL)
SUM3 = SUM3 + PYY3(I,MODAL)
SUM4 = SUM4 + PYY4(I,MODAL)
3000 CONTINUE
THMAX = 0.0
PYR(1) = SUM1 / NUMDPT
PYR(2) = SUM2 / NUMDPT
PYR(3) = SUM3 / NUMDPT
PYR(4) = SUM4 / NUMDPT
DO 3002 K = 1,361
SINA = SA(K)
COSA = CA(K)
SIND = SD(K)
COSD = CD(K)
SINDP = SDP(K)
SINAP = SAP(K)
GO TO (2711,2721),IFUNC
2710 CONTINUE
XREAL = 0.5 * PYR(1)*(FACT2*SINDP - FACT3*SINAP)
XIMAG = -0.5*PYR(2)*(FACT4*SIND - FACT5*SINA)
GO TO 2730
2720 CONTINUE
XREAL = FACT1*PYR(4) -PYR(1)*(FACT2*COSD*COSD - FACT3*COSA*COSA)
XIMAG = 0.5*(PYR(2)*PYR(3))*(FACT4*COSD-FACT5*COSA)
2730 CONTINUE
C THE OUTPUT OF EIGEN IS IN CGS UNITS THIS STEP IS TAKEN SO THAT
C THEORETICAL AMPLITUDES ARE NOT TOO SMALL FOR MACHINE .LT. 1E-19
XREAL = SRC * XREAL
XIMAG = SRC * XIMAG
DUM = SQRT(XREAL**2 + XIMAG**2)
IF(IFUNC.EQ.2.AND.IURUZ.GE.2) DUM = DUM * ABS(ELIP(MODAL))
XT(1) = DUM * SINAZ(K)
YT(1) = DUM * COSAZ(K)
IF(DUM.GT.THMAX) THMAX = DUM
3002 CONTINUE
C RADIATION PATTERNS ARE GIVEN WITHIN A 1 INCH RADIUS CIRCLE
VALMAX = DBMAX
IF(THMAX.GT.DBMAX) VALMAX = THMAX
DO 3500 K = 1*361
YT(K) =(YT(1) / VALMAX)+ YY
XT(K) =(XT(1) / VALMAX)+ XX
3500 CONTINUE
WRITE(61,20) AMODE(MODAL),AAA(1)
WRITE(61,15) XT(1),YT(1),XT(46),YT(46),XT(91),YT(91),XT(136),YT(13
16)
XT(362) = 0.0
YT(362) = 0.0
XT(363) = 1.0

```

```

YT(363) = 1.0
CALL LINE(XT,YT,361,1.0,0)
DO 3501 IN = 1,NUMSTA
  X=(XO(IN)/VALMAX)+ XX
  Y=(YO(IN)/VALMAX)+ YY
  CALL SYMBOL(X,Y,0.035,0.0,0,-1)
3501 CONTINUE
  Z = ALOG10(VALMAX)
  MM = Z
  FT = Z - MM
  FT = 10.*FT
  XXX = 1./FT
  IF(XXX.GT.1.) MM = MM - 1
  IF(XXX.GT.1.) XXX = XXX / 10.
  FT = MM
  WRITE(61,15) VALMAX,OBMAX,TMAX,XXX,FT,XX,YY,II,JJ
  CALL SYMBOL(XX-0.025,YY+1.0,0.10,1MM,0.0,1)
  CALL PLOT(XX,YY+0.95,3)
  CALL PLOT(XX,YY+0.82,2)
  CALL SYMBOL(XX,YY,0.03,3,0.0,-1)
  CALL PLOT(XX+XXX,YY-1.20,3)
  CALL PLOT(XX+XXX,YY-1.30,2)
  CALL PLOT(XX+XXX,YY-1.25,2)
  CALL PLOT(XX-XXX,YY-1.25,2)
  CALL PLOT(XX-XXX,YY-1.20,3)
  CALL PLOT(XX-XXX,YY-1.30,2)
  CALL SYMBOL(XX-0.35,YY-1.15,0.10,4H2 X +0.0,4)
  CALL SYMBOL(XX+0.05,YY-1.15,0.10,2H10,0.0,2)
  CALL NUMBER(XX+0.26,YY-1.07,0.07,FT,0.0,-1)
  CALL SYMBOL(XX-0.4,YY-1.50,0.10,4HT = ,0.0,4)
  CALL NUMBER(XX,YY-1.50,0.10,TT,0.0,1)
  GO TO 2999
9999 CONTINUE
REWIND 01
CALL PLOT(13.0,0.0,999)
STOP
END

SUBROUTINE POSIT(I,J,XX,YY)
DIMENSION X(3),Y(3)
X(1) = 1.5
X(2) = 4.25
X(3) = 7.00
Y(1) = 7.25
Y(2) = 4.50
Y(3) = 1.75
IF(I.GT.0.AND.J.GT.0) GO TO 100
CALL PLOT(13.0,0.0,-3)
CALL PLOT(0.0,-11.0,-3)
CALL PLOT(0.0,1.0,-3)

```

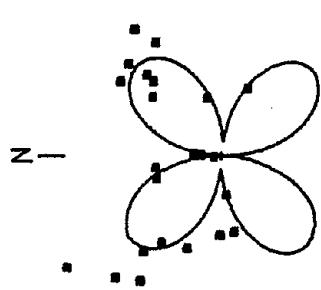
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```
I = IABS(I)
J = IABS(J)
100 CONTINUE
XX = X(J)
YY = Y(I)
RETURN
END
```

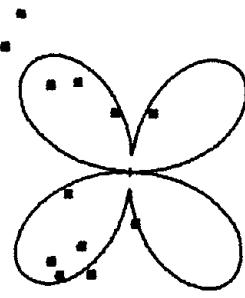
(1X,A4,5X,3E10.2,E11.4,9X,F10.4)

181,	52,	272,	89,			
12,						
16,	1	0	1	1.4000E-04	2.2918E 23	
2	1					
SLMT	3.2	335.3		8.4891E-04	3.3989	16.
LHCT	3.6	1425.30		6.8459E-04	3.4802	16.
AAMT	36.2	941.4		5.1925E-03	3.4276	16.
BLOT	40.3	526.4		4.4153E-03	3.4484	16.
OTTT	44.4	1630.		4.9422E-03	3.4836	16.
MNTT	47.2	1793.		4.5981E-03	3.4970	16.
CLET	47.4	1014.30		4.4010E-03	3.3878	16.
STJT	55.1	3374.1		4.2669E-03	3.4486	16.
WEST	59.8	1815.0		5.1962E-03	3.4810	16.
BLAT	76.0	917.1		2.3541E-03	3.5017	16.
ATLT	111.6	612.0		3.4088E-03	3.4991	16.
LUBT	261.1	1069.7		2.9095E-03	3.2037	16.
TUCT	264.6	1916.8		1.6224E-03	3.2604	16.
ALQT	271.7	1455.8		1.9990E-03	3.3494	16.
DUGT	291.2	2024.7		3.2214E-03	3.3583	16.
NEWT	303.7	2489.9		5.5236E-03	3.2151	16.
LONT	305.1	2881.8		3.3148E-03	3.3513	16.
MSOT	309.7	2316.3		5.3376E-03	3.3697	16.
PNTT	311.8	2811.9		5.7971E-03	3.3968	16.
EDMT	324.7	2651.9		7.7983E-03	3.3830	16.
FFCT	340.9	2307.3		2.4308E-03	3.4970	16.
MBCT	350.2	4758.5		1.7778E-03	3.4853	16.
HALT						
-1,						

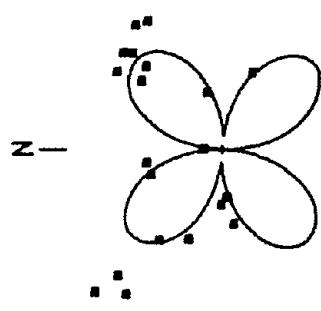
FUND	DEPTH	12.00	0.1400E-03	0.2292E 24
SLHIT	0	3.20	335.30	0. 8897E-03
LHCT	3.60	1425.30	0.	0.8358E-03
AAHT	36.120	941.40	0.	0.5924E-02
BLOIT	40.40	526.40	0.	0.4753E-02
OTTIT	44.40	1650.00	0.	0.6226E-02
MNTIT	47.20	1793.00	0.	0.5910E-02
CLET	47.40	1014.30	0.	0.5073E-02
STUT	55.10	3374.10	0.	0.6843E-02
WEST	59.80	1815.00	0.	0.6699E-02
BLATI	76.00	917.10	0.	0.2677E-02
ATLI	111.60	612.00	0.	0.3714E-02
LUBT	261.10	1069.70	0.	0.3380E-02
TUCIT	264.60	1916.80	0.	0.2122E-02
ALQT	271.70	1455.80	0.	0.2451E-02
DUGIT	291.40	2024.70	0.	0.4277E-02
NEWIT	303.70	2489.90	0.	0.7827E-02
LONT	305.10	2881.80	0.	0.4962E-02
MSOT	309.70	2316.30	0.	0.7382E-02
PNTIT	311.60	2813.90	0.	0.8596E-02
EDMT	324.70	2651.90	0.	0.1130E-01
FFCT	340.90	2307.30	0.	0.3358E-02
MBCIT	350.20	4758.50	0.	0.3461E-02
				3.3968
				3.2151
				3.3513
				3.3697
				3.3830
				3.4970
				3.4853
				3.3988
				3.4802
				3.4276
				3.4484
				3.4836
				3.4970
				3.3878
				3.4466
				3.4810
				3.5017
				3.4991
				3.2037
				3.2604
				3.3494
				3.3583
				3.3988
				3.4765E 01
				0.4765E 01
				0.1765E 01
				0.4235E 01
				0.4251E-02
				0.8846E 00
				-0.2000E 01
				1.5
				4.5
				2
				1



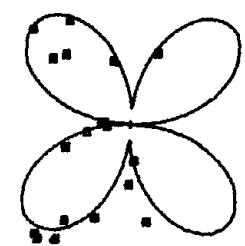
2×10^{-3}
 $T = 20.0$



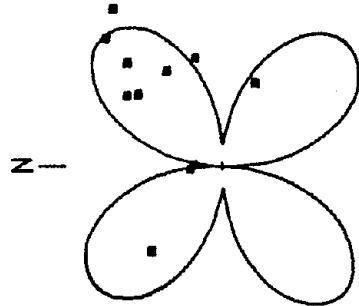
2×10^{-3}
 $T = 40.0$



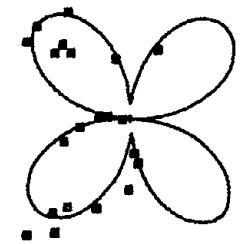
2×10^{-3}
 $T = 16.0$



2×10^{-3}
 $T = 28.0$



2×10^{-3}
 $T = 9.0$



2×10^{-3}
 $T = 24.0$

XI. QUESTION

PROGRAMMER: R. B. HERRMANN / May 1973

PURPOSE:

This program accepts observed Love and Rayleigh wave amplitude spectra for an earthquake. The program searches through a parameter space of focal mechanism orientations and focal depth. For each particular combination, the program computes several goodness of fit characteristics which are used by the user to determine the best focal mechanism and focal depth combination which satisfies the observations.

$$S(\omega) = M_o / i\omega .$$

INPUT/OUTPUT

The surface wave amplitude spectra are read from card on FILE 60. Two tapes are used. The Love wave eigenfunction tape is placed on FILE 01 and the Rayleigh wave eigenfunction tape is placed on FILE 02. Printer output is on FILE 61. Both tapes must have the same earth model specifications. The sample data for LEIGEN and REIGEN are consistent with the periods on eigenfunction tapes.

THEORY

The theory of Levshin and Yanson is used to specify the theoretical amplitude spectra. The formulation in terms of strike, dip, and slip is used. Since only the amplitudes will be considered, the strike need vary only over 180° , the dip over 90° , and the slip over 180° . Because of the symmetry of the theoretical surface wave amplitude spectra to a 180° rotation in strike, the strike is only known to within 180° . The slip varies over 180° instead of 360° since the amplitudes only are used. The choice of the proper strike (or strike + 180°) is made from the best fit to observed P wave first motion.

For each combination of strike, dip, slip, and focal depth the following parameters are computed.

MOMENT - Seismic moment from average of Love and Rayleigh wave estimates.

R-RAYL - Correlation coefficient between the observed and theoretical Rayleigh wave amplitude spectra for the totality of data from all azimuths and periods.

RES-RAYL - Sum of square residuals between observed and theoretical Rayleigh wave amplitude spectra using the average seismic moment estimate

R-LOVE - Correlation coefficient between the observed and theoretical Love wave amplitude spectra for the totality of data from all azimuths and periods

RES-LOVE - Sum of square residuals between observed and theoretical Love wave amplitude spectra using the average seismic moment estimate

RES-L+R - Square root of sum of Love and Rayleigh square residuals.

MOMENT-L - Seismic moment estimate from only Love wave data

MOMENT-R - Seismic moment estimate from only Rayleigh wave data

If either correlation coefficient is less than zero, none of the above information is listed. This saves considerable printer time during the program execution.

The last column of the printout is flagged by a ** if the two independent seismic moment estimates are within a multiplicative factor of 1.66 .

After the computations for a given focal depth, a summary is given of the various slip, dip, and strike combinations which gave the largest RR, RL, and least RES, RESR and RESL, as well as the largest value of BEST. BEST is defined as $RR \cdot RL \cdot (\text{ratio of seismic moment estimates chosen to be less than 1.0})$.

The best focal mechanism estimate has been found to be the one with the largest values of the correlation coefficients RR, RL and for which the two independent seismic moment estimates are as equal as possible. The output of all these parameters permits a robust search.

When beginning a search for the best focal mechanism, the following search range is used:

DIP 30° (15°) 90°

SLIP 90° (20°) 70°

STRIKE 0° (20°) 160°

DEPTH 5 (5) 20 km

The spectra of $\text{SLIP} = -90^\circ$ equals the spectra for $\text{SLIP} = +90^\circ$. Likewise $\text{STRIKE} = 0^\circ$ and 180° yield the same spectra. Hence the choice above.

After performing this initial search over a wide grid, an appropriate region is found for the solution and the program is rerun using a finer grid.

PROGRAM DESCRIPTION

PROGRAM QUESTION : This is the main input/output routine.

SUBROUTINE PRESS: This subroutine determines the correlation coefficients, residuals, and seismic moments.

(The following subroutines are required for operation but are found in the User Library. They are described elsewhere in this volume).

SUBROUTINE EIGEN

SUBROUTINE MODEL

SUBROUTINE NONPRP

SUBROUTINE EXCIT

REFERENCES

Herrmann, R. B. (1974). Surface-wave generation by central United States earthquakes, Ph. D. Dissertation, Saint Louis University.

Card Sequence	Column	Name	Format	Explanation
A.	1-10	DISTMX	F10.5	Data for distances greater than DISTMX km are not used in the computations
B.	2-5	NUMDPT	I4	Number of focal depths to be considered
C. Focal depth cards (NUMDPT cards)				
	1-10	AAA(I)	F10.5	Focal depths in km
D.	1-5	ND	I5	Number of dips to be used
	11-20	DI	F10.5	First dip in degrees
	21-30	DINC	F10.5	Dip increment in degrees
E.	1-5	NS	I5	Number of slips to be used
	11-20	SI	F10.5	Initial slip angle
	21-30	SINC	F10.5	Slip angle increment
F.	1-5	NF	I5	Number of strikes to be used
	11-20	FI	F10.5	Initial strike angle
	21-30	FINC	F10.5	Strike angle increment
(Love wave data first, arranged in order of increasing period.)				
G.	1-10	TT	F10.5	.LT.0 end of Love wave data .GT.0 period in seconds
	11-15	MODAL	I5	Mode of Love wave data at period TT
	16-20	NUMSTA	I5	Dummy variable-not used

G. (cont'd)

Card

Sequence	Column	Name	Format	Explanation
	21-25	IURUZ	I5	Not used for Love waves
	30-40	ATTEN	E11.4	Anelastic attenuation coefficient in km^{-1} for Love wave of period TT and mode MODAL
H.	2-5	ISTA	A4	Station identifier code
	11-20	AZ	F10.5	Azimuth from source to station in degrees, clockwise from North
	21-30	DIST	F10.5	Distance from source to station in km
	31-40	DUM1	F10.5	Dummy variable - no use in this program
	41-51	AMP	E11.4	Spectral amplitude
	61-70	DUM2	F10.5	DUM2 variable - not used by this program

The program corrects the amplitude for anelastic attenuation and returns to Point H to read another ISTA set. When it finds ISTA = HALT, the program then goes to Point G to read another TT header and corresponding ISTA data. All the Love wave data is read in until a TT = -1.0 (negative) is found. (This signifies the end of the Love wave data).

(The Rayleigh wave data is now read, arranged in order of increasing period.)

Card Sequence	Column	Name	Format	Explanation
I.	1-10	TT	F10.5	.LT.0 end of Rayleigh wave data .GT.0 period in seconds
	11-15	MODAL	I5	Mode of Rayleigh wave data Fundamental-- MODAL = 1
	16-20	NUMSTA	I5	Dummy variable-not used
	21-25	IURUZ	I5	.EQ.1 Vertical component spectra used .EQ.2 Radial component. Rayleigh wave spectra used
	30-40	ATTEN	E11.4	Anelastic attenuation coefficient in km^{-1} for Rayleigh waves of period TT and mode MODAL
J.	2-5	ISTA	A4	Station identification code
	11-20	AZ	F10.5	Azimuth from source to station in degrees, measured clockwise, from N
	21-30	DIST	F10.5	Distance from source to station in km
	31-40	DUM1	F10.5	Dummy variable-not used by program
	41-51	AMP	E11.4	Spectral amplitude in cm-sec
	61-70	DUM2	F10.5	Dummy variable - not used by program

The program corrects the spectral amplitude for anelastic attenuation and returns to Point J to read another ISTA set. When it finds an ISTA = HALT, the program recognizes the end of data for period TT and mode MODAL and proceeds to POINT I to read another TT header and the corresponding ISTA data. A value of TT = negative indicates the end of Rayleigh wave data. At this point the search process begins for the best solution.

C PROGRAM QUESTION

C THIS PROGRAM ACCEPTS RAYLEIGH AND LOVE SPECTRAL VALUES FOR VARIOUS
 C PERIODS AND MODES AND COMPUTES SEISMIC MOMENT ASSUMING
 C 1./OMEGA SOURCE. THE GOODNESS OF FIT AND THE MOMENT ARE
 C DETERMINED FOR ANY FOCAL MECHANISM

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COMMON/ONE/AY(100,4),CY(100,4),N,DEPTH(100),D(100)
COMMON/TWO/LAYER(100),A(100),B(100),RHO(100),XMU(100),XLAMB(100)
COMMON/THREE/AMODE(10),WVNO(10),ELIP(10),ABSORP(10),DENOM(10)
COMMON/FOUR/NUMDPT,AAA(15),PYY1(15,10),PYY2(15,10),PYY3(15,10),
1PYY4(15,10),IFUNC
DIMENSION STRK(15),CST(15),SST(15),S2ST(15),C2ST(15)
DIMENSION PIP(15),SD(15),CD(15),C2D(15),SLIP(15),SS(15),CS(15)
DIMENSION SINAZ(2,550),COSAZ(2,550),XL(550),XR(550),YL(550),YR(550)
1.),PYL1(10,30),PYL2(10,30),PYR1(10,30),PYR2(10,30),PYR3(10,30)
DIMENSION PYR4(10,30),NNUML(30),NNUMR(30)
DIMENSION IS(6),ID(6),IP(6),IBEST(6)
DIMENSION XBEST(6)
M1 = 4H
M2 = 4H **
IBEST(1) = 4H RR
IBEST(2) = 4H RL
IBEST(3) = 4H RES
IBEST(4) = 4HRESR
IBEST(5) = 4HRESL
IBEST(6) = 4HBEST
1 FORMAT(1H ,2(I4,6X))
2 FORMAT(10X,5H NR = ,I5,10X,5H NL = , I5 )
3 FORMAT(1H1,8HDISTMX = ,F10.2/)
4 FORMAT(8F11.5)
7 FORMAT(F10.5,3I5,4X,E11.4)
10 FORMAT(1X,A4,5X,3F10.5,E11.4,9X,F10.5)
11 FORMAT(I5,5X,2F+0.5)
12 FORMAT(1H0,1X,5HDEPTH,2X,4HSLIP,3X,3HDIP,6H STRK,2X,8H MOMENT ,
1 4X,8H R-RAYL,4X,8HRES-RAYL,4X,8H R-LOVE,4X,8HRES-LOVE,4X,
2 8HRES--L+R,4X,8HMOMENT-L,4X,8HMOMENT-R )
13 FORMAT(1H ,3H***,A4,3(F5.0,1X),1X,E10.3)
14 FORMAT(1H ,4F6.1,8(1X,E10.3,1X),4X,A4)
C LOGIC UNIT 01 = LOVE TAPE
C LOGIC UNIT 02 = RAYL TAPE
CALL MODEL(01)
CALL MODEL(02)
READ(60,4) DISTMX
WRITE(61,3) DISTMX
READ(60,1) NUMDPT
DO 28 I = 1,NUMDPT
28 READ(60,4) AAA(I)
IHALT = 4HHALT
XLOW = 1.0E-30
DEGRAD = 0.017452329
C ND = NUMBER DIPS      DI = INITIAL DIP      DINC = DIP INCREMENT
```

QUESTION PAGE 2

```
      READ(60,11) ND, DI, DINC
C      NS = NUMBER SLIPS   SI = INITIAL SLIP   SINC = SLIP INCREMENT
      READ(60,11) NS,SI, SINC
C      NF = NUMBER STRIKES FI = INITIAL STRIKE FINC = STRIKE INCREMENT
      READ(60,11) NF, FI, FINC
      DO 4001 J = 1,NF
      XJ = J
      DIP(J) = (DI + (XJ-1.)*DINC)
      RDIP = DEGRAD * DIP(J)
      SD(J) = SIN(RDIP)
      CD(J) = COS(RDIP)
      C2D(J) = COS(2.*RDIP)
 4001 CONTINUE
      DO 4002 K = 1,NS
      XJ = K
      SLIP(K) = (SI + (XJ-1.) * SINC)
      RSLIP = DEGRAD * SLIP(K)
      SS(K) = SIN(RSLIP)
      CS(K) = COS(RSLIP)
 4002 CONTINUE
      DO 4003 L = 1,NF
      XJ = L
      STRK(L) = (FI + (XJ-1.) * FINC)
      RSTRK = DEGRAD * STRK(L)
      CST(L) = COS(RSTRK)
      SST(L) = SIN(RSTRK)
      S2ST(L) = SIN(2.*RSTRK)
      C2ST(L) = COS(2.*RSTRK)
 4003 CONTINUE
C      SOURCE SPECTRUM IS ASSUMED TO BE 1./OMEGA
C      TT = PERIOD OF DATA
C      MODAL = MODE OF DATA  1 = FUND, 2 = 1ST, ETC
C      NUMSTA = DUMMY
C      IURUZ  1 = UZ DATA , 2 = UR DATA
C      ATTEN = APPROXIMATE VALUE OF ATTENUATION COEFFICIENT GAMMA
      T = 0
      NL = 0
C      LOVE WAVE DATA
      DO 293 J = 1,30
      READ(60,7) TT,MODAL,NUMSTA,IURUZ,ATTEN
      IF(TT,LT,0) GO TO 7997
      NUMSTA = 0
      DO 29 I = 1, 100
      READ(60,10) ISTA,AZ,DIST,DUM1,AMP,DUM2
      IF(DIST,GE,DISTMX) GO TO 29
      NUMSTA = NUMSTA + 1
      IF(ISTA,EQ,1HALT) GO TO 291
      NL = NL + 1
      IF(NL,GT,550) GO TO 9999
      XL(NL)=AMP*EXP(ATTEN*DIST)
```

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AZ = DEGRAD * AZ
SINAZ(1,NL) = SIN(AZ)
COSAZ(1,NL) = COS(AZ)
29 CONTINUE
GO TO 292
291 NUMSTA = NUMSTA - 1
292 CONTINUE
NNUML(J) = NUMSTA
FACT = TT/6.2831853
IF(T.EQ.TT) GO TO 296
CALL EIGEN(Q1,TT,T)
296 DO 297 IZ = 1,NUMDPT
PYL1(IZ,J) = PYY1(IZ,MODAL) * FACT
297 PYL2(IZ,J) = PYY2(IZ,MODAL) * FACT
293 CONTINUE
GO TO 7987
7997 J = J - 1
7987 JLMX = J
T = 0
NR = 0
C RAYLEIGH WAVE DATA
DO 313 J = 1,30
READ(60,7) TT,MODAL,NUMSTA,IURUZ,ATTEN
IF(TT.LT.0) GO TO 8998
NUMSTA = 0
DO 31 I = 1,100
READ(60,10) ISTA,AZ,DIST,DUM1,AMP,DUM2
IF(DIST.GE.DISTMX) GO TO 31
NUMSTA = NUMSTA + 1
IF(ISTA.EQ.1HALT) GO TO 311
NR = NR + 1
IF(NR.GT.550) GO TO 9999
XR(NR)=AMP*EXP(ATTEN*DIST)
AZ = DEGRAD * AZ
SINAZ(2,NR) = SIN(AZ)
COSAZ(2,NR) = COS(AZ)
31 CONTINUE
GO TO 312
311 NUMSTA = NUMSTA - 1
312 CONTINUE
NNUMR(J) = NUMSTA
FACT = TT/6.2831853
IF(T.EQ.TT) GO TO 315
CALL EIGEN(Q2,TT,T)
315 CONTINUE
IF(IURUZ.EQ.2) FACT = FACT + ABS(ELIP(MODAL))
316 DO 317 IZ = 1,NUMDPT
PYR1(IZ,J) = PYY1(IZ,MODAL) * FACT
PYR2(IZ,J) = PYY2(IZ,MODAL) * FACT
PYR3(IZ,J) = PYY3(IZ,MODAL) * FACT

```

QUESTION PAGE 4

```

317 PYR4(IZ,J) = PYY4(IZ,MODAL) * FACT
313 CONTINUE
      GO TO 8988
8998 J = J - 1
8988 JRMX = J
      WRITE(61,2) NR,NL
      DO 3000 I = 1,NUMDPT
      WRITE(61,12)
      SRESR=1.0/XLOW
      SRESL=1.0/XLOW
      SRES =1.0/XLOW
      GRR=0.0
      GRL=0.0
      PMRR = 0.0
      DO 3001 J = 1,ND
      SINL = SD(J)
      COSL = CD(J)
      COS2D = C2D(J)
      DO 3002 K = 1,NS
      SINS = SS(K)
      COSS = CS(K)
      DO 3003 L = 1,NF
      COSST = CST(L)
      SINST = SST(L)
      SIN2ST = S2ST(L)
      COS2ST = C2ST(L)
      FAC1 = SIND*COSS*COS2ST
      FAC2 = SIND*COSS*SIN2ST
      FAC3 = -SIND*COSD*SINS*COS2ST
      FAC4 = -SIND*COSD*SINS*SIN2ST
      FAC5 = COS2D*SINS*CCSST
      FAC6 = COS2D*SINS*SINST
      FAC7 = -COSD*COSS*CCSST
      FAC8 = -COSD*COSS*SINST
      FAC9 = 2.*SIND*COSD*SINS
      ML = 0
      DO 3008 KL = 1,JLMX
C       THE OUTPUT OF EIGEN IS IN CGS UNITS THIS STEP IS TAKEN SO THAT
C       THEORETICAL AMPLITUDES ARE NOT TOO SMALL FOR MACHINE ,LT. 1E-19
      PY1 = PYL1(I,KL)*1.0E20
      PY2 = PYL2(I,KL)*1.0E20
      NUM = NNUML(KL)
      DO 3006 M = 1,NUM
      ML = ML + 1
      SAZ = SINAZ(1,ML)
      CAZ = COSAZ(1,ML)
      CALL EXCIT(FAC1,FAC2,FAC3,FAC4,FAC5,FAC6,FAC7,FAC8,FAC9,
      1 PY1,PY2,PY3,PY4,1,SAZ,CAZ,AMP,1)
3006 YL(ML) = AMP
3008 CONTINUE

```

QUESTION PAGE 5

```

MR = 0
DO 3009 KR = 1,JRMX
C THE OUTPUT OF EIGEN IS IN CGS UNITS THIS STEP IS TAKEN SO THAT
C THEORETICAL AMPLITUDES ARE NOT TOO SMALL FOR MACHINE ,LT. 1E-19
PY1 = PYR1(I,KR)*1.0E20
PY2 = PYR2(I,KR)*1.0E20
PY3 = PYR3(I,KR)*1.0E20
PY4 = PYR4(I,KR)*1.0E20
NUM = NNUMR(KR)
DO 3007 M = 1,NUM
MR = MR + 1
SAZ = SINAZ(2,MR)
CAZ = COSAZ(2,MR)
CALL EXCIT(FAC1,FAC2,FAC3,FAC4,FAC5,FAC6,FAC7,FAC8,FAC9,
1 PY1,PY2,PY3,PY4,2,SAZ,CAZ,AMP,1)
3007 YR(MR) = AMP
3009 CONTINUE
CALL RRESS(AL,NR,XL,XR,YL,YR,RL,RR,S,RESL,RESR,RES,SL,SR)
C THE OUTPUT OF EIGEN IS IN CGS UNITS THIS STEP IS TAKEN SO THAT
C THEORETICAL AMPLITUDES ARE NOT TOO SMALL FOR MACHINE ,LT. 1E-19
S = S * 1.0E20
SL = SL * 1.0E20
SR = SR*1.0E20
ICROAK = M1
IF(RL,LT,0,CR,RR,LT,0) GO TO 421
SLDSR = SL/SR
IF(SLDSR.GT.0.6.AND.SLDSR.LT.1.7) ICROAK = M2
WRITE(61,14) AAA(I),SLIP(K),DIP(J),STRK(L),S,RR,RESR,RL,RESL,RES
1 ,SL,SR,ICROAK
IF(SLDSR.GT.1) SLDSR = 1./SLDSR
PROD = RR * RL * SLDSR
IF(RR+LE+GRR) GO TO 402
GRR=RR
IS(1)=K
ID(1)=J
IP(1)=L
402 IF(RL,LE,GRL) GO TO 403
GRL=RL
IS(2)=K
ID(2)=J
IP(2)=L
403 IF(RES.GE.SRES) GO TO 404
SRES=RES
IS(3)=K
ID(3)=J
IP(3)=L
404 IF(RESR.GE.SRESR) GO TO 405
SRESR=RESR
IS(4)=K
ID(4)=J

```

```

IP(4)=L
405 IF(PRESL.GE.SRESL) GO TO 410
SRESL=RESL
IS(5)=K
ID(5)=J
IP(5)=L
410 IF(PROD.LE.PMRR) GO TO 411
PMRR=PROD
IS(6)=K
ID(6)=J
IP(6)=L
411 CONTINUE
421 CONTINUE
3003 CONTINUE
3002 CONTINUE
3001 CONTINUE
XBEST(1) = GRR
XBEST(2) = GRL
XBEST(3) = SRES
XBEST(4) = SRESR
XBEST(5) = SRESL
XBEST(6) = PMRR
DO 420 NN = 1,6
KK=IS(NN)
JJ=ID(NN)
LL=IP(NN)
420 WRITE(61,13) IREST(NN),SLIP(KK),DIP(JJ),STRK(LL),XBEST(NN)
3000 CONTINUE
9999 CONTINUE
WRITE(61,2) NR,AL
REWIND 1
REWIND 2
END

SUBROUTINE RRESS(NL,NR,XL,XR,YL,YR,RL,RR,S,RESL,RESR,RES,SL,SR)
DIMENSION XL(1),XR(1),YL(1),YR(1)
SUML0 = 0
SUML1 = 0
SUML2 = 0
SUML3 = 0
SUML4 = 0
SUML5 = 0
SUMR0 = 0
SUMR1 = 0
SUMR2 = 0
SUMR3 = 0
SUMR4 = 0
SUMR5 = 0
DO 601 I = 1,NL
SUML0 = SUML0 + 1.

```

```
SUML1 = SUML1 + XL(I)
SUML2 = SUML2 + XL(I) * XL(I)
SUML3 = SUML3 + YL(I)
SUML4 = SUML4 + YL(I) * YL(I)
601 SUML5 = SUML5 + XL(I) * YL(I)
DO 602 I = 1, NR
SUMR0 = SUMR0 + 1,
SUMR1 = SUMR1 + XR(I)
SUMR2 = SUMR2 + XR(I) * XR(I)
SUMR3 = SUMR3 + YR(I)
SUMR4 = SUMR4 + YR(I) * YR(I)
602 SUMR5 = SUMR5 + XR(I) * YR(I)
RL = (SUML0*SUML5-SUML1*SUML3)/ SQRT((SUML0*SUML2-SUML1*SUML1) +
1 (SUML0*SUML4-SUML3*SUML3))
RR = (SUMR0*SUMR5-SUMR1*SUMR3)/SQRT((SUMR0*SUMR2-SUMR1*SUMR1) +
1 (SUMR0*SUMR4-SUMR3*SUMR3))
SL = SUML5/SUML4
SR = SUMR5/SUMR4
S = (SUML0 * SL + SUMR0 * SR)/(SUML0 + SUMR0)
RESL = SUML2 - 2.*S*SUML5 + S*S*SUML4
RESR = SUMR2 - 2.*S*SUMR5 + S*S*SUMR4
RES = SQRT(ABS(RESL + RESR))
RETURN
END
```

3500,					
3					
5.					
10,					
15.					
3	45.	15.			
6	-90.	30.			
6	0.0	30.			
6.0	1	0 1	6.0000E-04	0.8594E 23	
BLAT	76.0	917.1		2.7991E-03	3.4329
MBCT	350.2	4758.5		4.8066E-04	3.4202
HALT					
7.0	1	0 1	5.0000E-04	1.0027E 23	
AAMT	36.2	941.4		3.4337E-03	3.3160
BLOT	40.3	526.4		2.2932E-03	3.1493
MNTT	47.2	1793.		8.9092E-04	3.5107
CLET	47.4	1014.30		2.5317E-03	3.3267
WEST	59.8	1815.0		1.3002E-03	3.3263
BLAT	76.0	917.1		2.3604E-03	3.4752
MBCT	350.2	4758.5		9.1541E-04	3.4288
HALT					
8.0	1	0 1	4.0000E-04	1.1459E 23	
AAMT	36.2	941.4		3.0654E-03	3.2870
BLOT	40.3	526.4		2.6165E-03	3.2615
MNTT	47.2	1793.		1.8601E-03	3.4800
CLET	47.4	1014.30		3.7731E-03	3.3321
WEST	59.8	1815.0		1.8485E-03	3.3680
BLAT	76.0	917.1		2.3649E-03	3.4458
ATLT	111.6	612.0		1.6026E-03	3.5968
MBCT	350.2	4758.5		1.5448E-03	3.4338
FCCT	355.2	2591.60		4.6747E-04	3.4472
HALT					
9.0	1	0 1	3.0000E-04	1.2892E 23	
SLMT	3.2	335.3		9.0595E-04	3.2824
AAMT	36.2	941.4		3.0412E-03	3.3368
BLOT	40.3	526.4		3.2108E-03	3.4093
MNTT	47.2	1793.		2.7919E-03	3.4834
CLET	47.4	1014.30		4.3481E-03	3.0390
STJT	55.1	3374.1		2.3679E-03	3.3829
WEST	59.8	1815.0		2.1790E-03	3.3947
BLAT	76.0	917.1		2.8796E-03	3.4169
ATLT	111.6	612.0		2.8249E-03	3.5343
MSOT	309.7	2316.3		1.8667E-03	3.2355
MBCT	350.2	4758.5		2.0619E-03	3.4425
FCCT	355.2	2591.60		5.1259E-04	3.4472
HALT					
10.0	1	0 1	2.0000E-04	1.4324E 23	
SLMT	3.2	335.3		9.9339E-04	3.3480
CHCT	3.6	1425.30		6.4216E-04	3.4528
AAMT	36.2	941.4		3.4497E-03	3.3937

BLOT	40,3	526,4	3.6821E-03	3.4260	10.
OTTT	44,4	1650,	2.1226E-03	3.4473	10.
MNTT	47,2	1793,	3.4781E-03	3.4868	10.
CLET	47,4	1014,30	4.7529E-03	3.3486	10.
STJT	55,1	3374,1	3.8204E-03	3.3931	10.
WEST	59,8	1815,0	2.9500E-03	3.4300	10.
BLAT	76,0	917,1	2.8878E-03	3.4297	10.
ATLT	111,6	612,0	3.1447E-03	3.5396	10.
LONT	305,1	2881,8	1.1006E-03	3.2420	10.
MSOT	309,7	2316,3	3.2504E-03	3.2698	10.
PNTT	311,8	2813,9	1.6218E-03	3.3011	10.
EDMT	324,7	2651,9	6.9380E-03	3.2265	10.
MBCT	350,2	4758,5	2.5056E-03	3.4462	10.
FCCT	355,2	2591,60	6.2815E-04	3.4244	10.
HALT					
12,0	1	0 1	1.8000E-06	1.7189E 23	
SLMT	3,2	335,3	8.7832E-04	3.4162	12.
LHCT	3,6	1425,30	6.7490E-04	3.5058	12.
AAMT	36,2	941,4	4.6993E-03	3.4152	12.
BLOT	40,3	526,4	4.7110E-03	3.3874	12.
OTTT	44,4	1650,	3.5122E-03	3.4437	12.
CLET	47,4	1014,30	4.8991E-03	3.3821	12.
MNTT	47,2	1793,	4.5247E-03	3.4800	12.
STJT	55,1	3374,1	4.6678E-03	3.4051	12.
WEST	59,8	1815,0	4.7602E-03	3.4512	12.
BLAT	76,0	917,1	2.3993E-03	3.4426	12.
ATLT	111,6	612,0	3.3696E-03	3.5092	12.
LUBT	261,1	1069,7	3.5037E-03	3.0814	12.
ALQT	271,7	1455,8	1.6857E-03	3.3038	12.
DUGT	291,2	2024,7	2.2743E-03	3.2426	12.
NEWT	303,7	2489,9	3.3736E-03	3.1914	12.
LONT	305,1	2881,8	2.7395E-03	3.2791	12.
MSOT	309,7	2316,3	3.9906E-03	3.3071	12.
PNTT	311,8	2813,9	2.1071E-03	3.4174	12.
EDMT	324,7	2651,9	7.7152E-03	3.3050	12.
FFCT	340,9	2307,3	7.0549E-04	3.5590	12.
MBCT	350,2	4758,5	2.4532E-03	3.4663	12.
FCCT	355,2	2591,60	6.5769E-04	3.4438	12.
HALT					
14,	1	0 1	1.6000E-04	2.0054E 23	
SLMT	3,2	335,3	6.2858E-04	3.4603	14.
LHCT	3,6	1425,30	8.4011E-04	3.4217	14.
AAMT	36,2	941,4	5.1475E-03	3.4339	14.
BLOT	40,3	526,4	4.8098E-03	3.4204	14.
OTTT	44,4	1650,	4.2765E-03	3.4727	14.
MNTT	47,2	1793,	4.8216E-03	3.4868	14.
CLET	47,4	1014,30	4.8599E-03	3.3934	14.
STJT	55,1	3374,1	4.5359E-03	3.4224	14.
WEST	59,8	1815,0	5.1680E-03	3.4644	14.
BLAT	76,0	917,1	2.2099E-03	3.4297	14.

ATLT	111.6	612.0	3.1774E-03	3.4892	14.
LUBT	261.1	1069.7	3.4543E-03	3.1287	14.
TUCT	264.6	1916.8	2.0117E-03	3.2384	14.
ALQT	271.7	1455.8	1.2076E-03	3.3513	14.
DUGT	291.2	2024.7	3.0225E-03	3.3143	14.
NEWT	303.7	2489.9	4.9861E-03	3.2213	14.
LONT	305.1	2881.8	3.4397E-03	3.3185	14.
MSOT	309.7	2316.3	4.7667E-03	3.3599	14.
PNTT	311.8	2813.9	4.2536E-03	3.3906	14.
EDMT	324.7	2651.9	7.8316E-03	3.3615	14.
FFCT	340.9	2307.3	1.7451E-03	3.4838	14.
MBCT	350.2	4758.5	2.0222E-03	3.4789	14.
FCCT	355.2	2591.60	4.7894E-04	3.4575	14.
HALT					
16.	1	0 1	1.4000E-04	2.2918E 23	
SLMT	3.2	335.3	8.4891E-04	3.3989	16.
LHCT	3.6	1425.30	6.8459E-04	3.4802	16.
AAMT	36.2	941.4	5.1925E-03	3.4276	16.
BLOT	40.3	526.4	4.4153E-03	3.4484	16.
OTTT	44.4	1650.	4.9422E-03	3.4836	16.
MNTT	47.2	1793.	4.5981E-03	3.4970	16.
CLET	47.4	1014.30	4.4010E-03	3.3878	16.
STJT	55.1	3374.1	4.2669E-03	3.4486	16.
WEST	59.8	1815.0	5.1962E-03	3.4810	16.
BLAT	76.0	917.1	2.3541E-03	3.5017	16.
ATLT	111.6	612.0	3.4088E-03	3.4991	16.
LUBT	261.1	1069.7	2.9095E-03	3.2037	16.
TUCT	264.6	1916.8	1.6224E-03	3.2604	16.
ALQT	271.7	1455.8	1.9990E-03	3.3494	16.
DUGT	291.2	2024.7	3.2214E-03	3.3583	16.
NEWT	303.7	2489.9	5.5236E-03	3.2151	16.
LONT	305.1	2881.8	3.3148E-03	3.3513	16.
MSOT	309.7	2316.3	5.3376E-03	3.3697	16.
PNTT	311.8	2813.9	5.7971E-03	3.3968	16.
EDMT	324.7	2651.9	7.7983E-03	3.3830	16.
FFCT	340.9	2307.3	2.4308E-03	3.4970	16.
MBCT	350.2	4758.5	1.7778E-03	3.4853	16.
HALT					
18.	1	0 1	1.2000E-04	2.5783E 23	
SLMT	3.2	335.3	9.0823E-04	3.3564	18.
LHCT	3.6	1425.30	8.1978E-04	3.4908	18.
AAMT	36.2	941.4	5.2242E-03	3.4623	18.
BLOT	40.3	526.4	3.9577E-03	3.4826	18.
OTTT	44.4	1650.	4.9760E-03	3.4984	18.
MNTT	47.2	1793.	4.2965E-03	3.5004	18.
CLET	47.4	1014.30	3.9117E-03	3.4163	18.
STJT	55.1	3374.1	4.5630E-03	3.4645	18.
WEST	59.8	1815.0	5.1315E-03	3.4927	18.
BLAT	76.0	917.1	2.4508E-03	3.5253	18.
ATLT	111.6	612.0	3.2956E-03	3.5345	18.

LUBT	261.1	1069.7		2.7728E-03	3.1941	18.
TUCT	264.6	1916.8		1.4905E-03	3.2480	18.
ALQT	271.7	1455.8		2.7499E-03	3.3075	18.
DUGT	291.2	2024.7		3.3274E-03	3.3807	18.
NEWT	303.7	2489.9		5.5151E-03	3.2370	18.
LONT	305.1	2881.8		3.3145E-03	3.3670	18.
MSOT	309.7	2316.3		4.9811E-03	3.3820	18.
PNTT	311.8	2813.9		5.5888E-03	3.4174	18.
EDMT	324.7	2651.9		7.1691E-03	3.4096	18.
FFCT	340.9	2307.3		2.4461E-03	3.5076	18.
MBCT	350.2	4758.5		1.8372E-03	3.4802	18.
HALT						
20.	1	0	1	1.0000E-04	2.8648E 23	
SLMT	3.2	335.3		9.1387E-04	3.4603	20.
LHCT	3.6	1425.30		1.0844E-03	3.4444	20.
AAAMT	36.2	941.4		5.0996E-03	3.4655	20.
BLOT	40.3	526.4		3.8283E-03	3.4769	20.
OTTT	44.4	1650.		4.9274E-03	3.5021	20.
MNTT	47.2	1793.		4.0911E-03	3.5072	20.
CLET	47.4	1014.30		4.0533E-03	3.4395	20.
STJT	55.1	3374.1		4.8742E-03	3.4770	20.
WEST	59.8	1815.0		4.8487E-03	3.5198	20.
BLAT	76.0	917.1		2.4357E-03	3.4984	20.
ATLT	111.6	612.0		3.0189E-03	3.5602	20.
LUBT	261.1	1069.7		3.0454E-03	3.1917	20.
TUCT	264.6	1916.8		1.4179E-03	3.2425	20.
ALQT	271.7	1455.8		3.0171E-03	3.3094	20.
DUGT	291.2	2024.7		3.5706E-03	3.3920	20.
NEWT	303.7	2489.9		5.1862E-03	3.2592	20.
LONT	305.1	2881.8		3.5135E-03	3.3788	20.
MSOT	309.7	2316.3		4.3631E-03	3.4043	20.
PNTT	311.8	2813.9		5.4352E-03	3.4257	20.
EDMT	324.7	2651.9		6.5454E-03	3.4245	20.
FFCT	340.9	2307.3		2.4413E-03	3.4996	20.
MBCT	350.2	4758.5		1.8751E-03	3.4904	20.
FCCT	355.2	2591.60		2.6374E-04	3.3416	20.
HALT						
22.	1	0	1	0.8000E-04	3.1513E 23	
SLMT	3.2	335.3		9.0117E-04	3.4964	22.
LHCT	3.6	1425.30		1.1741E-03	3.4074	22.
AAAMT	36.2	941.4		4.5990E-03	3.4370	22.
BLOT	40.3	526.4		3.8602E-03	3.4093	22.
MNTT	47.2	1793.		3.8788E-03	3.5107	22.
CLET	47.4	1014.30		3.8511E-03	3.4221	22.
STJT	55.1	3374.1		5.2060E-03	3.4860	22.
WEST	59.8	1815.0		4.4698E-03	3.5387	22.
BLAT	76.0	917.1		2.3462E-03	3.4818	22.
ATLT	111.6	612.0		2.8630E-03	3.5706	22.
LUBT	261.1	1069.7		2.3160E-03	3.1681	22.
TUCT	264.6	1916.8		1.2224E-03	3.3168	22.

ALOT	271.7	1455.8	2.8285E-03	3.3417	22.
DUGT	291.2	2024.7	3.5712E-03	3.4092	22.
NEWT	303.7	2489.9	4.7496E-03	3.2753	22.
LONT	305.1	2881.8	3.5530E-03	3.4108	22.
MSOT	309.7	2316.3	4.0858E-03	3.4245	22.
PNTT	311.8	2813.9	5.3097E-03	3.4404	22.
EDMT	324.7	2651.9	5.9631E-03	3.4355	22.
FFCT	340.9	2307.3	2.4430E-03	3.5049	22.
MBCT	350.2	4758.5	1.6628E-03	3.5110	22.
FCCT	355.2	2591.60	2.5214E-04	3.3666	22.
HALT					
24.	1	0 1	0.8000E-04	3.4377E 23	
SLMT	3.2	335.3	8.9556E-04	3.5519	24.
LHCT	3.6	1425.30	1.1334E-03	3.4114	24.
AAMT	36.2	941.4	4.8287E-03	3.4433	24.
BLOT	40.3	526.4	3.8452E-03	3.3874	24.
OTTT	44.4	1650.	4.6535E-03	3.5321	24.
MNTT	47.2	1793.	3.4910E-03	3.5559	24.
CLET	47.4	1014.30	3.2497E-03	3.4163	24.
STJT	55.1	3374.1	5.8823E-03	3.4770	24.
WEST	59.8	1815.0	4.2839E-03	3.5422	24.
BLAT	76.0	917.1	2.2883E-03	3.5017	24.
ATLT	111.6	612.0	2.8225E-03	3.5758	24.
LUBT	261.1	1069.7	1.6599E-03	3.1241	24.
TUCT	264.6	1916.8	1.1884E-03	3.4327	24.
ALOT	271.7	1455.8	2.5445E-03	3.3666	24.
DUGT	291.2	2024.7	3.3038E-03	3.4440	24.
NEWT	303.7	2489.9	4.5242E-03	3.2861	24.
LONT	305.1	2881.8	3.4819E-03	3.4455	24.
MSOT	309.7	2316.3	4.1010E-03	3.4474	24.
PNTT	311.8	2813.9	5.0413E-03	3.4594	24.
EDMT	324.7	2651.9	5.1177E-03	3.4490	24.
FFCT	340.9	2307.3	2.3631E-03	3.5076	24.
MBCT	350.2	4758.5	1.4003E-03	3.5280	24.
FCCT	355.2	2591.60	2.5572E-04	3.2730	24.
HALT					
26.	1	0 1	0.8000E-04	3.7242E 23	
SLMT	3.2	335.3	9.1227E-04	3.6288	26.
LHCT	3.6	1425.30	9.6413E-04	3.4951	26.
AAMT	36.2	941.4	5.3947E-03	3.6014	26.
BLOT	40.3	526.4	3.6288E-03	3.4149	26.
OTTT	44.4	1650.	4.4517E-03	3.5473	26.
MNTT	47.2	1793.	3.1529E-03	3.6060	26.
CLET	47.4	1014.30	3.0750E-03	3.5602	26.
STJT	55.1	3374.1	6.0512E-03	3.4842	26.
WEST	59.8	1815.0	4.9293E-03	3.5613	26.
BLAT	76.0	917.1	2.2262E-03	3.5423	26.
ATLT	111.6	612.0	2.1307E-03	3.5758	26.
LUBT	261.1	1069.7	1.9948E-03	3.1751	26.
TUCT	264.6	1916.8	1.1783E-03	3.4889	26.

ALQT	271.7	1455.8	2.2749E+03	3.3824	26.
DUGT	291.2	2024.7	3.0513E-03	3.4675	26.
NEWT	303.7	2489.9	4.3957E-03	3.3025	26.
LONT	305.1	2881.8	3.4277E-03	3.4729	26.
MSOT	309.7	2316.3	4.2757E-03	3.4654	26.
PNTT	311.8	2813.9	4.6407E-03	3.4765	26.
EDMT	324.7	2651.9	5.2280E-03	3.4715	26.
FFCT	340.9	2307.3	2.2307E-03	3.5183	26.
MBCT	350.2	4758.5	1.2157E-03	3.5398	26.
HALT					
28.	1	0 1	0.8000E-04	4.0107E 23	
SLMT	3.2	335.3	1.0200E-03	3.4692	26.
LHCT	3.6	1425.30	8.3968E-04	3.5606	26.
AAMT	36.2	941.4	6.0470E-03	3.8362	26.
BLOT	40.3	526.4	3.4863E-03	3.4428	26.
OTTT	44.4	1650.	4.2973E-03	3.5664	26.
MNTT	47.2	1793.	2.9387E-03	3.6169	26.
STJT	55.1	3374.1	5.2375E-03	3.5427	26.
WEST	59.8	1815.0	3.7179E-03	3.6055	26.
BLAT	76.0	917.1	2.1691E-03	3.5944	26.
ATLT	111.6	612.0	2.6015E-03	3.6074	26.
LUBT	261.1	1069.7	3.2476E-03	3.3154	26.
TUCT	264.6	1916.8	1.1281E-03	3.4984	26.
ALQT	271.7	1455.8	1.9144E-03	3.3610	26.
DUGT	291.2	2024.7	3.0604E-03	3.4675	26.
NEWT	303.7	2489.9	4.8622E-03	3.3312	26.
LONT	305.1	2881.8	3.3213E-03	3.4978	26.
MSOT	309.7	2316.3	4.4054E-03	3.4863	26.
PNTT	311.8	2813.9	4.1707E-03	3.5025	26.
EDMT	324.7	2651.9	5.4612E-03	3.4967	26.
FFCT	340.9	2307.3	2.0361E-03	3.5535	26.
MBCT	350.2	4758.5	1.0819E-03	3.5623	26.
FCCT	355.2	2591.60	6.8163E-04	3.4541	26.
HALT					
30.	1	0 1	0.8000E-04	4.2972E 23	
LHCT	3.6	1425.30	8.0098E-04	3.5606	30.
AAMT	36.2	941.4	6.4815E-03	3.8998	30.
BLOT	40.3	526.4	3.3159E-03	3.4484	30.
OTTT	44.4	1650.	4.2371E-03	3.5858	30.
MNTT	47.2	1793.	2.8852E-03	3.6023	30.
STJT	55.1	3374.1	4.3225E-03	3.6071	30.
WEST	59.8	1815.0	3.4065E-03	3.6435	30.
BLAT	76.0	917.1	2.0655E-03	3.6335	30.
ATLT	111.6	612.0	2.4118E-03	3.6779	30.
LUBT	261.1	1069.7	4.5324E-03	3.3283	30.
TUCT	264.6	1916.8	1.9736E-03	3.4905	30.
DUGT	291.2	2024.7	3.2708E-03	3.4469	30.
NEWT	303.7	2489.9	3.6209E-03	3.3661	30.
LONT	305.1	2881.8	3.2187E-03	3.5170	30.
MSOT	309.7	2316.3	4.3618E-03	3.5021	30.

PNTT	311.8	2813.9	3.8042E+03	3.5444	30.
EDMT	324.7	2651.9	4.7444E+03	3.4990	30.
FFCT	340.9	2307.3	1.8617E+03	3.6063	30.
MBCT	350.2	4758.9	9.2581E+04	3.5959	30.
HALT					
35,	1	0 1	0.6000E+04	5.0134E 23	
LHCT	3.6	1425.30	9.1639E+04	3.6471	35.
AAAMT	36.2	941.4	7.0043E+03	4.1444	35.
BLOT	40.3	526.4	2.2832E+03	3.6902	35.
OTTT	44.4	1650.	4.2554E+03	3.7191	35.
MNTT	47.2	1793.	2.8678E+03	3.6132	35.
STJT	55.1	3374.1	3.7171E+03	3.7021	35.
WEST	59.8	1815.0	2.8168E+03	3.6860	35.
BLAT	76.0	917.1	1.7757E+03	3.7525	35.
ATLT	111.6	612.0	1.9882E+03	3.9006	35.
LUBT	261.1	1069.7	4.6987E+03	3.6029	35.
TUCT	264.6	1916.8	1.0858E+03	3.5986	35.
DUGT	291.2	2024.7	2.7073E+03	3.5969	35.
NEWT	303.7	2489.9	2.9423E+03	3.3638	35.
LONT	305.1	2881.8	2.9365E+03	3.5386	35.
MSQT	309.7	2316.3	3.1867E+03	3.4346	35.
PNTT	311.8	2813.9	3.2986E+03	3.6525	35.
EDMT	324.7	2651.9	2.5629E+03	3.5222	35.
FFCT	340.9	2307.3	1.7760E+03	3.7317	35.
HALT					
40,	1	0 1	0.6000E+04	5.7296E 23	
OTTT	44.4	1650.	3.9975E+03	3.6777	40.
MNTT	47.2	1793.	2.6279E+03	3.7027	40.
STJT	55.1	3374.1	3.9163E+03	3.8065	40.
WEST	59.8	1815.0	2.3149E+03	3.7919	40.
BLAT	76.0	917.1	1.4325E+03	3.7991	40.
ATLT	111.6	612.0	1.5001E+03	4.0157	40.
TUCT	264.6	1916.8	1.1493E+03	3.5290	40.
DUGT	291.2	2024.7	2.4258E+03	3.6033	40.
NEWT	303.7	2489.9	1.9112E+03	3.3672	40.
LONT	305.1	2881.8	2.6314E+03	3.5604	40.
PNTT	311.8	2813.9	2.5364E+03	3.7201	40.
FFCT	340.9	2307.3	1.4313E+03	3.7621	40.
HALT					
45,	1	0 1	0.6000E+04	6.4458E 23	
OTTT	44.4	1650.	3.9628E+03	3.7066	45.
MNTT	47.2	1793.	2.1168E+03	3.8008	45.
STJT	55.1	3374.1	3.5916E+03	3.8966	45.
WEST	59.8	1815.0	1.8029E+03	3.8749	45.
ATLT	111.6	612.0	1.1483E+03	4.3281	45.
LONT	305.1	2881.8	2.2098E+03	3.7165	45.
PNTT	311.8	2813.9	1.7006E+03	3.8952	45.
FFCT	340.9	2307.3	1.2329E+03	3.7837	45.
HALT					
50,	1	0 1	0.6000E+04	7.1620E 23	

MNTT	47.2	1793.		1.6703E+03	3.6837	50.
STJT	55.1	3374.1		3.2756E+03	4.0389	50.
PNTT	311.8	2813.9		1.5283E+03	3.8605	50.
FFCT	340.9	2307.3		1.2582E+03	3.8988	50.
HALT						
-1,						
5.0	1	0 1	7.0000E-04	0.7162E 23		
SLMZ	3.2	335.3		1.0161E+03	3.0099	5.
ATLZ	111.6	611.2		1.4279E+03	2.8878	5.
ROLZ	334.2	284.2		1.0276E+03	3.0299	5.
HALT						
6.0	1	0 1	6.0000E-04	0.8594E 23		
SLMZ	3.2	335.3		2.2808E+03	3.0166	6.
BLOZ	40.3	526.4		1.5474E+03	2.8585	6.
ATLZ	111.6	611.20		2.3801E+03	2.9829	6.
EDMZ	324.7	2651.9		0.3150E+04	2.8230	6.
ROLZ	334.2	284.2		1.4814E+03	2.9744	6.
HALT						
7.0	1	0 1	5.0000E-04	1.0027E 23		
SLMZ	3.2	335.3		3.9518E+03	3.0166	7.
LHCZ	3.6	1425.30		1.0498E+03	3.0847	7.
AAMZ	36.2	941.4		2.5370E+03	2.9590	7.
BLOZ	40.3	526.4		2.4028E+03	2.8979	7.
WES1	59.8	1815.0		1.0318E+03	2.8929	7.
BLAZ	76.0	917.1		1.3893E+03	3.0504	7.
ATLZ	111.6	611.2		2.1269E+03	3.0235	7.
MSOZ	309.7	2316.3		1.2152E+03	2.9231	7.
EDMZ	324.7	2651.9		1.3487E+03	2.8320	7.
ROLZ	334.2	284.2		2.3577E+03	3.0875	7.
MBCZ	350.2	4758.5		1.0460E+03	3.0724	7.
FCCZ	355.2	2591.60		7.7257E+04	3.1497	7.
HALT						
8.0	1	0 1	4.0000E-04	1.1459E 23		
SLMZ	3.2	335.3		4.5646E+03	3.0861	8.
LHCZ	3.6	1425.30		1.5264E+03	3.0998	8.
AAMZ	36.2	941.4		3.2005E+03	3.0014	8.
BLOZ	40.3	526.4		2.5102E+03	2.9631	8.
OTTZ	44.4	1650.9		1.3664E+03	3.1096	8.
MNTZ	47.2	1793.6		9.9463E+04	3.0074	8.
STJZ	55.1	3374.1		7.2545E+04	3.0527	8.
WES1	59.8	1815.0		1.3461E+03	2.9894	8.
BLAZ	76.0	917.1		1.6706E+03	3.0889	8.
ATLZ	111.6	611.2		3.1515E+03	3.0768	8.
MSOZ	309.7	2316.3		1.6655E+03	2.9213	8.
EDMZ	324.7	2651.9		1.4166E+03	2.8672	8.
ROLZ	334.2	284.2		2.4695E+03	3.1128	8.
MBCZ	350.2	4758.5		1.9471E+03	3.0813	8.
FCCZ	355.2	2591.60		1.5107E+03	3.1345	8.
HALT						
9.0	1	0 1	3.0000E-04	1.2092E 23		

SLMZ	3.2	335.3	4.3259E-03	3.1075	9.
LHCZ	3.6	1425.30	2.0393E-03	3.1066	9.
AAMZ	36.2	941.4	3.8213E-03	3.0353	9.
BLOZ	40.3	526.4	2.6020E-03	3.0270	9.
OTTZ	44.4	1650,	1.0704E-03	3.1482	9.
MNTZ	47.2	1793,	1.0225E-03	3.1253	9.
CLER	47.4	1014.30	1.0802E-03	2.9929	9.
STJZ	55.1	3374.1	1.6571E-03	3.0496	9.
WES1	59.8	1815.0	1.6408E-03	3.0205	9.
BLAZ	76.0	917.1	1.8356E-03	3.1418	9.
ATLZ	111.6	611.2	3.2410E-03	3.1319	9.
LUBZ	261.1	1069.7	2.0061E-03	2.8667	9.
TUCZ	264.6	1916.8	1.3115E-03	2.8069	9.
ALQZ	271.7	1455.8	8.1349E-04	3.0085	9.
MSOZ	309.7	2316.3	2.0064E-03	2.9398	9.
EDMZ	324.7	2651.9	1.6904E-03	2.9097	9.
ROLZ	334.2	284.2	2.1941E-03	3.0543	9.
FFCZ	340.9	2307.3	6.8574E-04	3.1508	9.
MBCZ	350.2	4758.5	2.9412E-03	3.0954	9.
FCCZ	355.2	2591.60	1.9887E-03	3.1392	9.
HALT					
10.0	1	0 1	2.0000E-04	1.4324E 23	
SLMZ	3.2	335.3	3.5582E-03	3.0932	10.
LHCZ	3.6	1425.30	2.1355E-03	3.1066	10.
AAMZ	36.2	941.4	2.8828E-03	3.0775	10.
BLOZ	40.3	526.4	2.3525E-03	3.0270	10.
OTTZ	44.4	1650,	1.3934E-03	3.1422	10.
MNTZ	47.2	1793,	1.1206E-03	3.1473	10.
CLER	47.4	1014.30	1.4342E-03	2.9929	10.
STJZ	55.1	3374.1	2.2226E-03	3.0538	10.
WES1	59.8	1815.0	1.8628E-03	3.0458	10.
BLAZ	76.0	917.1	1.9044E-03	3.1418	10.
ATLZ	111.6	611.2	2.9250E-03	3.1360	10.
LUBZ	261.1	1069.7	2.5717E-03	2.8708	10.
TUCZ	264.6	1916.8	1.1871E-03	2.8048	10.
ALQZ	271.7	1455.8	1.1165E-03	3.0131	10.
LONZ	305.1	2881.8	1.2889E-03	2.9712	10.
MSOZ	309.7	2316.3	2.0530E-03	2.9662	10.
EDMZ	324.7	2651.9	1.3119E-03	2.9338	10.
ROLZ	334.2	284.2	2.1992E-03	3.0218	10.
FFCZ	340.9	2307.3	8.8492E-04	3.1572	10.
FCCZ	355.2	2591.60	1.8552E-03	3.1488	10.
HALT					
12.0	1	0 1	1.8000E-04	1.7189E 23	
SLMZ	3.2	335.3	2.8389E-03	3.1220	12.
LHCZ	3.6	1425.30	1.9254E-03	3.1168	12.
AAMZ	36.2	941.4	2.2893E-03	3.0699	12.
BLOZ	40.3	526.4	1.7311E-03	3.0193	12.
OTTZ	44.4	1650,	1.9867E-03	3.1038	12.
MNTZ	47.2	1793,	1.4572E-03	3.1280	12.

CLER	47.4	1014.30	1.3160E-03	2.9973	12.
STJZ	55.1	3374.1	2.5086E-03	3.0565	12.
WES1	59.8	1815.0	2.3676E-03	3.0561	12.
BLAZ	76.0	917.1	2.2352E-03	3.0811	12.
ATLZ	111.6	611.2	2.4221E-03	3.1808	12.
TUCZ	264.6	1916.8	9.1576E-04	2.9642	12.
ALQZ	271.7	1455.8	2.1401E-03	2.9401	12.
DUGZ	291.2	2024.7	2.1988E-03	2.9736	12.
LONZ	305.1	2881.8	8.0948E-04	3.0420	12.
MSOZ	309.7	2316.3	2.0729E-03	3.0027	12.
EDMZ	324.7	2651.9	1.9007E-03	3.0036	12.
ROLZ	334.2	284.2	2.3067E-03	3.1214	12.
FFCZ	340.9	2307.3	1.2413E-03	3.1572	12.
MBCZ	350.2	4758.5	3.4356E-03	3.0863	12.
FCCZ	355.2	2591.60	1.1791E-03	3.1430	12.
HALT					
14.	1	0 1	1.6000E-04	2.0054E 23	
SLMZ	3.2	335.3	2.6050E-03	3.1513	14.
LHCZ	3.6	1425.30	1.8692E-03	3.1100	14.
AAMZ	36.2	941.4	1.9791E-03	3.0329	14.
BLOZ	40.3	526.4	1.6319E-03	3.0667	14.
OTTZ	44.4	1650,	1.8574E-03	3.0720	14.
MNTZ	47.2	1793,	1.4416E-03	3.1226	14.
CLER	47.4	1014.30	7.0880E-04	3.0196	14.
STJZ	55.1	3374.1	2.1843E-03	3.0469	14.
WES1	59.8	1815.0	2.0088E-03	3.0433	14.
BLAZ	76.0	917.1	1.9824E-03	3.0811	14.
ATLZ	111.6	611.2	1.4747E-03	3.0729	14.
LUBZ	261.1	1069.7	5.4938E-04	3.1402	14.
TUCZ	264.6	1916.8	7.1416E-04	3.0298	14.
ALQZ	271.7	1455.8	2.3197E-03	2.9416	14.
DUGZ	291.2	2024.7	2.7392E-03	2.9692	14.
LONZ	305.1	2881.8	5.4073E-04	3.0163	14.
MSOZ	309.7	2316.3	2.0042E-03	2.9969	14.
EDMZ	324.7	2651.9	2.0220E-03	2.9918	14.
ROLZ	334.2	284.2	2.1265E-03	3.0875	14.
FFCZ	340.9	2307.3	1.3113E-03	3.0979	14.
MBCZ	350.2	4758.5	2.8563E-03	3.0615	14.
HALT					
16.	1	0 1	1.4000E-04	2.2918E 23	
SLMZ	3.2	335.3	2.0586E-03	3.1662	16.
LHCZ	3.6	1425.30	1.8304E-03	3.0583	16.
AAMZ	36.2	941.4	2.1194E-03	3.0600	16.
BLOZ	40.3	526.4	1.5042E-03	3.0097	16.
OTTZ	44.4	1650,	1.6590E-03	3.0692	16.
MNTZ	47.2	1793,	1.1696E-03	3.0535	16.
CLER	47.4	1014.30	9.6655E-04	3.0514	16.
STJZ	55.1	3374.1	1.2497E-03	3.0277	16.
WES1	59.8	1815.0	1.7155E-03	3.0612	16.
BLAZ	76.0	917.1	1.5719E-03	3.0811	16.

ATLZ	111.6	611.2	1.1358E-03	3.1603	16.
LUBZ	261.1	1069.7	4.4781E-04	3.2400	16.
TUCZ	264.6	1916.8	4.7992E-04	3.0467	16.
ALQZ	271.7	1455.8	1.7883E-03	2.9446	16.
DUGZ	291.2	2024.7	2.0440E-03	2.9390	16.
LONZ	305.1	2881.8	7.9251E-04	2.9774	16.
MSOZ	309.7	2316.3	1.8302E-03	2.9892	16.
PNTZ	311.8	2813.9	4.6592E-04	2.9514	16.
EDMZ	324.7	2651.9	2.1366E-03	3.0002	16.
ROLZ	334.2	284.2	1.8834E-03	3.0379	16.
FFCZ	340.9	2307.3	1.4200E-03	3.0752	16.
HALT					
18,	1	0	1.2000E-04	2.5783E 23	
SLMZ	3.2	335.3	1.3914E-03	3.0440	18.
LHCZ	3.6	1425.30	1.4748E-03	3.0681	18.
AAMZ	36.2	941.4	1.9852E-03	3.0901	18.
BLOZ	40.3	526.4	1.4003E-03	3.0622	18.
OTTZ	44.4	1650.	1.5572E-03	3.0579	18.
MNTZ	47.2	1793.	1.3002E-03	3.0587	18.
CLER	47.4	1014.30	1.9039E-03	3.1707	18.
STJZ	55.1	3374.1	1.5582E-03	3.0250	18.
WES1	59.8	1815.0	1.2590E-03	3.0625	18.
BLAZ	76.0	917.1	1.1238E-03	3.0428	18.
ATLZ	111.6	611.2	9.5085E-04	3.2399	18.
TUCZ	264.6	1916.8	2.8667E-04	3.0576	18.
ALQZ	271.7	1455.8	1.2206E-03	2.9268	18.
DUGZ	291.2	2024.7	1.2671E-03	2.9137	18.
MSOZ	309.7	2316.3	1.7215E-03	2.9795	18.
PNTZ	311.8	2813.9	6.7301E-04	2.9654	18.
EDMZ	324.7	2651.9	2.2259E-03	3.0070	18.
ROLZ	334.2	284.2	1.6847E-03	3.1915	18.
FFCZ	340.9	2307.3	1.4597E-03	3.0528	18.
MBCZ	350.2	4758.5	2.0057E-03	3.0265	18.
FCCZ	355.2	2591.60	1.3393E-03	3.0713	18.
HALT					
20,	1	0	1.0000E-04	2.8648E 23	
SLMZ	3.2	335.3	1.3198E-03	2.9831	20.
LHCZ	3.6	1425.30	1.3614E-03	3.0599	20.
AAMZ	36.2	941.4	1.2744E-03	3.0926	20.
BLOZ	40.3	526.4	1.1678E-03	3.0445	20.
OTTZ	44.4	1650.	1.2368E-03	3.0353	20.
MNTZ	47.2	1793.	1.5036E-03	3.0639	20.
STJZ	55.1	3374.1	1.2832E-03	3.0662	20.
WES1	59.8	1815.0	9.8851E-04	3.0017	20.
BLAZ	76.0	917.1	9.1581E-04	2.9593	20.
ATLZ	111.6	611.2	6.5929E-04	3.3281	20.
TUCZ	264.6	1916.8	2.3236E-04	3.0430	20.
ALQZ	271.7	1455.8	8.4961E-04	2.9006	20.
DUGZ	291.2	2024.7	7.2643E-04	2.9221	20.
MSOZ	309.7	2316.3	1.5100E-03	2.9950	20.

PNTZ	311.8	2813.9	7.9996E+04	2.9811	20
EDMZ	324.7	2651.9	1.9193E+03	2.9867	20
ROLZ	334.2	284.2	1.4237E+03	3.2186	20
FFCZ	340.9	2307.3	1.1418E+03	3.0568	20
MBCZ	350.2	4758.5	1.5975E+03	3.0429	20
FCCZ	355.2	2591.60	1.3285E+03	3.0677	20
HALT					
22,	1	0 1	0.8000E+04	3.1513E 23	
SLMZ	3.2	335.3	1.1943E+03	3.4425	22
LHCZ	3.6	1425.30	1.4126E+03	3.0948	22
AAMZ	36.2	941.4	1.5003E+03	3.0575	22
BLOZ	40.3	526.4	8.9760E+04	2.9631	22
OTTZ	44.4	1650.	1.0495E+03	3.0465	22
MNTZ	47.2	1793.	1.4961E+03	3.0850	22
STJZ	55.1	3374.1	2.1183E+03	3.1058	22
WES1	59.8	1815.0	7.9566E+04	3.0306	22
BLAZ	76.0	917.1	8.6200E+04	2.9475	22
TUCZ	264.6	1916.8	3.6588E+04	3.0262	22
ALQZ	271.7	1455.8	6.0200E+04	2.8991	22
MSQZ	309.7	2316.3	1.3027E+03	3.0263	22
PNTZ	311.8	2813.9	8.2263E+04	3.0018	22
EDMZ	324.7	2651.9	1.5745E+03	2.9750	22
ROLZ	334.2	284.2	1.1667E+03	3.1128	22
FFCZ	340.9	2307.3	9.9161E+04	3.0793	22
MBCZ	350.2	4758.5	1.3398E+03	3.0773	22
FCCZ	355.2	2591.60	1.2300E+03	3.0943	22
HALT					
24,	1	0 1	0.8000E+04	3.4377E 23	
LHCZ	3.6	1425.30	1.4842E+03	3.1617	24
AAMZ	36.2	941.4	1.2350E+03	3.0476	24
OTTZ	44.4	1650.	9.6488E+04	3.1155	24
MNTZ	47.2	1793.	1.8752E+03	3.1390	24
STJZ	55.1	3374.1	1.2616E+03	3.2031	24
WES1	59.8	1815.0	8.2929E+04	3.1792	24
BLAZ	76.0	917.1	9.3695E+04	3.0227	24
LUBZ	261.1	1069.7	4.0991E+04	3.1611	24
TUCZ	264.6	1916.8	6.1767E+04	3.0576	24
ALQZ	271.7	1455.8	4.7773E+04	2.9298	24
MSQZ	309.7	2316.3	1.2286E+03	3.1095	24
PNTZ	311.8	2813.9	7.0261E+04	3.0342	24
EDMZ	324.7	2651.9	1.2166E+03	3.0644	24
ROLZ	334.2	284.2	1.0908E+03	2.9744	24
FFCZ	340.9	2307.3	1.0533E+03	3.1551	24
MBCZ	350.2	4758.5	9.2179E+04	3.1372	24
FCCZ	355.2	2591.60	1.2148E+03	3.1555	24
HALT					
26,	1	0 1	0.8000E+04	3.7242E 23	
LHCZ	3.6	1425.30	1.4377E+03	3.2062	26
AAMZ	36.2	941.4	9.1922E+04	3.1628	26
OTTZ	44.4	1650.	8.6854E+04	3.1663	26

MNTZ	47.2	1793.	1.6563E-03	3.1835
STJZ	55.1	3374.1	1.2198E-03	3.2889
WES1	59.8	1815.0	8.5719E-04	3.2504
BLAZ	76.0	917.1	1.0799E-03	3.0915
ATLZ	111.6	611.2	5.1008E-04	3.6403
LUBZ	261.1	1069.7	5.2211E-04	3.2013
TUCZ	264.6	1916.8	7.3327E-04	3.1021
ALQZ	271.7	1455.8	4.3560E-04	3.0320
DUGZ	291.2	2024.7	6.2091E-04	3.1178
MSQZ	309.7	2316.3	1.1163E-03	3.2064
PNTZ	311.8	2813.9	5.1329E-04	3.0606
EDMZ	324.7	2651.9	1.1945E-03	3.1706
ROLZ	334.2	284.2	1.2399E-03	2.9900
FFCZ	340.9	2307.3	1.0958E-03	3.2144
MBCZ	350.2	4758.5	8.6732E-04	3.2200
FCCZ	355.2	2591.60	1.2869E-03	3.2172
HALT				
28,	1	0 1	0.8000E-04	4.0107E 23
LHCZ	3.6	1425.30	1.3009E-03	3.2170
OTTZ	44.4	1650.	7.9199E-04	3.2125
MNTZ	47.2	1793.	1.5658E-03	3.2381
STJZ	55.1	3374.1	7.9106E-04	3.2970
WES1	59.8	1815.0	7.9624E-04	3.2650
BLAZ	76.0	917.1	1.2243E-03	3.1338
ATLZ	111.6	611.2	7.9483E-04	3.4212
LUBZ	261.1	1069.7	6.9404E-04	3.2949
TUCZ	264.6	1916.8	6.3387E-04	3.1992
ALQZ	271.7	1455.8	4.8914E-04	3.1897
MSQZ	309.7	2316.3	8.3720E-04	3.2514
PNTZ	311.8	2813.9	3.9470E-04	3.1234
EDMZ	324.7	2651.9	1.1529E-03	3.1974
ROLZ	334.2	284.2	1.4247E-03	3.0138
FFCZ	340.9	2307.3	1.0975E-03	3.2807
MBCZ	350.2	4758.5	7.6243E-04	3.2630
FCCZ	355.2	2591.60	1.3687E-03	3.3022
HALT				
30,	1	0 1	0.8000E-04	4.2972E 23
LHCZ	3.6	1425.30	1.1474E-03	3.2044
OTTZ	44.4	1650.	7.5288E-04	3.2441
MNTZ	47.2	1793.	1.2955E-03	3.3129
WES1	59.8	1815.0	7.1008E-04	3.2547
BLAZ	76.0	917.1	1.2940E-03	3.1526
ATLZ	111.6	611.2	9.8719E-04	3.3509
LUBZ	261.1	1069.7	8.4603E-04	3.3439
TUCZ	264.6	1916.8	6.0935E-04	3.3284
ALQZ	271.7	1455.8	5.8920E-04	3.2722
DUGZ	291.2	2024.7	6.4633E-04	3.1496
MSQZ	309.7	2316.3	5.9334E-04	3.2560
EDMZ	324.7	2651.9	1.0920E-03	3.1820
ROLZ	334.2	284.2	1.5339E-03	2.9209

FFCZ	340.9	2307.3	1.0937E-03	3.3400
MBCZ	350.2	4758.9	5.9855E-04	3.3153
FCCZ	355.2	2591.60	1.4307E-03	3.3798
HALT				
35,	1	0 1	0.6000E-04	5.0134E-23
LHCZ	3.6	1425.30	1.0099E-03	3.3751
OTTZ	44.4	1650.	6.1376E-04	3.1602
WES1	59.8	1815.0	6.5948E-04	3.4252
BLAZ	76.0	917.1	1.1449E-03	3.3120
ATLZ	111.6	611.2	9.4786E-04	3.2834
FFCZ	340.9	2307.3	8.8108E-04	3.4218
HALT				
40,	1	0 1	0.6000E-04	5.7296E-23
OTTZ	44.4	1650.	3.8902E-04	3.2377
BLAZ	76.0	917.1	9.1741E-04	3.5491
ATLZ	111.6	611.2	9.9368E-04	3.4164
DUGZ	291.2	2024.7	2.2214E-03	3.1082
EDMZ	324.7	2651.9	5.1449E-04	3.5890
FFCZ	340.9	2307.3	4.8586E-04	3.7775
HALT				
45,	1	0 1	0.6000E-04	6.4458E-23
OTTZ	44.4	1650.	2.5450E-04	3.4545
BLAZ	76.0	917.1	8.0422E-04	3.4523
ATLZ	111.6	611.2	8.7968E-04	3.6621
DUGZ	291.2	2024.7	2.0624E-03	3.2426
EDMZ	324.7	2651.9	6.6968E-04	3.6991
FFCZ	340.9	2307.3	4.1238E-04	3.8150
HALT				
-1,				



XII. WIGGLE

PROGRAMMER: R. B. HERRMANN / May 73

PURPOSE:

This program uses the eigenfunction tapes generated by the programs LEIGEN or REIGEN to generate synthetic seismograms and also the spectra of the ground motion. If the LEIGEN tape is used, transverse ground motion is generated. If the REIGEN tape is used, the vertical or radial components of ground motion can be generated for earthquakes or an explosion. The spectra can be written on tape and are normalized for geometrical spreading to a distance of 1000 km. This tape can then be read by the programs DATAPLT or FILTER and plotted. Card output of the time series in a format for direct input to EXSPEC is also available.

This program accepts an arbitrary source spectrum and medium attenuation model. There is an option to generate source spectrum by inputting M_o , $S(\omega) = M_o/i\omega$.

Spectra and/or seismograms of the following can be generated: ground velocity, ground acceleration, or displacement; velocity or acceleration traces through a 15-100, 30-100 WWSSN systems, or two LP LRSM systems.

Note: the computed ground displacements are positive for UZ when motion is up, UR when motion is away from the source, and UT when motion is clockwise from source. The instrument magnification routines are such that a positive ground impulse causes an initial upward motion of the instrument output.

THEORY

See Herrmann (1974) or Herrmann and Nuttli (1975a,b).

INPUT/OUTPUT

Card input is on FILE 60. Printer output is on FILE 61 and 42. CALCOMP plotting is on FILE 10. Punched cards are on FILE 43. Eigenfunctions are read from FILE 01. Spectra may be written on FILE 03.

PROGRAM DESCRIPTION

PROGRAM WIGGLE : This is the main routine.

SUBROUTINE EIGENF : This reads the eigenfunction tape and

yields the eigenfunctions for the proper depth.

SUBROUTINE TAPEWR : This writes the ground spectra on FILE 03 in a format compatible with DATAPLT and FILTER.

SUBROUTINE SEISMAG1 : Response of the 15-100 or 30-100 WWSSN long period seismograph system.

SUBROUTINE SEISM3 : Response of LRSM LP system with filter 6824-2.

SUBROUTINE SEISM7 : Response of LRSM LP system with filter 6823-13.

SUBROUTINE RESP : This subroutine corrects the spectra of the ground motion for the instrument response.

SUBROUTINE VELRES : This multiplies the spectra by iw, which is the frequency domain equivalent of time domain differentiation.

SUBROUTINE FOUR2 : This is the fast Fourier transform routine.

SUBROUTINE EXCITT : This accounts for the force system at the source, either double-couple or explosive.

SUBROUTINE SRCIN : This reads in the real and imaginary components of the source spectrum or generates real and imaginary parts by reading in seismic moment at $S(\omega) = M_o / i\omega$.

SUBROUTINE GAMMA : This reads in the medium attenuation coefficients.

SUBROUTINE INTER : This establishes the complex spectra of the ground motion by linear interpolation.

SUBROUTINE PNCH : Punches time history of seismic trace on FILE 43 in a format for direct input to EXSPEC. A copy of output is placed on FILE 42.

REFERENCES

Herrmann, R. B. (1974). Surface-wave generation by central United States earthquakes, Ph.D. Dissertation, Saint Louis University.

Herrmann, R. B. and O. W. Nuttli (1975a). Ground motion modeling at regional distances for earthquakes in a continental interior, I. Theory and observations, Int. J. Earthq. Engr. and Struc.Dyn. 4, 49-58.

Herrmann, R. B. and O. W. Nuttli (1975b). Ground motion modeling at regional distances for earthquakes in a continental interior, II. Effect of focal depth, azimuth and attenuation, Int. J. Earthq. Engr. and Struc. Dyn., 4, 59-72.

INPUT DATA

Card Sequence	Column	Name	Format	Explanation
A.	1-5	IEQEX	I5	.EQ.1 Earthquake .EQ.2 Explosion
	6-10	JSRC	I5	.LE.0 read in source spectrum .GT.0 read in seismic moment
	11-20	DT	F10.5	Digitizing interval of seismogram
B.	1-10	DIP	F10.3	Dip of one nodal plane measured from horizontal toward downward pointing positive z-axis
	11-20	SLIP	F10.3	Direction of slip vector on nodal plane measured counterclockwise from the horizontal pointing in strike direction.
	21-30	STRK	F10.3	Strike of fault plane measured clockwise from north
C.	1-10	AA	F10.3	.GE.0 focal depth in km .LE.-10 go to Point B. .GT.-10 .and. LT.0 end program
D. Source spectrum data				
1. If JSRC.LE.0 read source spectrum (up to 49 periods)				
	1-10	T	F10.2	Period of source spectrum value

Card Sequence	Column	Name	Format	Explanation
D. (cont'd)	11-21	XR	E11.4	Real part of source spectrum
	31-41	XI	E11.4	Imaginary part of source spectrum
(The source spectrum values are read in until a card is reached for which T.LE 0)				
2. If JSRC ,GT.0 read seismic moment (one card)				
	1-10	XMOM	E10.3	Seismic moment in dyne-cm
E. Anelastic attenuation coefficients (up to 41 sets, 2 cards per set) (NOTE: two data cards are required per period since 11 sets of values are placed in 8F10.5 format)				
	1-10	T(I)	F10.5	Period in seconds .LE.0 end of gamma values
	11-20	ATEN(I,J), J = 1, 10		attenuation values in km^{-1} . J = 1
	...			corresponds to fundamental mode.
	71-80			
	1-10			
	11-20			
	21-30			
F.	1-10	PHI	F10.5	Azimuth of receiver from source .LT.-400 go to Point C for AA
	12-15	STA	A4	Station identification
G.	1-10	R	F10.3	.GT.0 epicentral distance in km .LE.0 go to Point F for new PHI
	11-20	T0	F10.3	Time seismogram begins
	21-25	ID(1)	I5	Ground Disp LE 0 skip LE 2 seismogram GE 2 spectra

Card Sequence	Column	Name	Format	Explanation
G. (cont'd)	26-30	ID(2)	I5	Ground Vel LE 0 skip LE 2 seismogram GE 2 spectra
	31-35	ID(3)	I5	Ground Accl LE 0 skip LE 2 seismogram GE 2 spectra
	36-40	ID(4)	I5	WWSSN Disp LE 0 skip LE 2 seismogram GE 2 spectra
	41-45	ID(5)	I5	WWSSN Vel LE 0 skip LE 2 seismogram GE 2 spectra
	46-50	ID(6)	I5	WWSSN Accl LE 0 skip LE 2 seismogram GE 2 spectra
	51-55	IURUZ	I5	.EQ.1 Vertical component of Rayleigh wave .EQ.2 Radial component of Rayleigh wave
	56-60	NPY	I5	Number of time steps to plot on seismograms. NPY.LE.1020. If NPY.LE.0, NPY = 1020
	61-70	YAXLEN	F10.3	Length of time axis in inches LE 9. If YAXLEN .LE.0, YAXLEN = 5.12

Card Sequence	Column	Name	Format	Explanation
G. (cont'd)	71-75	INS	I5	Instrument response 0 15-100 WWSSN Resp 1 30-100 WWSSN Resp 2 LRSM 6824-13 Resp 3 LRSM 6824-2 Resp
	76-80	IPNCH		GT.O Card output of seismogram LE.O No card output

The program now returns to Point G to read
in a new set of R, etc.

WIGGLE PAGE 1

C PROGRAM WIGGLE
C THIS PROGRAM READS EIGENFUNCTION TAPES AND CALCULATES SYNTHETIC
C SURFACE WAVE TIME HISTORIES
C GROUND MOTION CONVENTION US VERT + = UP
C RADIAL + = AWAY FROM SOURCE
C TRANS + = CLOCKWISE FROM SOURCE
COMMON/ONE/ELIP(42,10),WVNUM(42,10)
COMMON/TWO/IMODE(50),N
COMMON/THREE/ZIP(50),ZIP(1050)
DIMENSION PYY1(42,10),PYY2(42,10),PYY3(42,10),PYY4(42,10)
DIMENSION SR(50),SI(50),ATTEN(42,10),XRL(42,10),XIM(42,10)
DIMENSION DATA(2,1024)
DIMENSION IBUF(800),ID(6),X(1024),Y(1024)
DIMENSION DATM(2048)
EQUIVALENCE (DATA(1,1),DATM(1))
EQUIVALENCE(DATM(1),X(1)),(DATM(1025),Y(1))
EQUIVALENCE (ZIP(1),SR(1)),(ZIP(51),SI(1)),(ZIP(101),XRL(1)),
1 (ZIP(521),XIM(1))
CHARACTER IDATE(3),ITYP(2),IEX(2),IST,ICOM,IA
CHARACTER STA
IEX(1) = 4HEQ
IEX(2) = 4HEXPL
ITYP(1) = 4HLOVE
ITYP(2) = 4HRAYL
CALL PLOTS(IBUF,800,10)
CALL PLOT(0,0,-11.0,-3)
CALL PLOT(0,0,1,0,-3)
DEGRAD = 0.017452329
NWRITE = 1
3 FORMAT(1H ,8E15.8)
4 FORMAT(8F10.3)
5 FORMAT(2F10.3,8I5,F10.3,2I5)
6 FORMAT(2I5,F10.5)
7 FORMAT(10X,4HMAX ,A4,2H ==,E11.4)
8 FORMAT(F10.5,1X,A4)
11 FORMAT(1H1,A4,6X,4HDOT ==,F10.3)
12 FORMAT(1H0,7X,3HDIP,6X,4HSLIP,4X,6HSTRIKE)
13 FORMAT(1H0,7HDEPTH == ,F10.3)
14 FORMAT(1H0,6HPHI == ,F10.3 ,1X,A4)
15 FORMAT(1H ,A4)
C IEQEX = 1 EARTHQUAKE = 2 EXPLOSION AND DIP SLIP STRK HAVE NO
C MEANING
C JSRC ,GT, 0 READ IN SEISMIC MOMENT IN SRCIN AND NOT SOURCE
C SPECTRUM
C DT = DIGITIZING INTERVAL FOR RESULTANT SEISMOGRAM
READ(60,6) IEQEX,JSRC,DT
WRITE(61,11) IEx(IEQEX),DT
DO 500 K = 1,4
DO 501 I = 1,2
DO 501 J = 1,1024

```

501 DATA(I,J) = 0.0
      DATA(1,2) = 1./(4.*DT)
      DATA(1,3) = 1./(2.*DT)
      DATA(1,4) = 1./(4.*DT)
      CALL FOUR(DATA,1024,-1,DT,DF),
      INS = K - 1
      CALL RESP(DATA,DF,INS)
      DO 502 I = 513,1024
         J = 1026 - I
         DATA(1,I) = DATA(1,J)
502 DATA(2,I)=-DATA(2,J)
      DATA(2,513) = 0.0
      CALL FOUR(DATA,1024,+1,DT,DF),
C      SPECIAL STEPS ARE TAKEN BECAUSE OF EQUIVALENCED ARRAYS
      DO 504 I = 1,1020
         J = 2*I - 1
504 DATM(I) = DATM(J)
      DO 505 I = 1,1020
505 Y(I) = DATM(I)
      DO 506 I = 1,1020
506 X(I) = (I-1)*DT
      CALL SCALE(Y,1.0,1020,1)
      CALL SCALE(X,5.10,1020,1)
      CALL AXIS(0.25,0.0,4HSEC , -4,5.10,90.0,X(1021),X(1022))
      CALL AXIS(0.0,-0.25,4HINST,4,1.0,180.0,Y(1021),Y(1022))
      Y(1022) = - Y(1022)
      CALL LINE(Y,X,1020,1.0,0)
      CALL PLOT(3.0,0.0,0,-3)
500 CONTINUE
      DF = 1./(1024.*DT)
4321 CONTINUE
      READ(60,4) DIP,SLIP,STRK
      WRITE(61:12)
      WRITE(61:4) DIP,SLIP,STRK
C      DIP      DIP OF FAULT PLANE MEASURED FROM HORIZONTAL CLOCKWISE
C      Z DOWN
C      SLIP     DIRECTION OF SLIP VECTOR MEASURED CCL FROM HORIZONTAL
C      STRK    STRIKE OF FAULT PLANE MEASURED FROM NORTH
      XDIP = DEGRAD * DIP
      XSLP = DEGRAD * SLIP
      XSTR = DEGRAD * STRK
      SIND = SIN(XDIP)
      COSD = COS(XDIP)
      COS2D = COS(2.*XDIP)
      SINS = SIN(XSLP)
      COSS = COS(XSLP)
      COSST = COS(XSTR)
      SINST = SIN(XSTR)
      SIN2ST = SIN(2.*XSTR)
      COS2ST = COS(2.*XSTR)

```

```

FAC1 = SIND*COSS*COS2ST
FAC2 = SIND*COSS*SIN2ST
FAC3 = -SIND*COSD*SINS*COS2ST
FAC4 = -SIND*COSD*SINS*SIN2ST
FAC5 = COS2D*SINS*COSS
FAC6 = COS2D*SINS*SINST
FAC7 = -COSD*COSS*COSS
FAC8 = -COSD*COSS*SINST
FAC9 = 2.*SIND*COSD*SINS
2001 READ(60,4) AA
      WRITE(61,13) AA
      IF(AA.LE.-10.) GO TO 4321
      IF(AA.LT.0) GO TO 9998
      CALL EIGEN(PYY1,PYY2,PYY3,PYY4,IFUNC,AA,NPER)
      WRITE(61,15) ITYP(IFUNC)
      DO 1001 I = 1,1050
1001 ZIP(I) = 0.0
      CALL SRCIN(SR,SI,TT,NPER,JSRC)
      CALL GAMMA(ATTEM,TT,NPER)
      CALL SYMBOL(0.0,2.0,0.14,ITYP(IFUNC),90.0,4)
      CALL SYMBOL(0.0,4.0,0.14,4HH = ,90.0,4)
      CALL NUMBER(0.0+4.5+0.14*AA,90.0+1)
      CALL PLOT(3.0,0.0,-3)
2112 CONTINUE
      READ(60,8) PHI,STA
      C   PHI      STATION AZIMUTH MEASURED FROM NORTH
      IF(PHI.LT.-400.) GO TO 2001
      WRITE(61,14) PHI,STA
      SAZ = SIN(DEGRAD * PHI)
      CAZ = COS(DEGRAD * PHI)
      DO 1000 I = 1,NPER
      MODE = IMODE(I)
      DO 1000 J = 1,MODE
      PY1 = PYY1(I,J)
      PY2 = PYY2(I,J)
      PY3 = PYY3(I,J)
      PY4 = PYY4(I,J)
      CALL EXCITT(FAC1,FAC2,FAC3,FAC4,FAC5,FAC6,FAC7,FAC8,FAC9,PY1,
      1 PY2,PY3,PY4,IFUNC,SAZ,CAZ,XR,XI,IEQEX)
      XRL(I,J) = SR(I)*XR - SI(I)*XI
      XIM(I,J) = SR(I)*XI + SI(I)*XR
1000 CONTINUE
3113 CONTINUE
      C   READ IN EPICENTRAL DISTANCE AND STARTING TIME OF SEISMOGRAM TRACE
      READ(60,5) R,T0,ID(1),ID(2),ID(3),ID(4),ID(5),ID(6),IURUZ
      1 NPY,YAXLEN,INS,IPNCH
      IF(R.LE.0) GO TO 2112
      IF(NPY.LE.0) NPY = 1020
      IF(YAXLEN.LE.0) YAXLEN = 5.12
      YAXLE=YAXLEN

```

```

NOTE=YAXLEN
WRITE(61,5)R,T0, ID(1),ID(2),ID(3),ID(4),ID(5),ID(6),IURUZ
1,NPY,YAXLEN,INS,IPNCH
C   R EPICENTRAL DISTANCE IN KILOMETERS
C   TO TIME SEISMOGRAM BEGINS
C   ID CONTROLS THE OUTPUT OPTIONS
C   ID(1) GROUND DISP    LE 0 SKIP  LE 2 SEISMOGRAM  GE 2 SPECTRA
C   ID(2) GROUND VEL    LE 0 SKIP  LE 2 SEISMOGRAM  GE 2 SPECTRA
C   ID(3) GROUND ACCL   LE 0 SKIP  LE 2 SEISMOGRAM  GE 2 SPECTRA
C   ID(4) SEISMO DISP   LE 0 SKIP  LE 2 SEISMOGRAM  GE 2 SPECTRA
C   ID(5) SEISMO VEL   LE 0 SKIP  LE 2 SEISMOGRAM  GE 2 SPECTRA
C   ID(6) SEISMO ACCL   LE 0 SKIP  LE 2 SEISMOGRAM  GE 2 SPECTRA
C   IURUZ EQ 1 VERTICAL COMP OF RAYL  EQ 2 RADIAL COMPONENT OF RAYL
C   NPY = NUMBER OF TIME STEPS TO PLOT FOR SEISMOGRAM  NPY,LE,1024
C   YAXLEN = LENGTH OF TIME AXIS IN INCHES
C   INS = 0 15-100 WWSSA RESP
C         = 1 30-100 WWSSA RESP
C         = 2 LRSN 6824-13 RESP
C         = 3 LRSN 6824-2 RESP
C   IPNCH GT 0 CARD OUTPUT OF SEISMOGRAM
CALL SYMBOL(0.0,0.0,0.14,4HDIST,90.0,4)
CALL NUMBER(0.0,0.7,0.14,R,90.0,-1)
CALL SYMBOL(0.0,2.0,0.14,4H PHI,90.0,4)
CALL NUMBER(0.0,2.7,0.14,PHI,90.0,-1)
CALL SYMBOL(0.0,3.5,0.14,STA,90.0,4)
CALL PLOT(1.5,0.0,-3)
CALL INTER(TT,DATA,NPER,DF,IMODE,XRL,XIM,WVNUM,ATTEN,[FUNC,
1 R,IURUZ,ELIP,TG)
WRITE(54,3) (DATM(I),I=1,1026)
REWIND 54
DO 5000 INK = 1,6
IF(ID(INK).LE.0) GO TO 5000
READ(54,3) (DATM(I),I=1,1026)
REWIND 54
GO TO (5001,5001,5001,5002,5002,5002),INK
5002 CALL RESP(DATA,DF,INS)
5001 CONTINUE
GO TO (5011,5012,5013,5011,5012,5013),INK
5013 CALL VELRES(DATA,DF)
5012 CALL VELRES(DATA,DF)
5011 CONTINUE
IF(ID(INK).LT.2) GO TO 4115
IST = 4HSYNT
ICOM = 4HHETI
IDATE(1) = 4HC SE
IDATE(2) = 4HISMO
IDATE(3) = 4HGRAM
CALL TAPEWR(DATA,128,NWRITE,IST,ICOM,R,PHI,0.0,T0,DT,DATE,1024)
4115 CONTINUE
IF(ID(INK).GT.2) GO TO 5100

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DO 5500 I = 513,1024
J = 1026 - I
DATA(1,I) = DATA(1,J)
5500 DATA(2,I) = - DATA(2,J)
DATA(2,513) = 0.0
CALL FOUR(DATA,1024,+1,DT,DF)
FACC = SQRT(1000./R)
DO 5600 I = 1,1024
J = 2 * I - 1
5600 DATM(I) = DATM(J) * FACC
Y(1) = T0
DO 5700 I = 2,1024
5700 Y(I) = Y(I-1) + DT
GO TO (5701,5702,5703,5701,5702,5703),INK
5701 IA = 4HAMPL
GO TO 5704
5702 IA = 4HVEL
GO TO 5704
5703 IA = 4HACCL
5704 CONTINUE
CALL SCALE(X,1.0,NPY,1)
CALL SCALE(Y,YAXLEN,NPY,1)
CALL AXIS(0.25,0.0,4HSEC,-4,YAXLE,90.0,Y(NPY+1),Y(NPY+2))
CALL AXIS(0.0,-0.25,IA,4,1.0,180.0,X(NPY+1),X(NPY+2))
NPY2 = NPY+2
X(NPY2) = - X(NPY2)
CALL LINE(X,Y,NPY,1,0,0)
CALL PLOT(3.0,0.0,-3)
XMAX = 0
DO 5800 I = 1,NPY
IF(XMAX.LT.ABS(X(I))) XMAX = ABS(X(I))
5800 CONTINUE
WRITE(61,7) IA,XMAX
IF(INK.LT.4) INSTM = 6
IF(INK.GT.3.AND.INS.EQ.0) INSTM = 1
IF(INK.GT.3.AND.INS.GT.0) INSTM = 2
IF(IPNCH.GT.0) CALL PNCH(X,NPY,DT,XMAX,PHI,R,T0,INSTM)
5100 CONTINUE
5000 CONTINUE
GO TO 3113
9998 CONTINUE
CALL PLOT(10.0,0.0,0,999)
STOP
END

SUBROUTINE EIGENF(PYY1,PYY2,PYY3,PYY4,IFUNC,AA,NPER)
COMMON/ONE/ELIP(42,10),WVNUM(42,10)
COMMON/TWO/IMODE(50),N
COMMON/THREE/TT(50),ZIP(1050)
DIMENSION AY(70,4),CY(70,4),DEPTH(70),D(70)

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DIMENSION A(70),B(70),RHO(70),XMU(70),XLAMB(70)
DIMENSION PYY1(42,10),PYY2(42,10),PYY3(42,10),PYY4(42,10)
DIMENSION PYR(4)
1 FORMAT(1H ,2(I4,6X))
2 FORMAT(1H ,I4,5X,7F10.,4)
5 FORMAT(1H ,A4,3F10.5,5E11.,4)
6 FORMAT(1H ,7(4X,E11.4))
DO 37 I = 1,42
IMODE(I) = 1
DO 37 J=1,10
ELIP(I,J) = 0
WVNUM(I,J) = 0.0
PYY1(I,J) = 0
PYY2(I,J) = 0
PYY3(I,J) = 0
37 PYY4(I,J) = 0
READ(01,1)N
C READ IN MODEL
DO 100 I=1,N
100 READ(01,2) LAYER,DEPTH(I),D(I),A(I),B(I),RHO(I),XMU(I),XLAMB(I)
NPER = 0
2000 READ(01,1) IFUNC, MODE
IF(NPER.EQ.0) IFUN = IFUNC
MAXY = 2 * IFUNC
IF(IFUNC.LT.0) GO TO 9999
NPER = NPER + 1
IMODE(NPER) = MODE
DO 2999 IM = 1,MODE
READ(01,5) MODE,T,C,L,SUM0,WVNUM(NPER,IM),AL,ELIP(NPER,IM),ARSOR
TT(NPER) = 1
C THE SPECTRA ARE COMPUTED AT A DISTANCE OF 1000 KILOMETERS
WVNO = WVNUM(NPER,IM)
DENOM = SQRT(WVNO)*0.3989423*1.0E+20
DENOM = DENOM * 6.2831853 * 31.622776 / AL
DO 2100 IN=1,N
2100 READ(01,6) AY(IN,1),AY(IN,2),AY(IN,3),AY(IN,4),DCDA,DCDB,DCDR
WVN02 = WVNO * WVNO
C RAYLEIGH WAVES
C INITIALLY AY(IN,1) = UR, AY(IN,2) = UZ, AY(IN,3) = TZ,
C AY(IN,4) = TR
C NOW AY(IN,1) = UR, AY(IN,2) = UZ, AY(IN,3) = DURDZ,
C AY(IN,4) = DUZDZ
C AND CY(IN,1) = DURDZ, CY(IN,2) = DUZDZ, CY(IN,3) = D2URDZ2,
C CY(IN,4) = D2UZDZ2
C GO TO (2400,2200),IFUNC
2200 DO 2300 IN = 1,N
H = RHO(IN)*C*C*WVN02
XL2M = XLAMB(IN) + 2.*XMU(IN)
CY(IN,1) = -WVNO      *AY(IN,2) + AY(IN,4)/XMU(IN)
CY(IN,2) = (WVNO      *XLAMB(IN)*AY(IN,1)+AY(IN,3))/XL2M

```

```

CY(IN,3) = (-H/XMU(IN) + WVNO2*(3,*XL2M-2,*XMU(IN))/XL2M)*AY(IN,1)
1 -WVNO * (1,+XLAMB(IN)/XMU(IN))*AY(IN,3)/XL2M
CY(IN,4) = (-(H+WVNC2*XLAMB(IN))*AY(IN,2)+WVNO * (1,+XLAMB(IN)-
1 XMU(IN))*AY(IN,4))/XL2M
AY(IN,3) = CY(IN,1)
AY(IN,4) = CY(IN,2)
2300 CONTINUE
GO TO 2501
C LOVE WAVES
C INITIALLY AY(IN,1) = UT, AY(IN,2) = TT
C NOW AY(IN,1) = UT, AY(IN,2) = DUTDZ
C AND CY(IN,1) = DUTDZ, CY(IN,2) = D2UTDZ2
2400 DO 2500 IN=1,N
CY(IN,1) = AY(IN,2)/XMU(IN)
CY(IN,2) = WVNO2*(XMU(IN)- C*C*RHO(IN))*AY(IN,1)/XMU(IN)
AY(IN,2) = CY(IN,1)
2500 CONTINUE
2501 CONTINUE
DO 103 I = 1,N
IFOCA = I
D1=DEPTH(I)-0.5*D(I)
D2=DEPTH(I)+0.5*D(I)
IF(AA.GE.D1.AND.AA.LE.D2) GO TO 101
103 CONTINUE
101 DO 110 JJ = 1,MAXY
AAA = AY(IFOCA,JJ)
CC = CY(IFOCA,JJ)
AAA = AAA - CC * DEPTH(IFOCA)
110 PYR(JJ) = AAA + CC * AA
GO TO(2551,2552),IFUNC
2551 PYY1(NPER,IM) = WVNO * PYR(1)/DENOM
PYY2(NPER,IM) = PYR(2) / DENOM
GO TO 2555
2552 CONTINUE
PYY1(NPER,IM) = WVNC * PYR(1) / DENOM
PYY2(NPER,IM) = WVNC * PYR(2) / DENOM
PYY3(NPER,IM) = PYR(3)/DENOM
PYY4(NPER,IM) = PYR(4)/DENOM
2555 CONTINUE
2999 CONTINUE
GO TO 2000
9999 CONTINUE
IFUNC = IFUN
REWIND 1
RETURN
END

SUBROUTINE TAPEWR(X,M,NWRITE,IST,ICOM,DIST,DEG,BACKAZ,T0,DT,
1 IDATE,NPTS)
DIMENSION X(1),IDATE(3)

```

WIGGLE PAGE 8

```
      WRITE(03,310) M,NPTS,IST,ICOM,DIST,DEG,BACKAZ,T0,DT,
1 IDATE(1),IDATE(2),IDATE(3)
310 FORMAT(2I5,A4,1X,A4,1X,5F10.2,3A4)
DO 312 I = 1,M
K = I * 8
J = K - 7
312 WRITE(03,311) NWRITE,(X(L),L=J,K)
311 FORMAT(I5,8E15.0)
NWRITE = NWRITE + 1
RETURN
END

SUBROUTINE SEISMAG1(FREQ,PEAK,XR,XI,INS)
C INS EQ 0 15-100 WWSSN
C PEAK MAGNIFICATIONS ARE 350,700,1400,2800,5600
C INS EQ 1 30-100 WWSSN
WE=6.2831853*FREQ
INDEX=(PEAK+1)/375
IF(INS.GT,0) GO TO 200
100 CONTINUE
GO TO(1,2,2,3,3,3,4,4,4,4,4,4,4,4,5),INDEX
1 FMAG=270,
SIGMA=0.003
GO TO 6
2 FMAG=556.0
SIGMA=0.013
GO TO 6
3 FMAG=1110,
SIGMA=0.047
GO TO 6
4 FMAG=2190,
SIGMA=0.204
GO TO 6
5 FMAG=3950,
SIGMA=0.805
6 ZETA=0.93
ZETA1-i
WN=-418879
WN1=.062831853
GO TO 300
200 CONTINUE
GO TO(10,20,20,30,30,30,30,40,40,40,40,40,40,40,40,50),INDEX
10 FMAG = 251.9
SIGMA = 0.003
GO TO 60
20 FMAG = 503.1
SIGMA = 0.012
GO TO 60
30 FMAG = 1001.5
SIGMA = 0.044
```

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```
GO TO 60
40 FMAG = 1941.9
SIGMA = 0.195
GO TO 60
50 FMAG = 2241.8
SIGMA = .767
60 ZETA = 1.5
ZETA1 = 1.0
WN = .2094395
WN1 = .062831853
30 CONTINUE
AR = (WE*WE-WN*WN)*(WE*WE-WN1*WN1)-4.*ZETA*ZETA1*WN*WN1*(1.-SIGMA)
1*WE*WE
AI=2.*WE*(ZETA1*WN1*(WN*WN-WE*WE)+ZETA*WN*(WN1*WN1-WE*WE))
FACTOR = FMAG*WE*WE*WE / (AI*AI + AR*AR)
XR = -AI * FACTOR
XI = -FACTOR * AR
RETURN
END
```

SUBROUTINE SEISM3(FREQ,PEAK,XR,XI)

```
C LRSW RESPONSE FOR LP SYSTEM WITH FILTER 6824-2
C PHASE RESPONSE OBTAINED FROM HILBERT TRANSFORM OF AMPLITUDE
C RESPONSE. GAIN NORMALIZED TO 1.0 AT 25 SECONDS.
DIMENSION FRE(28),P(28),PHI(28)
DATA FRE/.001,.002,.003,.004,.005,.006,.007,.008,.009,.01,.02,.03,
1.04,.05,.06,.07,.08,.09,.1,.2,.3,.4,.5,.6,.7,.8,.9,1./
DATA PHI/263.9,257.9,251.9,245.6,239.3,233.1,227.1,221.1,214.6,
120.3,134.9,73.2,15.3,-33.2,-71.2,-100.1,-122.4,-140.0,-153.8,-213
2.1,-232.6,-242.8,-249.2,-253.0,-256.9,-259.3,-261.2,-262.9/
DATA F/.00005,.00040,.00135,.00321,.00625,.01077,.01706,.02546,
1.03631,.05006,.34117,.73904,1.0000,.97633,.79807,.61417,.46105,
2.34464,.26315,.03583,.01097,.00468,.00240,.00139,.00087,.00058,
3.00040,.00029/
M = 28
DEGRAD = 0.01745329
IF(FREQ.GT.0.5) FREQ = 0.5
IF(FREQ.LT.0.005) FREQ = 0.005
DO 200 I = 1,M
IF(FREQ.GE.FRE(I)).AND.FREQ.LE.FRE(I+1)) GO TO 160
200 CONTINUE
160 PF = PHI(I)+(PHI(I+1)-PHI(I))/(FRE(I+1)-FRE(I))*(FREQ-FRE(I))
PH = P(I) + (P(I+1) - P(I))/(FRE(I+1)-FRE(I))*(FREQ-FRE(I))
PF = PF * DEGRAD
XR = COS(PF) * PH * PEAK
XI = SIN(PF) * PH * PEAK
RETURN
END
```

SUBROUTINE SEISM7(FREQ,PEAK,XR,XI)

```

C      LRSM RESPONSE FOR LP SYSTEM WITH FILTER 6824-13
C      PHASE RESPONSE OBTAINED FROM HILBERT TRANSFORM OF AMPLITUDE
C      RESPONSE. GAIN NORMALIZED TO 1.0 AT 25 SECONDS.
C      DIMENSION FRE(28),P(28),PHI(28)
C      DATA FRE/.001,.002,.003,.004,.005,.006,.007,.008,.009,.01,.02,.03,
1.04,.05,.06,.07,.08,.09,.1,.2,.3,.4,.5,.6,.7,.8,.9,1./
C      DATA P/.000015,.000391,.00131,.00310,.00609,.01068,.01713,.02556,
1.03618,.04848,.28030,.67553,1.0,1.09448,1.0044,.86969,.74511,.6355
27,.54818,.12153,.04830,.02741,.01827,.01328,.01017,.008052,.006545
3,.005404/
C      DATA PHI/265.0,259.1,253.4,248.1,242.8,236.9,230.3,223.4,216.4,
1209.7,153.4,102.0,54.6,14.4,-16.8,-40.4,-58.2,-73.3,-85.4,-146.4,
-153.9,-155.3,-156.4,-157.7,-159.2,-160.6,-162.3,-163.8/
M = 28
DEGRAD = 0.01745329
IF(FREQ.GT.0.5) FREQ = 0.5
IF(FREQ.LT.0.005) FREQ = 0.005
DO 200 I = 1,M
IF(FREQ.GE.FRE(I).AND.FREQ.LE.FRE(I+1)) GO TO 160
200 CONTINUE
160 PF = PHI(I)+(PHI(I+1)-PHI(I))/(FRE(I+1)-FRE(I))*(FREQ-FRE(I))
PH = P(I) + (P(I+1) - P(I))/(FRE(I+1)-FRE(I))*(FREQ-FRE(I))
PF = PF * DEGRAD
XR = PH*PEAK*COS(PF)
XI = PH*PEAK*SIN(PF)
RETURN
END

SUBROUTINE RESP(DATA,DF,INS)
DIMENSION DATA(1)
DATA(1)=0.0
DATA(2) = 0.0
DO 10 I = 2,513
XI = I - 1
FREQ = XI * DF
J = 2 * I - 1
K = 2 * I
IF(INS.LT.2) CALL SEISMAG1(FREQ,3000.,XR,XI,INS)
IF(INS.EQ.2) CALL SEISM7(FREQ,3000.,XR,XI)
IF(INS.EQ.3) CALL SEISM3(FREQ,3000.,XR,XI)
TEMPI = XR*DATA(J) - XI*DATA(K)
TEMPR = XI*DATA(K) + XI*DATA(J)
DATA(J) = TEMPI
10 DATA(K) = TEMPR
RETURN
END

SUBROUTINE VELRES(X,dF)
DIMENSION X(1)
C      THIS IS FREQUENCY DOMAIN EQUIVALENT OF TIME DOMAIN DIFFERENTIATION

```

```

DOM = 6.2831853 * DF
DO 10 I = 1,513
X1 = I - 1
FREQ = XI * DOM
J = 2 * I - 1
K = 2 * I
TEMPIR = - FREQ * X(K)
TEMPII = FREQ * X(J)
X(J) = TEMPIR
10 X(K) = TEMPII
RETURN
END

SUBROUTINE FOUR(DATA,NN,ISIGN,DT,DF)
C THE COOLEY-TOUKEY FAST FOURIER TRANSFORM IN USASI BASIC FORTRAN
C TRANSFORM(J) = SUM(DATA(I)*W**((I-1)(J-1)), WHERE I AND J RUN
C FROM 1 TO NN AND W = EXP(ISIGN*2*PI*SQRT(-1)/NN). DATA IS A ONE-
C DIMENSIONAL COMPLEX ARRAY (I.E., THE REAL AND IMAGINARY PARTS OF
C DATA ARE LOCATED IMMEDIATELY ADJACENT IN STORAGE, SUCH AS
C FORTRAN IV PLACES THEM) WHOSE LENGTH NN IS A POWER OF TWO. ISIGN
C IS +1 OR -1, GIVING THE SIGN OF THE TRANSFORM. TRANSFORM VALUES
C ARE RETURNED IN ARRAY DATA, REPLACING THE INPUT DATA. THE TIME IS
C PROPORTIONAL TO N*LOG2(N), RATHER THAN THE USUAL N**2
C RMS RESOLUTION ERROR BEING BOUNDED BY 6*SQRT(1)*LOG2(NN)*2**(-B),
C WHERE B IS THE NUMBER OF BITS IN THE FLOATING POINT FRACTION.
C PROGRAM AUTOMATICALLY DIVIDES TRANSFORM BY NN FOR INVERSE
C TRANSFORM
C DIMENSION DATA(1)
C N = 2 * NN
C IF(DT.EQ.0.0) DT=1./(NN*DF)
C IF(DF.EQ.0.0) DF=1./(NN*DT)
C IF(DT.NE.(NN*DF)) DF = 1./(NN*DT)
C J = 1
C DO 5 I=1,N,2
C IF(I-J)1,2,2
1 TEMPIR = DATA(J)
TEMPII = DATA(J+1)
DATA(J) = DATA(I)
DATA(J+1)=DATA(I+1)
DATA(I) = TEMPIR
DATA(I+1) = TEMPII
2 M = N/2
3 IF(J-M)5,5,4
4 J = J-M
M = M/2
IF(M-2)5,3,3
5 J=J+M
MMAX = 2
6 IF(MMAX-N)7,10,10
7 ISTEP= 2 *MMAX

```

```

THETA = 6.283185307/FLOAT(ISIGN*MMAX)
SINTH=SIN(THETA/2.)
WSTPR=-2.*SINTH*SINTH
WSTPI=SIN(THETA)
WR=1.0
WI=0.0
DO 9 M=1,MMAX,2
DO 8 I=M,N,ISTEP
J=I+MMAX
TEMPR=WR*DATA(J)-WI*DATA(J+1)
TEMPI=WR*DATA(J+1)+WI*DATA(J)
DATA(J)=DATA(I)-TEMPR
DATA(J+1)=DATA(I+1)-TEMPI
DATA(I)=DATA(I)+TEMPR
8 DATA(I+1) = DATA(I+1)+TEMPI
TEMPR = WR
WR = WR+WSTPR-WI*WSTPI + WR
9 WI = WI*WSTPR+TEMPR*WSTPI + WI
MMAX = ISTEP
GO TO 6
11 CONTINUE
IF(ISIGN.LT.0) GO TO 1002
C   FREQUENCY TO TIME DOMAIN
DO 1001 IIII=1,N
1001 DATA(IIII)=DATA(IIII)*DF
RETURN
1002 CONTINUE
C   TIME TO FREQUENCY DOMAIN
DO 1003 IIII=1,N
1003 DATA(IIII)=DATA(IIII)*DT
RETURN
END

SUBROUTINE EXCITT(FAC1,FAC2,FAC3,FAC4,FAC5,FAC6,FAC7,FAC8,FAC9,
1 PY1,PY2,PY3,PY4,IFUNC,SAZ,CAZ,XR,XI,IEQEX)
GO TO (100,200),IEQEX
100 CONTINUE
S2AZ = 2.*SAZ*CAZ
C2AZ = CAZ*CAZ - SAZ*SAZ
GO TO (1,2),IFUNC
1 CONTINUE
XR = PY1*((FAC1-FAC4)*C2AZ + (FAC2+FAC3)*S2AZ )
XI = - PY2*((FAC5+FAC8)*CAZ + (FAC6-FAC7)*SAZ)
RETURN
2 CONTINUE
XR = PY1*((FAC3+FAC2)*C2AZ + (FAC4-FAC1)*S2AZ) + FAC9 + PY4
1 + 0.5*FAC9*PY1
XI = (PY3+PY2)*((FAC7-FAC6)*CAZ+(FAC8+FAC5)*SAZ)
RETURN
200 CONTINUE

```

```

      GO TO (10,20),IFUNC
10  XR = 0
    XI = 0
    RETURN
20  XR = PY4 - PY1
    XI = 0
    RETURN
    END

SUBROUTINE SRCIN(SR,SI,TT,NPER,JSRC)
DIMENSION SR(1),SI(1),TT(1)
DIMENSION XR(50),XI(50),T(50)
IF(JSRC.GT.0) GO TO 400
C THIS READS IN THE VALUES OF THE SOURCE SPECTRUM AT PERIODS T WHICH
C ARE ALIGNED IN ASCENDING ORDER. THE SOURCE VALUES CORRESPONDING
C TO THE PERIODS TT OF THE EIGENFUNCTIONS ARE FOUND BY
C INTERPOLATION
DO 100 I = 1,50
READ(60,2) T(I),XR(I),XI(I)
2 FORMAT(F10.2,E11.4,9X,E11.4)
IF(T(I).LE.0) GO TO 101
N = I
100 CONTINUE
101 CONTINUE
N = N - 1
DO 300 K = 1,NPER
PER = TT(K)
DO 210 I = 1,N
IF(PER.GE.T(I).AND.PER.LE.T(I+1)) GO TO 260
210 CONTINUE
260 SR(K) = XR(I) + (XR(I+1) - XR(I))/(T(I+1)-T(I))*(PER-T(I))
SI(K) = XI(I) + (XI(I+1) - XI(I))/(T(I+1)-T(I))*(PER-T(I))
300 CONTINUE
GO TO 401
400 CONTINUE
C THIS READS IN THE SEISMIC MOMENT, ASSUMES A STEP DISLOCATION
C AND COMPUTES THE SOURCE SPECTRUM
READ(60,10) XMOM
10 FORMAT(E10.3)
DO 200 K = 1,NPER
SR(K) = 0.0
SI(K) = - XMOM*TT(K)/6.2831853
200 CONTINUE
401 CONTINUE
WRITE(61,3)
3 FORMAT(1H0,4X,6HPERIOD,4X,7HS(REAL),13X,7HS(IMAG) )
WRITE(61,1)((TT(I),SR(I),SI(I)),I=1,NPER)
1 FORMAT(1H ,F9.2,E11.4,9X,E11.4)
RETURN
END

```

```

SUBROUTINE GAMMA(ATEN,TT,NPER)
DIMENSION ATEN(42,10),TT(50),ATEN(42,10),T(42)
1 FORMAT(8F10.5)
DO 100 I = 1,42
READ(60+1) T(I), (ATEN(I,J),J=1,10)
IF(T(I),LE,0) GO TO 101
N = I
100 CONTINUE
101 CONTINUE
N = N - 1
DO 200 K = 1,NPER
PER = TT(K)
DO 210 I = 1,N
IF(PER.GE.T(I).AND.PER.LE.T(I+1)) GO TO 260
210 CONTINUE
260 DO 270 J = 1,10
270 ATEN(K,J) = ATEN(I,J) + (ATEN(I+1,J)-ATEN(I,J))/(T(I+1)-T(I))
1 * (PER - T(I))
200 CONTINUE
WRITE(61,3)
3 FORMAT(1H0,5X,6HPERIOD,1D5X,5HGAMMA ) )
WRITE(61,2) ((T(I),(ATEN(I,J),J=1,10)),I=1,NPER)
2 FORMAT(1H ,F9.2,10F10.5)
RETURN
END

SUBROUTINE INTER(TT,DATA,NPER,DF,IMODE,XRL,XIM,WVNUM,ATTEN,IFUNC,
1 R,IURU4,ELIP,TJ)
DIMENSION DATA(2,1024),XRL(42,10),XIM(42,10),IMODE(50),ATTEN(42,10
1),WVNUM(42,10),TT(50),ELIP(42,10)
PI4 = 0.7853982
TLW = TT(1)
TMX = TT(NPER)
DATA(1,1) = 0.0
DATA(2,1) = 0.0
DO 200 II = 2,513
DATA(1,II)=0.0
DATA(2,II)=0.0
XK = II - 1
PER = 1.0/(XK*DF)
IF(PER.GT.TMX) GO TO 200
IF(PER.LT.TLW) GO TO 200
202 DO 210 I = 1,NPER-1
IF(PER.GE.TT(I).AND.PER.LE.TT(I+1)) GO TO 260
210 CONTINUE
C THIS USES FACT THAT IMODE(I+1) IS LE TO IMODE(I)
C
260 MODE = IMODE(I+1)
DO 270 J = 1,MODE

```

```

XR = XRL(I,J) + (XRL(I+1,J)-XRL(I,J))/(TT(I+1)-TT(I))*(PER-TT(I))
XI = XIM(I,J) + (XIM(I+1,J)-XIM(I,J))/(TT(I+1)-TT(I))*(PER-TT(I))
WVNO = WVNUM(I,J)+(WVNUM(I+1,J)-WVNUM(I,J))/(TT(I+1)-TT(I))*(PER-
1 TT(I))
ATTEN = ATTEN(I,J)+(ATTEN(I+1,J)-ATTEN(I,J))/(TT(I+1)-TT(I))*(PER-
1 - TT(I))
FACT = (6.2831853*T0/PER) - WVNO * R
GO TO (301,302),IFUNC
301 FACT = FACT + PI4
GO TO 305
302 GO TO (303,304),ILRLZ
303 FACT = FACT - PI4
GO TO 305
304 FACT = FACT - 2.3561945
ELLIP = ELIP(I,J) + (ELIP(I+1,J)-ELIP(I,J))/(TT(I+1)-TT(I))*(PER-
1 - TT(I))
305 CONTINUE
CF = COS(FACT)
SF = SIN(FACT)
FACT = 0.0
ATN = ATEN*R
IF(ATN.LT.80.0) FACT = EXP(-ATN)
IF(IFUNC.EQ.2.AND.IURUZ.EQ.2) FACT = FACT * ELLIP
IF(IFUNC.EQ.2.AND.IURUZ.EQ.1) FACT = - FACT
DATA(1,II) = DATA(1,II) + FACT*(XR*CF - XI*SF)
DATA(2,II) = DATA(2,II) + FACT *(XR*SF + XI*CF)
270 CONTINUE
200 CONTINUE
RETURN
END

```

```

SUBROUTINE PNCH(X,NPY,DT,XMAX,PHI,R,T0,INSTM)
DIMENSION X(1),IX(1025),IY(1025)
DEG = R/111.195
BACKAZ = 360. - PHI
TMAX = (NPY-1)*DT
CPM = 60./DT
CPMM=500./XMAX
WRITE(43,6) INSTM
WRITE(42,6) INSTM
6 FORMAT(15,4X,1H1,4X,1H1,3X,2H-1,4X,1H2,4X,1H2)
WRITE(43,5)
WRITE(42,5)
5 FORMAT(1H )
WRITE(43,3) R,DEG,BACKAZ,T0,DT,CPMM
WRITE(42,3) R,DEG,BACKAZ,T0,DT,CPMM
WRITE(43,4) TMAX,CPM
WRITE(42,4) TMAX,CPM
3 FORMAT(10X,5F10.2,1PE10,3,6H 3000.)
4 FORMAT(25X,F10.2,F10.3)

```

```
DO 100 I = 1,NPY
IX(I) = I-1
IY(I) = 5000*X(I)/XMAX
100 CONTINUE
NPY1 = NPY+1
IX(NPY1)=-9999
IY(NPY1)=-9999
WRITE(43,2)((IX(I),IY(I)),I=1,NPY1)
2 FORMAT(16I5)
WRITE(42,2)((IX(I),IY(I)),I=1,NPY1)
RETURN
END
```

DATA PAGE 1

1 1 0.488
85. 145. 350.
7.5
0.250E 23
1.5 0.01

2.0 0.01

2.5 0.005

5.0 0.0007

10. 0.0002

20. 0.0001

22. 0.00008

30. 0.00003

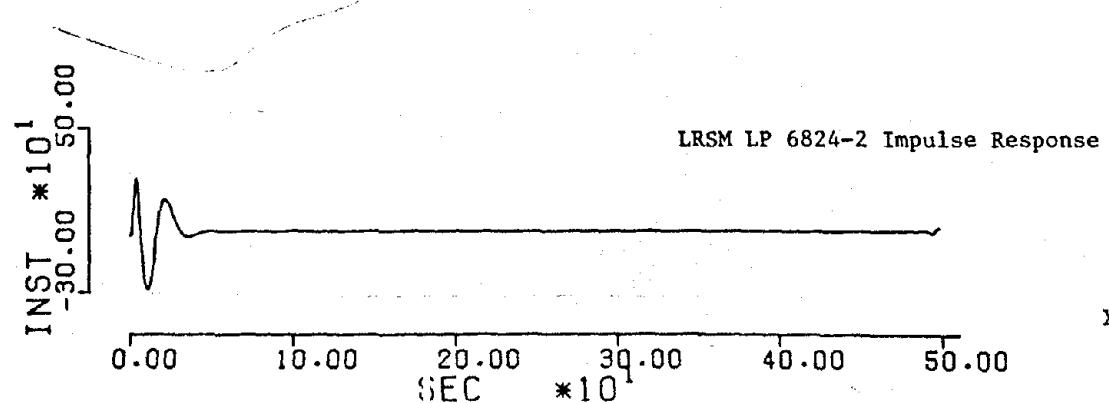
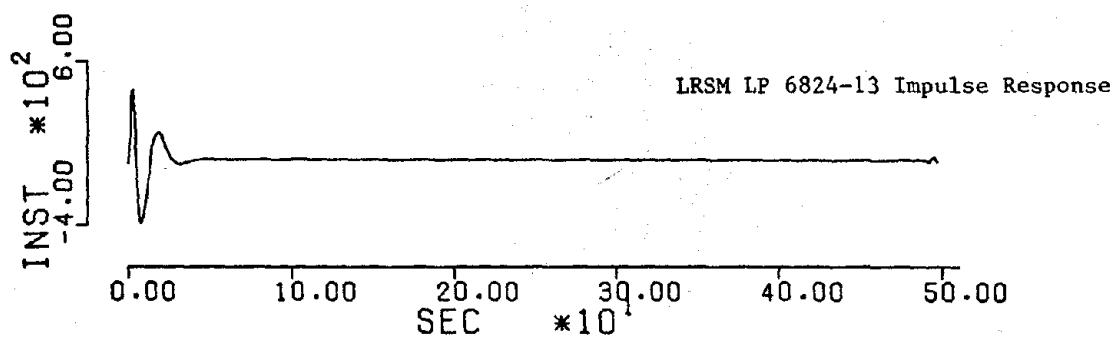
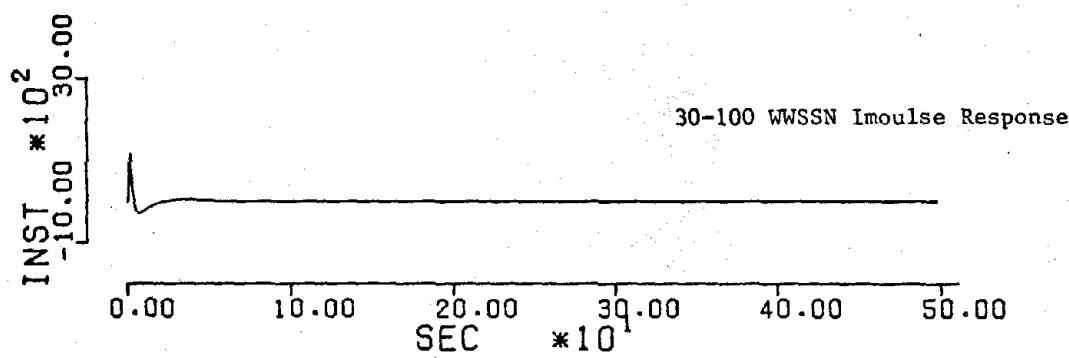
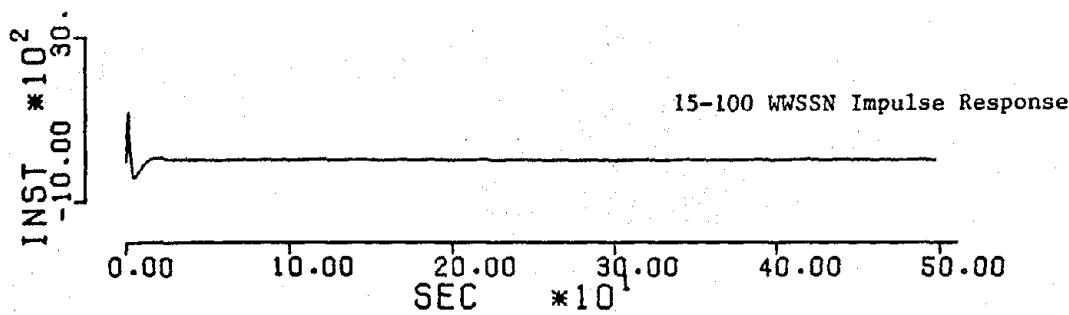
35. 0.00006

100. 0.00006

352.3	DBQ						
854.3	0.0	2		1	480	5.0	0
346.7	SLM						
244.9	0.0	2		1	480	5.0	0
345.4	FLO						
265.2	0.0	2		1	480	5.0	0
76.4	BLWV						
749.0	0.0	2		1	480	5.0	0
262.9	HBOOK						
854.3	0.0	3		1	480	5.0	0
-500.							
-1.							

LOVE

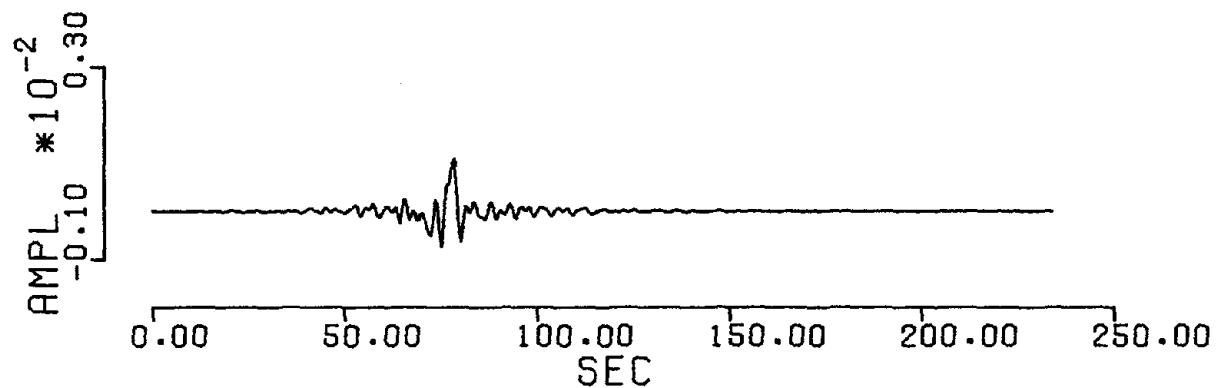
H = 7.5



DIST 265

PHI 345

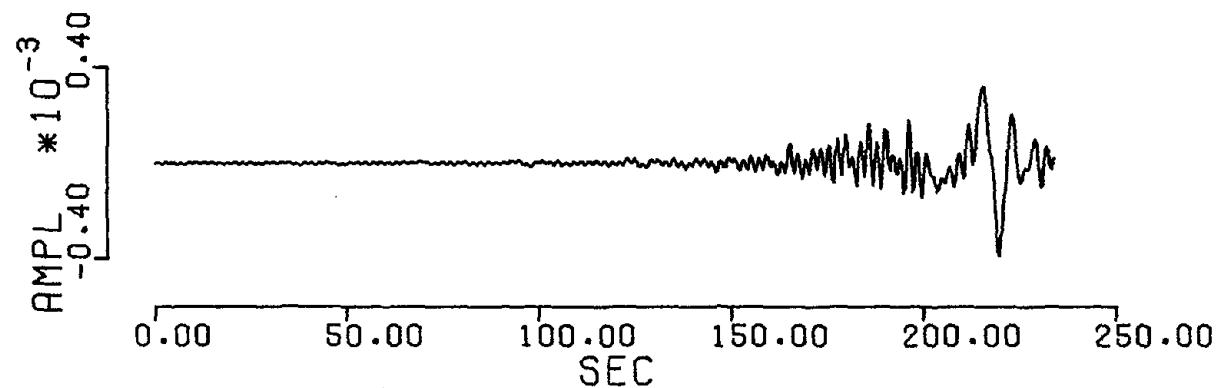
FL0



DIST 750

PHI 76

BLW



EQ DT = 0.488
 DIP SLIP STRIKE
 85.000 145.000 350,000
 DEPTH = 7.500
 LOVE

PERIOD S(REAL)

	S(IMAGE)
2.00 0.	-0.7958E 22
2.50 0.	-0.9947E 22
3.00 0.	-0.1194E 23
3.50 0.	-0.1393E 23
4.00 0.	-0.1592E 23
5.00 0.	-0.1989E 23
6.00 0.	-0.2387E 23
7.00 0.	-0.2785E 23
8.00 0.	-0.3183E 23
9.00 0.	-0.3581E 23
10.00 0.	-0.3979E 23
11.00 0.	-0.4377E 23
12.00 0.	-0.4775E 23
13.00 0.	-0.5173E 23
14.00 0.	-0.5571E 23
15.00 0.	-0.5969E 23
16.00 0.	-0.6366E 23
17.00 0.	-0.6764E 23
18.00 0.	-0.7162E 23
19.00 0.	-0.7560E 23
20.00 0.	-0.7958E 23
22.00 0.	-0.8754E 23
24.00 0.	-0.9549E 23
26.00 0.	-0.1035E 24
28.00 0.	-0.1114E 24
30.00 0.	-0.1194E 24
35.00 0.	-0.1393E 24
40.00 0.	-0.1592E 24
45.00 0.	-0.1791E 24
50.00 0.	-0.1989E 24
60.00 0.	-0.2387E 24
70.00 0.	-0.2785E 24
80.00 0.	-0.3183E 24
90.00 0.	-0.3581E 24
100.00 0.	-0.3979E 24
150.00 0.	-0.5968E 24
200.00 0.	-0.7958E 24

PERIOD	GAMMA	GAMMA	GAMMA	GAMMA
2.00	0.01000	0.	0.	0.
2.50	0.00500	0.	0.	0.
3.00	0.00414	0.	0.	0.
3.50	0.00328	0.	0.	0.
4.00	0.00242	0.	0.	0.
5.00	0.00170	0.	0.	0.
6.00	0.00060	0.	0.	0.
7.00	0.00050	0.	0.	0.
8.00	0.00040	0.	0.	0.
9.00	0.00030	0.	0.	0.
10.00	0.00020	0.	0.	0.
11.00	0.00019	0.	0.	0.
12.00	0.00018	0.	0.	0.
13.00	0.00017	0.	0.	0.
14.00	0.00016	0.	0.	0.
15.00	0.00015	0.	0.	0.

PERIOD	GAMMA	GAMMA	GAMMA	GAMMA
2.00	0.	0.	0.	0.
2.50	0.	0.	0.	0.
3.00	0.	0.	0.	0.
3.50	0.	0.	0.	0.
4.00	0.	0.	0.	0.
5.00	0.	0.	0.	0.
6.00	0.	0.	0.	0.
7.00	0.	0.	0.	0.
8.00	0.	0.	0.	0.
9.00	0.	0.	0.	0.
10.00	0.	0.	0.	0.
11.00	0.	0.	0.	0.
12.00	0.	0.	0.	0.
13.00	0.	0.	0.	0.
14.00	0.	0.	0.	0.
15.00	0.	0.	0.	0.

Appendix A.

A. L. Levshin and Z. A. Yanson

Surface waves in vertically and radially inhomogeneous media.

Algorithms for the interpretation of seismic data. Computational Seismology Series, Volume 5, pp 147-177, Moskva, 1971.

Translated by R. B. Herrmann, Saint Louis University, 1972.

INTRODUCTION

In recent years, the data from surface waves generated by earthquakes and explosions have been used for a detailed investigation of the properties of the earth and of the source of the disturbance (for example, to discern zones of reduced velocity or large velocity gradients, to estimate the distribution of absorption with depth, to determine the mechanism and time function of sources, etc.). In order to solve this relatively difficult mathematical problem, simplifications are usually made in the model of the medium (layered homogeneous medium) and of the source (point source) (5,12,16-20); these simplifications can be shown to be insufficient. In our work, based on results obtained in (2,4,15), we will expound the basic elements of a more general theory of surface waves, valid for very minimal restrictions on the model and source.

We will consider vertically (radially) inhomogeneous media with arbitrary rules for the variation of the elastic constants and density with depth (radius); to find the solution of the equation of motion in such a medium, the spectral theory of operators will be used (6,9). The seismic sources will be treated as fields of volume forces, localized spatially and temporally; the only restrictions imposed on the properties of these fields are those of physical existence. An exact solution is constructed for such a source; asymptotic expansions are made at large distances from the source where the field of the disturbance will separate into propagating Love and Rayleigh waves. Then formulae are obtained for some elementary force fields.

(Some results of Saito's work (22) are used here, but that work is marked by approaches to the final solution which seem quick to us. Hence, a different treatment of the seismic source and the asymptotic expansion is given).

1. Displacement field in an elastic halfspace

Statement of the problem. We will consider an elastic halfspace with coordinates z, r, ϕ ($0 \leq z < \infty, 0 \leq r < \infty, 0 \leq \phi < 2\pi$). The equations of motion are (7):

$$\begin{aligned} \frac{\partial \hat{r}_z}{\partial r} + \frac{1}{r} \frac{\partial \hat{\varphi}_z}{\partial \varphi} + \frac{\partial \hat{z}_z}{\partial z} + \frac{\hat{r}_z}{r} &= \rho \frac{\partial^2 u_z}{\partial t^2} - F_z, \\ \frac{\partial \hat{r}_r}{\partial r} + \frac{1}{r} \frac{\partial \hat{r}_\varphi}{\partial \varphi} + \frac{\partial \hat{r}_z}{\partial z} + \frac{\hat{r}_r - \hat{\varphi}_\varphi}{r} &= \rho \frac{\partial^2 u_r}{\partial t^2} - F_r, \quad (1.1) \\ \frac{\partial \hat{r}_\varphi}{\partial r} + \frac{1}{r} \frac{\partial \hat{\varphi}_\varphi}{\partial \varphi} + \frac{\partial \hat{\varphi}_z}{\partial z} + \frac{2\hat{r}_\varphi}{r} &= \rho \frac{\partial^2 u_\varphi}{\partial t^2} - F_\varphi. \end{aligned}$$

Here \hat{rz} , \hat{rr} , $\hat{r\phi}$, $\hat{\phi\phi}$, \hat{zz} are components of the stress tensor; u_z , u_r , u_ϕ are components of the displacement vector $u(t, z, r, \phi)$ in the directions a_z , a_r , a_ϕ respectively; F_z , F_r , F_ϕ are components of the vector volume force $F(t, z, r, \phi)$ acting at the source along the same directions; t is the time.

For the components of stress we have the following relations:

$$\begin{aligned}\hat{rz} &= \mu \left(\frac{\partial u_z}{\partial r} + \frac{\partial u_r}{\partial z} \right), & \hat{\phi z} &= \mu \left(\frac{1}{r} \frac{\partial u_z}{\partial \phi} + \frac{\partial u_\phi}{\partial z} \right), \\ \hat{rr} &= \lambda \Delta + 2\mu \frac{\partial u_r}{\partial r}, & \hat{\phi\phi} &= \lambda \Delta + 2 \frac{\mu}{r} \left(\frac{\partial u_\phi}{\partial \phi} + u_r \right), \\ \hat{r\phi} &= \mu \left(\frac{\partial u_\phi}{\partial r} - \frac{u_\phi}{r} + \frac{1}{r} \frac{\partial u_r}{\partial \phi} \right), & \hat{zz} &= \lambda \Delta + 2\mu \frac{\partial u_z}{\partial z}.\end{aligned}\quad (1.2)$$

where Δ - dilation.

$$\Delta = \frac{\partial u_z}{\partial z} + \frac{1}{r} \frac{\partial u_\phi}{\partial \phi} + \frac{\partial u_r}{\partial r} + \frac{u_r}{r}. \quad (1.3)$$

The Lamé constants λ and μ and the density ρ are piecewise continuous positive functions of a single coordinate z ; for $z > Z$, λ, μ, ρ are constant, and the velocity of transverse waves $b = \sqrt{\mu/\rho}$ and of compressional waves $a = \sqrt{(\lambda+2\mu)/\rho}$ are maximal (this is necessary since Z can be as large as desired):

$$b(Z+0) = \max b(z), \quad a(Z+0) = \max a(z)^1.$$

The components of displacement and stress are continuous and bounded everywhere in the region $0 \leq z < \infty$; the surface $z = 0$ of stress, i.e.

$$\hat{rz} = \hat{\phi z} = \hat{zz} = 0 \quad \text{при } z = 0. \quad (1.4)$$

Initial conditions:

$$u = \frac{\partial u}{\partial t} = 0 \quad \text{при } t < 0. \quad (1.4a)$$

Source. The force field $F(t, z, r, \phi)$ is described as a real source localized in space and time. The following conditions are imposed on F :

- 1) $F(t, z, r, \phi) = 0$ for $t < 0$)
- 2) $F(t, z, r, \phi)$ is absolutely integrable and satisfies the Dirichlet conditions for all arguments. Then we are permitted the following representations:

$$F(t, z, r, \varphi) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} e^{ipt} \int_0^{\infty} \left[\sum_{m=-\infty}^{\infty} \sum_{i=1}^3 f_m^{(i)} A_m^{(i)} \right] \xi d\xi dp, \quad (1.5)$$

where

$$\begin{aligned} A_m^{(1)} &= a_z Y_m, & A_m^{(2)} &= a_r \frac{\partial Y_m}{\partial r} \frac{1}{\xi} + a_\varphi \frac{\partial Y_m}{\partial \varphi} \frac{1}{\xi r}, \\ A_m^{(3)} &= a_r \frac{\partial Y_m}{\partial \varphi} \frac{1}{\xi r} - a_\varphi \frac{\partial Y_m}{\partial r} \frac{1}{\xi}, & Y_m &= e^{im\varphi} J_m(\xi r). \end{aligned} \quad (1.6)$$

Here J_m is a Bessel function of the first kind of order m . The system of vector functions $A_m^{(i)}$ are fields and mutually orthogonal. The coefficients $f_m^{(i)}$ (z, ξ, p) are found using the orthogonality relation

$$\int_0^{\infty} \int_0^{2\pi} [A_m^{(i)}(\lambda r), \bar{A}_m^{(j)}(\gamma r)] r d\varphi dr = 2\pi \delta_{ij} \delta_{ml} \frac{\delta(\gamma - \lambda)}{\sqrt{\lambda}}$$

(δ_{ij} is the Kronecker delta, $\delta(\gamma - \lambda)$ is the Dirac delta function¹) and equal

$$f_m^{(i)} = \frac{1}{2\pi} \int_{-\infty}^{+\infty} e^{-ipr} \int_0^{\infty} \int_0^{2\pi} (F, \bar{A}_m^{(i)}) r d\varphi dr dt.$$

Specifically for $i = 1, 2, 3$ we have:

$$\begin{aligned} f_m^{(1)} &= \frac{1}{2\pi} \int_{-\infty}^{+\infty} e^{-ipr} \int_0^{\infty} \int_0^{2\pi} [F_z \bar{Y}_m(\xi r)] r d\varphi dr dt, \\ f_m^{(2)} &= \frac{1}{2\pi} \int_{-\infty}^{+\infty} e^{-ipr} \int_0^{\infty} \int_0^{2\pi} \left[\left(F_r \frac{\partial \bar{Y}_m}{\partial r} + F_\varphi \frac{\partial \bar{Y}_m}{\partial \varphi} \frac{1}{r} \right) \right] \frac{r}{\xi} d\varphi dr dt, \\ f_m^{(3)} &= \frac{1}{2\pi} \int_{-\infty}^{+\infty} e^{-ipr} \int_0^{\infty} \int_0^{2\pi} \left[\left(F_r \frac{\partial \bar{Y}_m}{\partial \varphi} \frac{1}{r} - F_\varphi \frac{\partial \bar{Y}_m}{\partial r} \right) \right] \frac{r}{\xi} d\varphi dr dt, \\ \bar{Y}_m &= e^{-im\varphi} J_m(\xi r). \end{aligned} \quad (1.7)$$

For the force components of F_z , F_r , F_φ we have been using (1.5), (1.6):

$$\begin{aligned} F_z &= \frac{1}{2\pi} \int_{-\infty}^{+\infty} e^{ipt} \int_0^{\infty} \left[\sum_{m=-\infty}^{\infty} f_m^{(1)} Y_m \right] \xi d\xi dp, \\ F_r &= \frac{1}{2\pi} \int_{-\infty}^{+\infty} e^{ipt} \int_0^{\infty} \left[\sum_{m=-\infty}^{\infty} \left(f_m^{(2)} \frac{\partial Y_m}{\partial r} + f_m^{(3)} \frac{\partial Y_m}{\partial \varphi} \frac{1}{r} \right) \right] d\xi dp, \\ F_\varphi &= \frac{1}{2\pi} \int_{-\infty}^{+\infty} e^{ipt} \int_0^{\infty} \left[\sum_{m=-\infty}^{\infty} \left(f_m^{(2)} \frac{\partial Y_m}{\partial \varphi} \frac{1}{r} - f_m^{(3)} \frac{\partial Y_m}{\partial r} \right) \right] d\xi dp. \end{aligned} \quad (1.8)$$

$\bar{A}_1^{(j)}$ is complex conjugate of $A_1^{(j)}$

Formulae for displacement. The solutions arising in non-stationary problems of the theory of elasticity have the form

$$u(t, z, r, \varphi) = \frac{1}{2\pi} \int_{-\infty}^{+\infty} e^{ipt} U dp, \quad (1.9)$$

where

$$U(p, z, r, \varphi) = v \cdot \left[\int_0^{+\infty} \sum_{m=-\infty}^{\infty} \sum_{i=1}^3 V_m^{(i)}(z, \xi, p) A_m^{(i)} \xi d\xi \right]$$

with a solution similar to the stationary problem of elasticity theory, which is derived in the same way with conditions that $V_m^{(i)}$, (z, ξ, p) be square integrable on the interval $z \in (0, \infty)$. Here and later on, V indicates that the integration contour goes along the real axis, and around the poles of the integrand with small semicircles above the pole (for integration with ξ) and below the pole (with integration with p). Hence for the components of displacement along the directions a_z, a_r, a_φ we obtain:

$$\begin{aligned} u_z &= \frac{1}{2\pi} \int_{-\infty}^{+\infty} e^{ipt} \bar{v} \cdot \int_0^{\infty} \left[\sum_{m=-\infty}^{\infty} V_m^{(1)} Y_m \right] \xi d\xi dp, \\ u_r &= \frac{1}{2\pi} \int_{-\infty}^{+\infty} e^{ipt} \bar{v} \cdot \int_0^{\infty} \left[\sum_{m=-\infty}^{\infty} \left(V_m^{(2)} \frac{\partial Y_m}{\partial r} + \frac{V_m^{(3)}}{r} \frac{\partial Y_m}{\partial \varphi} \right) \right] d\xi dp, \quad (1.10) \\ u_\varphi &= \frac{1}{2\pi} \int_{-\infty}^{+\infty} e^{ipt} \bar{v} \cdot \int_0^{\infty} \left[\sum_{m=-\infty}^{\infty} \left(\frac{V_m^{(2)}}{r} \frac{\partial Y_m}{\partial \varphi} - V_m^{(3)} \frac{\partial Y_m}{\partial r} \right) \right] d\xi dp. \end{aligned}$$

Placing (1.8), (1.10) in equation (1.1) and in the boundary condition (1.4), and accepting the permissibility of moving the double differentiation under the integral sign, we obtain the following equations:

1. For $V_m^{(1)}, V_m^{(2)}$:

$$\begin{aligned} l_1(V_m^{(1)}, V_m^{(2)}) &\equiv \frac{d}{dz} \left[(\lambda + 2\mu) \frac{dV_m^{(1)}}{dz} - \xi \lambda V_m^{(2)} \right] - \\ &- \xi \mu \frac{dV_m^{(2)}}{dz} + V_m^{(1)} (p^2 \rho - \xi^2 \mu) = - f_m^{(1)}, \quad (1.11) \end{aligned}$$

$$l_2(V_m^{(1)}, V_m^{(2)}) \equiv \frac{d}{dz} \left[\mu \frac{dV_m^{(2)}}{dz} + \xi \mu V_m^{(1)} \right] +$$

$$\text{with boundary conditions } + \xi \lambda \frac{dV_m^{(1)}}{dz} + V_m^{(2)} (p^2 \rho - \xi^2 \lambda - 2 \xi^2 \mu) = - f_m^{(2)},$$

$$\begin{aligned} \sigma_{zz} &\equiv (\lambda + 2\mu) \frac{dV_m^{(1)}}{dz} - \xi \lambda V_m^{(2)} = 0, \quad \tau_{rz} \equiv \mu \frac{dV_m^{(2)}}{dz} + \xi \mu V_m^{(1)} = 0 \\ \text{при } z &= 0. \end{aligned} \quad (1.12)$$

The functions $v_m^{(1)}$, $v_m^{(2)}$, σ_{zz} , and τ_{rz} are continuous and bounded for all z .

2. For $v_m^{(3)}$

$$l_3(V_m^{(3)}) \equiv \frac{d}{dz} \left(\mu \frac{dV_m^{(3)}}{dz} \right) + V_m^{(3)} (p^2 \rho - \xi^2 \mu) = -f_m^{(3)}, \quad (1.13)$$

with boundary condition

$$\tau_{\phi z} \equiv \mu \frac{dV_m^{(3)}}{dz} = 0 \quad \text{up to } z = 0. \quad (1.14)$$

$v_m^{(3)}$ and $\tau_{\phi z}$ are continuous and bounded for all z . The left side of equation (1.11) and the boundary condition (1.12) define a self-adjoint operator L in the region of integration with the square integrable vector function

$$\begin{vmatrix} v_m^{(1)} \\ v_m^{(2)} \end{vmatrix} \quad \begin{array}{l} \text{The left side of (1.13) and the boundary condition (1.14) defines} \\ \text{a self-adjoint operator } L_3 \text{ in the region of integration with the} \\ \text{function } v_m^{(3)}. \text{ If the vector function } \begin{vmatrix} f_m^{(1)} \\ f_m^{(2)} \end{vmatrix} \text{ and the function} \end{array}$$

$f_m^{(3)}$ are also square integrable on the interval $z \in (0, \infty)$, which follows from the conditions imposed on the function $F(t, z, r, \phi)$, then the following relations between the functions $v_m^{(1)}$, $v_m^{(2)}$, $v_m^{(3)}$ and the eigenfunctions of the above described operators are valid (2, 6, 10).

Expression of V_m in terms of eigenfunctions. $v_m^{(i)}$ ($i = 1, 2$) can be expressed in the following manner:

$$V_m^{(i)} = \sum_{k=1}^{k_R(\xi)} c_{km}^R \tilde{V}_k^{(i)} + \int_{p_0^2}^{\infty} c_m^R(\beta) \tilde{V}^{(i)}(\beta, z) d\beta. \quad (1.15)$$

Here, for the coefficients c_{km}^R , c_m^R we have:

$$c_{km}^R = \frac{1}{p_{kR}^2(\xi) - p^2} \frac{D_{km}^R}{I_{kR}}, \quad c_m^R = \frac{D_m^R}{\beta(\xi) - p^2}, \quad (1.16)$$

$$D_{km}^R = \int_0^{\infty} (f_m^{(1)} \tilde{V}_k^{(1)} + f_m^{(2)} \tilde{V}_k^{(2)}) dz, \quad D_m^R = \int_0^{\infty} (f_m^{(1)} \tilde{V}^{(1)} + f_m^{(2)} \tilde{V}^{(2)}) dz,$$

$$I_{kR} = \int_0^{\infty} \rho [(\tilde{V}_k^{(1)})^2 + (\tilde{V}_k^{(2)})^2] dz, \quad p_0^2 = \xi^2 b^2 (Z + 0).$$

Here $\tilde{V}_k^{(1)}$ and $\tilde{V}_k^{(2)}$, $\tilde{V}^{(1)}$ and $\tilde{V}^{(2)}$ are the eigenfunctions of the operators constituting the left-hand sides of (1.11) with the boundary conditions (1.12).

The first part of the function corresponds to the discrete spectra of eigenvalues p_{kR}^2 ($k = 1, 2, \dots, k(\xi)$); $\xi^2 v_R^2 < p_{kR}^2 < p_0^2$, where v_R is the minimal velocity of Rayleigh waves in a halfspace with constants equal $a(z)$, $b(z)$, $\rho(z)$ for some value of depth (1,2). The second part corresponds to the continuous spectra of eigenvalues $\beta (p_0^2 \leq \beta < \infty)$. Here the wave number ξ plays the role of a free parameter. Formula (1.16) is obtained using the orthogonality relations:

$$\int_0^\infty \rho [V_i^{(1)} V_j^{(1)} + V_i^{(2)} V_j^{(2)}] dz = 0 \quad \text{при } i \neq j,$$

$$\int_0^\infty \rho [V^{(1)}(\beta) \tilde{V}(p^2) + V^{(2)}(\beta) \tilde{V}^{(2)}(p^2)] dz = \delta(p^2 - \beta),$$

where $\tilde{V}^{(i)}$ is the complex conjugate of $V^{(i)}$.

Similarly $V_m^{(3)}$ can be expressed as

$$V_m^{(3)} = \sum_{k=1}^{k_L(\xi)} c_{km}^L V_k^{(3)} + \int_{p_0^2}^\infty c_m^L(\beta) V^{(3)}(z, \beta) d\beta, \quad (1.17)$$

where

$$c_{km}^L = \frac{1}{p_{kL}^2(\xi) - p^2} \frac{D_{km}^L}{I_{kL}}, \quad c_m^L = \frac{D_m^L}{\beta(\xi) - p^2};$$

$$D_{km}^L = \int_0^\infty f_m^{(3)} V_k^{(3)} dz, \quad D_m^L = \int_0^\infty f_m^{(3)} \tilde{V}^{(3)} dz, \quad (1.18)$$

$$I_{kL} = \int_0^\infty \rho (V_k^{(3)})^2 dz.$$

For this derivation the following orthogonality conditions were used:

$$\int_0^\infty \rho V_i^{(3)} V_j^{(3)} dz = 0 \quad \text{при } i \neq j,$$

$$\int_0^\infty \rho V^{(3)}(p^2) \tilde{V}^{(3)}(\beta) dz = \delta(p^2 - \beta).$$

Placing these expressions for V_m into (1.10), we obtain the complete precise formulae for displacements.

$$u_z = \frac{1}{2\pi} \int_{-\infty}^{+\infty} e^{ipt} \left\{ v \cdot \int_0^{\infty} \left[\sum_{m=-\infty}^{\infty} \left(\sum_{k=1}^{K_R(\xi)} c_{km}^R V_k^{(1)} + \int_{p_0^2}^{\infty} c_m^R(\beta) V^{(1)} d\beta \right) Y_m \right] d\xi \right\} dp,$$

$$\begin{aligned} u_r = & \frac{1}{2\pi} \int_{-\infty}^{+\infty} e^{ipt} \left\{ v \cdot \int_0^{\infty} \left[\sum_{m=-\infty}^{\infty} \left(\left(\sum_{k=1}^{K_R(\xi)} c_{km}^R V_k^{(2)} + \int_{p_0^2}^{\infty} c_m^R(\beta) V^{(2)} d\beta \right) \frac{\partial Y_m}{\partial r} + \right. \right. \right. \right. \\ & \left. \left. \left. \left. + \frac{1}{r} \left(\sum_{k=1}^{K_L(\xi)} c_{km}^L V_k^{(3)} + \int_{p_0^2}^{\infty} c_m^L(\beta) V^{(3)} d\beta \right) \frac{\partial Y_m}{\partial \varphi} \right) \right] d\xi \right\} dp, \quad (1.19) \end{aligned}$$

$$\begin{aligned} u_\varphi = & -\frac{1}{2\pi} \int_{-\infty}^{+\infty} e^{ipt} \left\{ v \cdot \int_0^{\infty} \left[\sum_{m=-\infty}^{\infty} \left(\frac{1}{r} \left(\sum_{k=1}^{K_R(\xi)} c_{km}^R V_k^{(2)} + \int_{p_0^2}^{\infty} c_m^R(\beta) V^{(2)} d\beta \right) \times \right. \right. \right. \right. \\ & \left. \left. \left. \left. \times \frac{\partial Y_m}{\partial \varphi} - \left(\sum_{k=1}^{K_L(\xi)} c_{km}^L V_k^{(3)} + \int_{p_0^2}^{\infty} c_m^L(\beta) V^{(3)} d\beta \right) \frac{\partial Y_m}{\partial r} \right) \right] d\xi \right\} dp. \end{aligned}$$

It can be shown that the derived solution for non-stationary problem satisfies the null initial conditions (1.4a). Thus for fixed ξ the functions of p in (1.19), multiplied by $\exp(ip t)$, are only the coefficients c_{km}^Q , c_m^Q which are analytic everywhere except at a finite number of points for $\text{Im } p \geq 0$, decreasing as $p \rightarrow \infty$ not slower than $O(1/p^{1+|\xi|})$. Changing the order of integration among p and ξ taking the integration contour with respect to p in the lower halfspace, we obtain the null condition on $u(t)$ and du/dt for $t < 0$.

Asymptotic expressions for large r . At large distances r , which are not commensurable with the length dimensions of the seismic source, the main part of the disturbance given by formula (1.19) becomes the Rayleigh and Love surface waves. Their contribution equals the sum of residues of the poles in the integrand, contained in the sum $\sum_{k=1}^{K_Q(\xi)}$.

Keeping only parts, decreasing not faster than r^{-1} , we obtain the following asymptotic formula for displacements at large r (5):

$$\begin{aligned}
u_z(t, z, r, \varphi) &= \frac{1}{\sqrt{2\pi r}} \operatorname{Re} \int_{\frac{-\infty}{p}}^{\infty} \frac{\exp i(p t - \frac{\pi}{4})}{p} \times \\
&\quad \times \left[\sum_{k=1}^{K_R(p)} U_{kR}(p, \varphi) \tilde{V}_k^{(1)}(p, z) \frac{\sqrt{\xi_{kR}}}{C_{kR} I_{kR}} \exp(-i\xi_{kR} r) \right] dp, \\
u_r(t, z, r, \varphi) &= \frac{1}{\sqrt{2\pi r}} \operatorname{Re} \int_{\frac{-\infty}{p}}^{\infty} \frac{\exp i(p t - \frac{3\pi}{4})}{p} \times \\
&\quad \times \left[\sum_{k=1}^{K_R(p)} U_{kR}(p, \varphi) \tilde{V}_k^{(2)}(p, z) \frac{\sqrt{\xi_{kR}}}{C_{kR} I_{kR}} \exp(-i\xi_{kR} r) \right] dp, \quad (1.20) \\
u_\varphi(t, z, r, \varphi) &= \frac{1}{\sqrt{2\pi r}} \operatorname{Re} \int_{\frac{-\infty}{p}}^{\infty} \frac{\exp i(p t + \frac{\pi}{4})}{p} \times \\
&\quad \times \left[\sum_{k=1}^{K_L(p)} U_{kL}(p, \varphi) \tilde{V}_k^{(3)}(p, z) \frac{\sqrt{\xi_{kL}}}{C_{kL} I_{kL}} \exp(-i\xi_{kL} r) \right] dp.
\end{aligned}$$

Here for Q = R, L $K_Q(p)$ is the maximum number of harmonics of Rayleigh (R) and Love (L) waves which exist for a given p.

$$U_{kQ}(p, \varphi) = \sum_{m=-\infty}^{\infty} D_{km}^Q \exp im\left(\varphi + \frac{\pi}{2}\right), \quad (1.21)$$

ξ_{kQ} is the wave number, a root of the equation $p_{kQ}^2(\xi) - p^2 = 0$. The phase and group velocities of the k'th harmonic v_{kQ} and C_{kQ} are related to ξ_{kQ} by

$$v_{kQ} = \frac{p}{\xi_{kQ}}, \quad C_{kQ} = \frac{dp_{kQ}}{d\xi} (\xi = \xi_{kQ}). \quad (1.22)$$

\bar{p} is the limiting frequency; oscillations with frequencies less than \bar{p} make up the quasi-static part of the disturbance and are not of interest to us.

For the derivation of formula (1.20) the following asymptotic relations were used:

$$v. \int_0^\infty \frac{\Phi(\xi) J_m(\xi r) \xi d\xi}{p_{kQ}^2(\xi) - p^2} \approx \sqrt{\frac{\pi \xi_{kQ}}{2r}} \frac{\Phi(\xi_{kQ})}{p C_{kQ}} \times \\ \times \exp \left\{ -i \left[\xi_{kQ} r - \left(m - \frac{1}{2} \right) \frac{\pi}{2} \right] \right\},$$

$$v. \int_0^\infty \frac{\Phi(\xi) \frac{dJ_m(\xi r)}{dr} d\xi}{p_{kQ}^2(\xi) - p^2} \approx \sqrt{\frac{\pi \xi_{kQ}}{2r}} \frac{\Phi(\xi_{kQ})}{p C_{kQ}} \times \\ \times \exp \left\{ -i \left[\xi_{kQ} r - \left(m - \frac{3}{2} \right) \frac{\pi}{2} \right] \right\}.$$

The vertical u_z and radial u_r components of displacement make up the Rayleigh wave; the tangential component u_ϕ makes up the Love wave.

Expressions for p_{kQ} , C_{kQ} by integrals of eigenfunctions. The calculus of variation can be used to obtain formulae for p_{kQ} and C_{kQ} (2, 15, 23):

$$p_{kQ}^2 = (\xi^2 G_{1k}^Q + 2\xi G_{2k}^Q + G_{3k}^Q)/I_{kQ}, \quad C_{kQ} = (\xi G_{1k}^Q + G_{2k}^Q)/p_{kQ} I_{kQ}, \quad (1.23)$$

where G_{jk}^Q are the following integrals:

$$G_{1k}^R = \int_0^\infty [(\lambda + 2\mu)(\tilde{V}_k^{(2)})^2 + \mu(\tilde{V}_k^{(1)})^2] dz, \\ G_{2k}^R = \int_0^\infty \left[\mu \frac{d\tilde{V}_k^{(2)}}{dz} \tilde{V}_k^{(1)} - \lambda \frac{d\tilde{V}_k^{(1)}}{dz} \tilde{V}_k^{(2)} \right] dz, \\ G_{3k}^R = \int_0^\infty \left[(\lambda + 2\mu) \left(\frac{d\tilde{V}_k^{(1)}}{dz} \right)^2 + \mu \left(\frac{d\tilde{V}_k^{(2)}}{dz} \right)^2 \right] dz, \quad (1.24)$$

$$G_{1k}^L = \int_0^\infty \mu (\tilde{V}_k^{(3)})^2 dz,$$

$$G_{2k}^L \equiv 0,$$

$$G_{3k}^L = \int_0^\infty \mu \left(\frac{d\tilde{V}_k^{(3)}}{dz} \right)^2 dz.$$

Methods for calculating p_{kR} , $\tilde{V}_k^{(1)}$ and $\tilde{V}_k^{(2)}$ are described in (11, 21), and the methods for p_{kL} , $\tilde{V}_k^{(3)}$ in (2, 15).

SURFACE WAVES FROM ELEMENTARY SOURCES. We will consider what form equation (1.20) takes for some elementary sources: axially symmetric vertical and radial impacts, torsional impact, field of horizontal forces, dipoles, center of compression. The fields of many more complex sources can be obtained by adding the fields of these elementary sources with respective constants of proportionality. Because the function U_{kQ} in equation (1.20) is the only term which is affected by the source, it is only necessary to consider the expressions for U_{kQ} for various impacts.

1. Vertical axially symmetric impact. Let

$$\mathbf{F} = F_z(t, z, r) \mathbf{a}_z.$$

In this case we get from (1.7):

$$f_0^{(1)}(z, \xi, p) = \int_{-\infty}^{+\infty} e^{-ipz} \int_0^{\infty} F_z J_0(\xi r) r dr dt;$$

$$f_m^{(1)} = 0 \quad \text{при } m \neq 0; \quad f_m^{(2)} \equiv f_m^{(3)} \equiv 0;$$

$$D_{km}^R = \int_0^{\infty} f_0^{(1)}(z, \xi_{kR}, p) \tilde{V}_k^{(1)}(z, \xi_{kR}) dz;$$

$$D_{km}^L = 0 \quad \text{при } m \neq 0; \quad D_{km}^L \equiv 0.$$

In summary,

$$U_{kR} = \int_0^{\infty} f_0^{(1)} \tilde{V}_k^{(1)} dz, \quad U_{kL} = 0. \quad (1.25)$$

In particular, for an ideally concentrated vertical force at the point $z = h$, $r = 0$, $\mathbf{F} = \delta(z - h) \frac{\delta(r)}{r} \Phi(t) \mathbf{a}_z$

$$U_{kR} = \tilde{V}_k^{(1)}(h, p) S(p). \quad (1.26)$$

Here and later on $S(p) = \int_{-\infty}^{\infty} \Phi(t) \exp(-ipt) dt$ — the source time spectrum.

II. Radial axially symmetric impact. Let

$$\mathbf{F} = F_r(t, z, r) \mathbf{a}_r.$$

In this case we obtain from (1.7):

* Translator's note: these concentrated forces are not normalized by the factor $1/2\pi$.

$$f_0^{(2)} = - \int_{-\infty}^{+\infty} e^{-ipr} \int_0^{\infty} F_r J_1(\xi r) r dr dt; \quad f_m^{(2)} = 0 \text{ при } m \neq 0;$$

$$f_m^{(1)} \equiv f_m^{(3)} \equiv 0;$$

$$D_{k0}^R = - \int_0^{\infty} f_0^{(2)} V_k^{(2)} dz; \quad D_{km}^R = 0 \text{ при } m \neq 0; D_{km}^L \equiv 0.$$

Hence,

$$U_{kR} = - \int_0^{\infty} f_0^{(2)} V_k^{(2)} dz, \quad U_{kL} = 0. \quad (1.27)$$

In the ideal case of a radial source concentrated at the point $z = h$, $r = 0$, $\mathbf{F} = 2\delta(z-h) \frac{\delta(r)}{r^2} \varphi(t) \mathbf{a}_r$

$$U_{kR} = - \xi_{kR} V_k^{(2)}(h, p) S(p). \quad (1.28)$$

III. Rotational impact. Let

$$\mathbf{F} = F_{\phi}(t, z, r) \mathbf{a}_{\phi}.$$

From (1.7) we get

$$f_0^{(3)} = \int_{-\infty}^{+\infty} e^{-ipr} \int_0^{\infty} F_{\phi} J_1(\xi r) r dr dt; \quad f_m^{(3)} = 0 \text{ при } m \neq 0;$$

$$f_m^{(1)} \equiv f_m^{(2)} \equiv 0; \quad D_{km}^R \equiv 0; \quad D_{k0}^L = \int_0^{\infty} f_0^{(3)} V_k^{(3)} dz;$$

$$D_{km}^L = 0 \text{ при } m \neq 0.$$

The final result is

$$U_{kR} = 0, \quad U_{kL} = \int_0^{\infty} f_0^{(3)} V_k^{(3)} dz. \quad (1.29)$$

In the case of a concentrated torque at $z = h$, $r = 0$ and

$$\mathbf{F} = 2\delta(z-h) \frac{\delta(r)}{r^2} \varphi(t) \mathbf{a}_{\phi}$$

$$U_{kL} = \xi_{kL} V_k^{(3)}(h, p) S(p). \quad (1.30)$$

IV. Field of a horizontal force with fixed orientation. We will consider only the particular case of this force field

$$\mathbf{F} = F_T(t, z, r) \mathbf{a}_T,$$

* Translator's note: these concentrated forces are not normalized by the factor $1/2\pi$

where \mathbf{a}_T is a horizontal unit vector of fixed orientation:

$$(\mathbf{a}_T, \mathbf{a}_z) = 0, (\mathbf{a}_T, \mathbf{a}_r) = \cos(\delta - \varphi),$$

$$(\mathbf{a}_T, \mathbf{a}_\theta) = \sin(\delta - \varphi).$$

Then,

$$\mathbf{F} = F_T [\cos(\delta - \varphi) \mathbf{a}_r + \sin(\delta - \varphi) \mathbf{a}_\theta]$$

and from (1.7) we have

$$\begin{aligned} f_m^{(1)} &\equiv 0; \quad f_m^{(2)} = f_m^{(3)} = 0 \quad \text{при } m \neq \pm 1; \\ f_1^{(2)} &= \frac{e^{-i\delta}}{2} f_T; \quad f_{-1}^{(2)} = -\frac{e^{-i\delta}}{2} f_T; \\ f_1^{(3)} &= \frac{e^{-i\delta}}{2i} f_T; \quad f_{-1}^{(3)} = -\frac{e^{i\delta}}{2i} f_T; \\ f_T &= \int_{-\infty}^{+\infty} e^{-ipr} \int_0^{\infty} F_T J_0(\xi r) r dr dt. \end{aligned}$$

Summing over m , we obtain the following relation:

$$U_{kR} = i \cos(\delta - \varphi) \int_0^{\infty} f_T V_k^{(2)} dz, \quad (1.31)$$

$$U_{kL} = -i \sin(\delta - \varphi) \int_0^{\infty} f_T V_k^{(3)} dz.$$

In the case of an ideal concentrated force acting at the point $r = 0$, $z = h$,

$$\mathbf{F} = \delta(z - h) \frac{\delta(r)}{r} \varphi(t) \mathbf{a}_T;$$

$$U_{kR} = i \cos(\delta - \varphi) V_k^{(2)}(h, p) S(p), \quad (1.32)$$

$$U_{kL} = -i \sin(\delta - \varphi) V_k^{(3)}(h, p) S(p).$$

V. An arbitrarily oriented concentrated force. Let

$$\mathbf{F} = \delta(z - h) \frac{\delta(r)}{r} [\mathbf{a}_z \cos \beta + \mathbf{a}_T \sin \beta] \varphi(t).$$

Combining (1.25) and (1.32), we get

$$U_{kR} = [\cos \beta V_k^{(1)}(h, p) + i \sin \beta \cos(\delta - \varphi) V_k^{(2)}(h, p)] S(p), \quad (1.33)$$

$$U_{kL} = -i \sin \beta \sin(\delta - \varphi) V_k^{(3)}(h, p) S(p).$$

VI. Dipole without moment. The field of a dipole is made up of a pair of forces without moment

$$\mathbf{F}_{\pm} = \pm \delta(z - h) \frac{\delta(r)}{r} \varphi(t) [\mathbf{a}_z \cos \beta + \mathbf{a}_T \sin \beta],$$

we obtain the result by using the following operator on U_{kQ} IN (1.33) (keeping only those parts decreasing slower than r^{-1}):

$$\left[\cos \beta \frac{d}{dh} + i \xi_{kQ} \cos(\delta - \varphi) \sin \beta \right].$$

We obtain the results:

$$U_{kR} = \left[\cos^2 \beta \frac{d\tilde{V}_k^{(1)}}{dh}(h_1 p) - \xi_{kR} \sin^2 \beta \cos^2(\delta - \varphi) \tilde{V}_k^{(3)}(h_1 p) + \right. \\ \left. + \frac{i}{2} \sin 2\beta \left(\cos(\delta - \varphi) \xi_{kR} \tilde{V}_k^{(1)}(h_1 p) + \frac{d\tilde{V}_k^{(2)}}{dh}(h_1 p) \right) \right] S(p), \quad (1.34)$$

$$U_{kL} = -i \sin \beta \sin(\delta - \varphi) \left[\cos \beta \frac{d\tilde{V}_k^{(3)}}{dh}(h_1 p) + \right. \\ \left. + i \xi_{kL} \cos(\delta - \varphi) \sin \beta \tilde{V}_k^{(3)}(h_1 p) \right] S(p).$$

VII. Dipole with moment. Let a pair of forces act in the same direction as in VI, but allow the force to have a moment; the axis of the dipole, i.e., the line associated with the points where the force acts, is given by the vector a_n , where

$$(a_n, a_z) = \cos \gamma, (a_n, a_r) = \sin \gamma \cos(\alpha - \varphi);$$

$$(a_n, a_\varphi) = \sin \gamma \sin(\alpha - \varphi).$$

The angles γ , δ , β , α are not independent, and are related by the relation $\operatorname{ctg} \gamma \operatorname{ctg} \beta = -\cos(\delta - \alpha)$. The displacement field of this dipole valid for large distances r is obtained by using the following operator on U_{kQ} in (1.33):

$$\left[\cos \gamma \frac{d}{dh} + i \xi_{kQ} \cos(\alpha - \varphi) \sin \gamma \right].$$

$$U_{kR} = \left[\cos \gamma \cos \beta \frac{d\tilde{V}_k^{(1)}}{dh}(h_1 p) - \xi_{kR} \sin \beta \sin \gamma \cos(\delta - \varphi) \cos(\alpha - \varphi) \times \right. \\ \times \tilde{V}_k^{(3)}(h_1 p) + i \xi_{kR} \sin \gamma \cos \beta \cos(\alpha - \varphi) \tilde{V}_k^{(1)}(h_1 p) + \\ \left. + i \cos \gamma \sin \beta \cos(\delta - \varphi) \frac{d\tilde{V}_k^{(2)}}{dh}(h_1 p) \right] \times S(p),$$

$$U_{kL} = -i \sin \beta \sin(\delta - \varphi) \left[\cos \gamma \frac{d\tilde{V}_k^{(3)}}{dh}(h_1 p) + i \xi_{kL} \sin \gamma \cos(\alpha - \varphi) \times \right. \\ \left. \times \tilde{V}_k^{(3)}(h_1 p) \right] \times S(p).$$

VIII. Center of expansion. The center of expansion can be considered to be equivalent to a source composed of three orthogonal dipoles without moment. We will take one dipole to be vertical ($\beta = 0$), and the other two horizontal ($\beta = \pi/2$, $\delta = 0$ and $\pi/2$). The resultant field is

$$U_{kR} = \left[\frac{d\tilde{V}_k^{(1)}}{dh}(h_1 p) - \xi_{kR} \tilde{V}_k^{(2)} \right] \times S(p), \quad U_{kL} = 0. \quad (1.35)$$

2. Displacement fields in an elastic sphere

Statement of the problem. We will consider an elastic sphere with

coordinates R, θ, φ ($0 \leq R \leq R_0, 0 \leq \theta \leq \pi, 0 \leq \varphi < 2\pi$). motion in this coordinate system are (7):

$$\frac{1}{R} \frac{\partial \hat{R}}{\partial \theta} + \frac{1}{R \sin \theta} \frac{\partial \hat{\theta} R}{\partial \varphi} + \frac{\partial \hat{R} R}{\partial R} + \frac{1}{R} (2 \hat{R} R - \hat{\theta} \hat{\theta} - \hat{\varphi} \hat{\varphi} + \hat{\theta} \hat{R} \times \\ \times \operatorname{ctg} \theta) = \rho \frac{\partial^2 u_R}{\partial t^2} - F_R,$$

$$\frac{1}{R} \frac{\partial \hat{\theta}}{\partial \theta} + \frac{1}{R \sin \theta} \frac{\partial \hat{\theta} \hat{\varphi}}{\partial \varphi} + \frac{\partial \hat{\theta} R}{\partial R} + \frac{1}{R} [3 \hat{\theta} \hat{R} + (\hat{\theta} \hat{\theta} - \hat{\varphi} \hat{\varphi}) \times \\ \times \operatorname{ctg} \theta] = \rho \frac{\partial^2 u_\theta}{\partial t^2} - F_\theta, \quad (2.1)$$

$$\frac{1}{R} \frac{\partial \hat{\varphi}}{\partial \theta} + \frac{1}{R \sin \theta} \frac{\partial \hat{\varphi} \hat{\varphi}}{\partial \varphi} + \frac{\partial \hat{\varphi} R}{\partial R} + \\ + \frac{1}{R} [3 \hat{\varphi} \hat{R} + 2 \hat{\theta} \hat{\varphi} \operatorname{ctg} \theta] = \rho \frac{\partial^2 u_\varphi}{\partial t^2} - F_\varphi.$$

Here $\hat{R} R, \hat{\theta} \hat{\theta}, \hat{\varphi} \hat{\varphi}, \hat{R} \hat{R}$ are components of the stress tensor; u_R, u_θ, u_φ are components of the displacement vector u along the directions a_R, a_θ, a_φ ; F_R, F_θ, F_φ are components of the vector volume force acting at the source along the same directions.

For the components of stress we have the following relations

$$\begin{aligned} \hat{R} R &= \mu \left(\frac{\partial u_\theta}{\partial R} - \frac{u_\theta}{R} + \frac{1}{R} \frac{\partial u_R}{\partial \theta} \right), \\ \hat{\theta} \hat{\theta} &= \lambda \Delta + 2 \frac{\mu}{R} \left(\frac{\partial u_\theta}{\partial \theta} + u_R \right), \\ \hat{\varphi} \hat{\varphi} &= \frac{\mu}{R} \left(\frac{\partial u_\varphi}{\partial \theta} - u_\varphi \operatorname{ctg} \theta + \frac{1}{\sin \theta} \frac{\partial u_\theta}{\partial \varphi} \right), \\ \hat{\varphi} R &= \mu \left(\frac{\partial u_\varphi}{\partial R} - \frac{u_\varphi}{R} + \frac{1}{R \sin \theta} \frac{\partial u_R}{\partial \varphi} \right), \\ \hat{\varphi} \hat{\varphi} &= \lambda \Delta + 2 \frac{\mu}{R} \left(u_R + u_\theta \operatorname{ctg} \theta + \frac{1}{\sin \theta} \frac{\partial u_\varphi}{\partial \varphi} \right), \\ \hat{R} \hat{R} &= \lambda \Delta + 2 \mu \frac{\partial u_R}{\partial R}. \end{aligned} \quad (2.2)$$

Here the dilatation has the form

$$\Delta = \frac{\partial u_R}{\partial R} + \frac{2 u_R}{R} + \frac{1}{R \sin \theta} \frac{\partial u_\varphi}{\partial \varphi} + \frac{1}{R} \frac{\partial u_\theta}{\partial \theta} + \frac{u_\theta}{R} \operatorname{ctg} \theta. \quad (2.3)$$

The Lamé constants λ, μ and the density ρ are piecewise continuous positive functions of a single coordinate R ; the components of displacement and stress are continuous and bounded for all points in $0, R_0$. The surface of the sphere is free from stress, i.e.

$$\hat{\theta} \hat{R} = \hat{\varphi} \hat{R} = \hat{R} \hat{R} = 0 \text{ при } R = R_0. \quad (2.4)$$

The initial conditions are

$$u = \frac{\partial u}{\partial t} = 0 \quad \text{при } t < 0. \quad (2.4a)$$

The force field $F(t, R, \theta, \phi)$ is described as a real source localized in space and time. The following limitations are laid upon F :

- 1) $F(t, R, \theta, \phi) = 0$ при $t < 0$;
- 2) $F(t, R, \theta, \phi)$

is absolutely integrable with respect to t and satisfies the Dirichlet conditions for all arguments. Then we are permitted the following representation:

$$F = \frac{1}{2\pi} \int_{-\infty}^{+\infty} e^{ipl} \left[\sum_{n=0}^{\infty} \sum_{m=-n}^n \sum_{i=1}^3 f_{mn}^{(i)}(R, p) A_{mn} \right] dp, \quad (2.5)$$

where

$$\begin{aligned} A_{mn}^{(1)} &= a_R Y_{mn}, \\ A_{mn}^{(2)} &= \left(a_\theta \frac{\partial Y_{mn}}{\partial \theta} + a_\phi \frac{1}{\sin \theta} \frac{\partial Y_{mn}}{\partial \phi} \right) \frac{1}{N}, \\ A_{mn}^{(3)} &= \left(a_\theta \frac{\partial Y_{mn}}{\partial \phi} \frac{1}{\sin \theta} - a_\phi \frac{\partial Y_{mn}}{\partial \theta} \right) \frac{1}{N}, \\ Y_{mn}(\theta, \phi) &= e^{im\phi} P_n^m(\cos \theta), \\ N &= \sqrt{n(n+1)}. \end{aligned} \quad (2.6)$$

$P_n^m(x)$ is the associated Legendre polynomial, defined by the following formulae (8):

$$P_n^m(x) = \frac{(1-x^2)^{m/2}}{2^n \cdot n!} \frac{d^{n+m}(x^2-1)^n}{dx^{n+m}} \quad \text{при } m \geq 0,$$

$$P_n^{[m]}(x) = (-1)^m \frac{(n-|m|)!}{(n+|m|)!} P_n^{|m|}(x) \quad \text{при } m < 0.$$

The system of spherical vector functions $A_{mn}^{(i)}$ are a field of a system of vectors satisfying the following orthogonality conditions on a unit sphere:

$$\int_0^{2\pi} \int_0^\pi (A_{mn}^{(i)}, \bar{A}_{lq}^{(j)}) \sin \theta d\theta d\phi = 4\pi \delta_{ij} \delta_{ml} \delta_{nq} \frac{(n+m)!}{(n-m)!(2n+1)}.$$

¹ For $n = 0$ $N = 0$, $Y_{mn} = \text{const}$ and (2.6) loses meaning, we will consider that $A_{00}^{(2)} = A_{00}^{(3)} = 0$.

Another system of vector functions was suggested for the solution of a similar problem in elastic wave theory in the work of G. I. Petrashin (11), which are linear combinations of the system $A_{mn}^{(1)}$. In distinction with (11) where the wave field in the sphere is given in the form of a summation of potentials and solenoidal fields, the analysis of the wave fields according to the system $A_{mn}^{(1)}$ permits a separate study of the spheroidal and torsional oscillations of a sphere (22). All proofs on the orthogonality and completeness of the system, studied in (11) are easily applied to the system $A_{mn}^{(1)}$.

The coefficients $f_{mn}^{(i)}$ making up the impact F in the system $A_{mn}^{(i)}$ are

$$f_{mn}^{(i)}(R, p) = \frac{2n+1}{4\pi} \frac{(n-m)!}{(n+m)!} \int_{-\infty}^{+\infty} e^{-ipt} \iint_{0^0}^{2\pi\pi} (F, \bar{A}_{mn}^{(i)}) \sin \theta d\theta d\varphi dt.$$

Specifically for $i = 1, 2, 3$ we have:

$$\begin{aligned} f_{mn}^{(1)} &= \frac{(n-m)!}{(n+m)!} \frac{2n+1}{4\pi} \int_{-\infty}^{+\infty} e^{-ipt} \iint_{0^0}^{2\pi\pi} F_R \bar{Y}_{mn} \sin \theta d\theta d\varphi dt, \\ f_{mn}^{(2)} &= \frac{(n-m)!}{(n+m)!} \frac{2n+1}{4\pi N} \int_{-\infty}^{+\infty} e^{-ipt} \iint_{0^0}^{2\pi\pi} \left[F_\theta \frac{\partial \bar{Y}_{mn}}{\partial \theta} + F_\varphi \frac{1}{\sin \theta} \frac{\partial \bar{Y}_{mn}}{\partial \varphi} \right] \times \\ &\quad \times \sin \theta d\theta d\varphi dt, \\ f_{mn}^{(3)} &= \frac{(n-m)!}{(n+m)!} \frac{2n+1}{4\pi N} \int_{-\infty}^{+\infty} e^{-ipt} \iint_{0^0}^{2\pi\pi} \left[F_\theta \frac{1}{\sin \theta} \frac{\partial \bar{Y}_{mn}}{\partial \varphi} - F_\varphi \frac{\partial \bar{Y}_{mn}}{\partial \theta} \right] \times \\ &\quad \times \sin \theta d\theta d\varphi dt, \\ \bar{Y}_{mn}(\theta, \varphi) &= e^{-im\varphi} p_n^m (\cos \theta). \end{aligned} \quad (2.7)$$

For the force components F_R , F_θ , F_φ we obtain from (2.5), (2.6):

$$\begin{aligned} F_R &= \frac{1}{2\pi} \int_{-\infty}^{+\infty} e^{ipt} \left[\sum_{n=0}^{\infty} \sum_{m=-n}^n f_{mn}^{(1)} Y_{mn} \right] dp, \\ F_\theta &= \frac{1}{2\pi} \int_{-\infty}^{+\infty} e^{ipt} \left[\sum_{n=1}^{\infty} \frac{1}{N} \sum_{m=-n}^n \left(f_{mn}^{(2)} \frac{\partial Y_{mn}}{\partial \theta} + f_{mn}^{(3)} \frac{1}{\sin \theta} \frac{\partial Y_{mn}}{\partial \varphi} \right) \right] dp, \\ F_\varphi &= \frac{1}{2\pi} \int_{-\infty}^{+\infty} e^{ipt} \left[\sum_{n=1}^{\infty} \frac{1}{N} \sum_{m=-n}^n \left(f_{mn}^{(2)} \frac{1}{\sin \theta} \frac{\partial Y_{mn}}{\partial \varphi} - f_{mn}^{(3)} \frac{\partial Y_{mn}}{\partial \theta} \right) \right] dp. \end{aligned} \quad (2.8)$$

Formulae for displacements. We will seek displacements in the form

$$u(t, R, \theta, \varphi) = \frac{1}{2\pi} v \cdot \int_{-\infty}^{+\infty} e^{ipt} \left[\sum_{n=0}^{\infty} \sum_{m=-n}^n \sum_{i=1}^3 V_{mn}^{(i)} A_{mn}^{(i)} \right] dp, \quad (2.9)$$

where $v_{mn}^{(i)} = v_{mn}^{(i)}(R, p)$. Hence for the projection of displacements along the directions a_R, a_θ, a_ϕ we obtain:

$$\begin{aligned} u_R &= \frac{1}{2\pi} v \cdot \int_{-\infty}^{+\infty} e^{ipR} \left[\sum_{n=0}^{\infty} \sum_{m=-n}^n V_{mn}^{(1)} Y_{mn} \right] dp, \\ u_\theta &= \frac{1}{2\pi} v \cdot \int_{-\infty}^{+\infty} e^{ipR} \left[\sum_{n=1}^{\infty} \frac{1}{N} \sum_{m=-n}^n \left(V_{mn}^{(2)} \frac{\partial Y_{mn}}{\partial \theta} + V_{mn}^{(3)} \frac{1}{\sin \theta} \frac{\partial Y_{mn}}{\partial \varphi} \right) \right] dp, \\ u_\phi &= \frac{1}{2\pi} v \cdot \int_{-\infty}^{+\infty} e^{ipR} \left[\sum_{n=1}^{\infty} \frac{1}{N} \sum_{m=-n}^n \left(V_{mn}^{(2)} \frac{1}{\sin \theta} \frac{\partial Y_{mn}}{\partial \varphi} - V_{mn}^{(3)} \frac{\partial Y_{mn}}{\partial \theta} \right) \right] dp. \end{aligned} \quad (2.10)$$

Placing (2.8), (2.10) into equation (2.1) and the boundary conditions, we arrive at the following equation for $V_{mn}^{(1)}$.

1. For $V_{mn}^{(1)}, V_{mn}^{(2)}$:

$$\begin{aligned} l_1(V_{mn}^{(1)}, V_{mn}^{(2)}) &\equiv \frac{d}{dR} \left[(\lambda + 2\mu) \frac{dV_{mn}^{(1)}}{dR} + \frac{2\lambda}{R} V_{mn}^{(1)} - \frac{\lambda N}{R} V_{mn}^{(2)} \right] + \\ &\quad + \frac{\mu}{R^2} \left[4 \frac{dV_{mn}^{(1)}}{dR} R - 4V_{mn}^{(1)} + \right. \\ &\quad \left. + N \left(3V_{mn}^{(2)} - R \frac{dV_{mn}^{(2)}}{dR} - NV_{mn}^{(1)} \right) \right] + p^2 \rho V_{mn}^{(1)} = -f_{mn}^{(1)}, \\ l_2(V_{mn}^{(1)}, V_{mn}^{(2)}) &\equiv \frac{d}{dR} \left[\mu \left(\frac{dV_{mn}^{(2)}}{dR} - \frac{V_{mn}^{(2)}}{R} + \frac{NV_{mn}^{(1)}}{R} \right) \right] + \\ &\quad + \frac{\lambda N}{R} \left(\frac{dV_{mn}^{(1)}}{dR} + \frac{2}{R} V_{mn}^{(1)} - \frac{N}{R} V_{mn}^{(2)} \right) + \\ &\quad + \frac{\mu}{R^2} \left(5NV_{mn}^{(1)} + 3R \frac{dV_{mn}^{(2)}}{dR} - V_{mn}^{(2)} - 2N^2 V_{mn}^{(2)} \right) + p^2 \rho V_{mn}^{(2)} = -f_{mn}^{(2)} \end{aligned} \quad (2.11)$$

with the boundary conditions:

$$\begin{aligned} \sigma_{RR} &\equiv (\lambda + 2\mu) \frac{dV_{mn}^{(1)}}{dR} + \frac{2\lambda}{R} V_{mn}^{(1)} - \frac{\lambda N}{R} V_{mn}^{(2)} = 0, \\ \tau_{\theta R} &\equiv \mu \left(\frac{dV_{mn}^{(2)}}{dR} - \frac{V_{mn}^{(2)}}{R} + \frac{NV_{mn}^{(1)}}{R} \right) = 0 \quad \text{при } R = R_0, \quad (2.12) \\ V_{mn}^{(1)} &= V_{mn}^{(2)} = 0 \quad \text{при } R = 0. \end{aligned}$$

The functions $V_{mn}^{(1)}, V_{mn}^{(2)}, \sigma_{RR}$ and $\tau_{\theta R}$ are continuous and bounded everywhere on $0, R_0$.

2. For $V_{mn}^{(3)}$ we have

$$\begin{aligned} l_3(V_{mn}^{(3)}) &\equiv \frac{d}{dR} \left[\mu \left(\frac{dV_{mn}^{(3)}}{dR} - \frac{V_{mn}^{(3)}}{R} \right) \right] + \frac{3\mu}{R} \frac{dV_{mn}^{(3)}}{dR} - \\ &\quad - \frac{\mu}{R^2} (N^2 + 1) V_{mn}^{(3)} + p^2 \rho V_{mn}^{(3)} = -f_{mn}^{(3)} \end{aligned} \quad (2.13)$$

and the boundary condition:

$$\begin{aligned}\tau_{\phi R} &\equiv \mu \left(\frac{dV_{mn}^{(3)}}{dR} - \frac{V_{mn}^{(3)}}{R} \right) = 0 \quad \text{при } R = R_0, \\ V_{mn}^{(3)} &= 0 \quad \text{при } R = 0.\end{aligned}\quad (2.14)$$

The functions $V_{mn}^{(3)}$, $\tau_{\phi R}$ continuous and bounded everywhere on $[0, R_0]$.

Expressions for V_{mn} in terms of eigenfunctions. $V_{mn}^{(i)}$ ($i = 1, 2$) can be expressed in the following manner:

$$V_{mn}^{(i)} = \frac{1}{R_0^2} \sum_{k=1}^{\infty} c_{kmn}^S \tilde{V}_{kn}^{(i)}, \quad (2.15)$$

where for the coefficients c_{kmn}^S we have:

$$\begin{aligned}c_{kmn}^S &= \frac{1}{p_{knS}^2 - p^2} \frac{D_{kmn}^S}{I_{knS}}, \\ D_{kmn}^S &= \int_0^{R_0} (f_{mn}^{(1)} \tilde{V}_{kn}^{(1)} + f_{mn}^{(2)} \tilde{V}_{kn}^{(2)}) R^2 dR, \\ I_{knS} &= \frac{1}{R_0^2} \int_0^{R_0} \rho R^2 [(\tilde{V}_{kn}^{(1)})^2 + (\tilde{V}_{kn}^{(2)})^2] dR.\end{aligned}\quad (2.16)$$

Here $\tilde{V}_{kn}^{(i)}$ (R), $i = 1, 2$ is the eigenfunction, p_{knS}^2 is the eigenvalue of the operator consisting of the left hand side of (2.11) together with the boundary conditions (2.12), for a given integer value of the parameter n .

Similarly

$$V_{mn}^{(3)} = \frac{1}{R_0^2} \sum_{k=1}^{\infty} c_{kmn}^T \tilde{V}_{kn}^{(3)}, \quad (2.17)$$

where

$$\begin{aligned}c_{kmn}^T &= \frac{1}{p_{knT}^2 - p^2} \frac{D_{kmn}^T}{I_{knT}}, \\ D_{kmn}^T &= \int_0^{R_0} f_{mn}^{(3)} \tilde{V}_{kn}^{(3)} R^2 dR, \\ I_{knT} &= \frac{1}{R_0^2} \int_0^{R_0} \rho R^2 (\tilde{V}_{kn}^{(3)})^2 dR.\end{aligned}\quad (2.18)$$

where $\tilde{V}_{kn}^{(3)}$ is the eigenfunction, a_{knT}^2 is the eigenvalue of the operator consisting of the left-hand side of (2.13) and the boundary conditions (2.14).

To derive (2.15) - (2.18) these orthogonality conditions were used:

$$\int_0^{R_0} \rho R^2 (\tilde{V}_{kn}^{(1)} \tilde{V}_{ln}^{(1)} + \tilde{V}_{kn}^{(2)} \tilde{V}_{ln}^{(2)}) dR = 0,$$

$$\int_0^{R_0} \rho R^2 (\tilde{V}_{kn}^{(3)} \tilde{V}_{ln}^{(3)}) dR = 0 \quad \text{при } k \neq l.$$

Placing the expressions developed for $V_{mn}^{(i)}(R, p)$ into (2.10) the following formulae for displacements are obtained:

$$u_R = \frac{1}{2\pi R_0^2} v. \int_{-\infty}^{+\infty} e^{ipt} \left[\sum_{n=0}^{\infty} \sum_{m=-n}^n \sum_{k=1}^{\infty} c_{kmn}^S(p) \tilde{V}_{kn}^{(1)}(R) Y_{mn} \right] dp,$$

$$u_\theta = \frac{1}{2\pi R_0^2} v. \int_{-\infty}^{+\infty} e^{ipt} \left[\sum_{n=1}^{\infty} \frac{1}{N} \sum_{m=-n}^n \sum_{k=1}^{\infty} (c_{kmn}^S(p) \tilde{V}_{kn}^{(2)}(R) \frac{\partial Y_{mn}}{\partial \theta} + \right. \\ \left. + c_{kmn}^T(p) \tilde{V}_{kn}^{(3)}(R) \frac{1}{\sin \theta} \frac{\partial Y_{mn}}{\partial \varphi}) \right] dp, \quad (2.19)$$

$$u_\varphi = \frac{1}{2\pi R_0^2} v. \int_{-\infty}^{+\infty} e^{ipt} \left[\sum_{n=1}^{\infty} \frac{1}{N} \sum_{m=-n}^n \sum_{k=1}^{\infty} \left(c_{kmn}^S(p) \tilde{V}_{kn}^{(2)}(R) \frac{1}{\sin \theta} \frac{\partial Y_{mn}}{\partial \varphi} - \right. \right. \\ \left. \left. - c_{kmn}^T(p) \tilde{V}_{kn}^{(3)}(R) \frac{\partial Y_{mn}}{\partial \theta} \right) \right] dp.$$

The proof that this constructed solution (2.19) satisfies the initial conditions (2.4a) proceeds analogously with that of the plane case for each term in the series $\sum_{m,n}$.

Expression of displacements in the form of a sum of proper oscillations of a sphere. In formula (2.19) only products of the type $c_{mn}^Q \exp(ip\tau)$ depend on p . Thus calculation of integrals with respect to p reduces to evaluating a number of integrals of the type

$$v. \int_{-\infty}^{+\infty} \frac{e^{ipt} \psi(p)}{\alpha^2 - p^2} dp.$$

By introducing the conditions on the source, $\psi(p)$ can have singularities for $\operatorname{Im} p \geq 0$ and must be analytic in the lower halfspace. As $p \rightarrow \infty$, $\psi(p) \rightarrow 0$. Besides this $\psi(-p) = \overline{\psi(p)}$. Hence it is not difficult to obtain

$$\text{v. } \int_{-\infty}^{+\infty} e^{ipt} \frac{\psi(p)}{a^2 - p^2} dp = 2 \operatorname{Re} \left[\text{v.} \int_0^{\infty} e^{ipt} \frac{\psi(p)}{a^2 - p^2} dp \right] = \\ = \frac{2\pi}{a} \operatorname{Re} [\psi(a) e^{i(a t - \pi/2)}] + a(t).$$

The function $a(t)$ is associated with the particular source function $f_{mn}^{(i)}$ in the upper halfplane for $p = 0$. Its form, either exponential decay with time or static impact, is not of interest to us.

Placing this evaluation into (2.19) and ignoring $a(t)$ we get

$$u_R = \frac{1}{R_0^2} \operatorname{Re} \sum_{n=0}^{\infty} \sum_{m=-n}^n \sum_{k=1}^{\infty} \exp \left[i \left(p_{knS} t - \frac{\pi}{2} \right) \right] D_{kmn}^S \frac{\tilde{V}_{kn}^{(1)}(R)}{p_{knS} I_{knS}} \frac{Y_{mn}}{I_{knS}}, \\ u_{\theta} = \frac{1}{R_0^2} \operatorname{Re} \sum_{n=1}^{\infty} \frac{1}{N} \sum_{m=-n}^n \sum_{k=1}^{\infty} \left\{ \exp \left[i \left(p_{knS} t - \frac{\pi}{2} \right) \right] D_{kmn}^S \frac{\tilde{V}_{kn}^{(2)}(R)}{p_{knS} I_{knS}} \frac{\partial Y_{mn}}{\partial \theta} + \right. \\ \left. + \exp \left[i \left(p_{knT} t - \frac{\pi}{2} \right) \right] D_{kmn}^T \frac{\tilde{V}_{kn}^{(3)}(R)}{p_{knT} I_{knT} \sin \theta} \frac{\partial Y_{mn}}{\partial \varphi} \right\}, \\ u_{\phi} = \frac{1}{R_0^2} \operatorname{Re} \sum_{n=1}^{\infty} \frac{1}{N} \sum_{m=-n}^n \sum_{k=1}^{\infty} \left\{ \exp \left[i \left(p_{knS} t - \frac{\pi}{2} \right) \right] D_{kmn}^S \times \right. \\ \left. \times \frac{\tilde{V}_{kn}^{(2)}(R)}{I_{knS} p_{knS} \sin \theta} \frac{\partial Y_{mn}}{\partial \varphi} \right\} - \exp \left[i \left(p_{knT} t - \frac{\pi}{2} \right) \right] D_{kmn}^T \frac{\tilde{V}_{kn}^{(3)}(R)}{I_{knT} p_{knT}} \frac{\partial Y_{mn}}{\partial \theta}. \quad (2.19a)$$

Thus we have expressed the displacements in the form of proper S spheroidal and toroidal T oscillations of a sphere with discrete frequencies p_{knS} and p_{knT} . The radial component u_R is composed only of spheroidal oscillations; the components u_{θ} and u_{ϕ} are made up of both spheroidal and torsional components.

Asymptotic values of displacement for $n \sin \theta \gg |m|$. We will follow (13) for separating the propagating surface waves from the displacements described in (2.19a). We will change the order of summation

(i.e., change the sum $\sum_{n=0}^{\infty} \sum_{m=-n}^n \sum_{k=1}^{\infty}$ to $\sum_{k=1}^{\infty} \sum_{m=-\infty}^{\infty} \sum_{n=0}^{\infty}$) and

we will change the summation with respect to n to an integration along a contour L, enclosing positive part of the real axis in the plane of the complex variable v, where $n = \text{Integer} (\operatorname{Re} v)$:

$$\sum_{n=0}^{\infty} f_n P_n^m(\cos \theta) = -\frac{1}{2i} \oint_L \frac{f(v) P_v^m(\cos(\pi - \theta)) (-1)^m dv}{\cos \pi(v + \frac{1}{2})}.$$

Expressing $(\cos \pi(v+1/2))^{-1}$ for $\operatorname{Im} v \geq 0$ by the series

$$2 \sum_{l=0}^{\infty} (-1)^l \exp [i(2l+1)(v + \frac{1}{2})\pi], \quad \text{and for } \operatorname{Im} v \leq 0 \text{ the series}$$

2 $\sum_{l=0}^{\infty} (-1)^l \exp [-i(2l+1)(v + \frac{1}{2})\pi]$, fixing the contour of integration to the real axis and introducing a new variable $p = p_{kQ}(v)$ we have

$$\sum_{n=0}^{\infty} f_n P_n^m(\cos \theta) \approx 2 \sum_{l=0}^{\infty} (-1)^{l+|m|} \int_0^{\infty} f(p) P_l^m(\cos(\pi - \theta)) \times \\ \times \sin[(2l+1)\left(v + \frac{1}{2}\right)\pi] \frac{dv}{dp} dp.$$

Excluding from consideration the low frequency part of the field with $p < p$ and applying the asymptotic expansion of the associated Legendre polynomial for large values $n \sin \theta \gg |m|$, we arrive at

$$\sum_{\text{Ent}[v(p)]}^{\infty} f_n P_n^m(\cos \theta) \approx \frac{2\sqrt{2}}{\sqrt{\pi \sin \theta}} \sum_{l=0}^{\infty} (-1)^l \int_p^{\infty} v^{m+1/2} f(p) \frac{dv}{dp} \times \\ \times \cos\left[\left(v + \frac{1}{2}\right)(\pi - \theta) - \frac{\pi}{4} + \frac{m\pi}{2}\right] \sin\left[(2l+1)\left(v + \frac{1}{2}\right)\pi\right] dp.$$

$$\sum_{\text{Ent}[v(p)]}^{\infty} f_n \frac{dP_n^m(\cos \theta)}{d\theta} \approx \frac{2\sqrt{2}}{\sqrt{\pi \sin \theta}} \sum_{l=0}^{\infty} (-1)^l \int_p^{\infty} v^{m+1/2} f(p) \frac{dv}{dp} \times \\ \times \cos\left[\left(v + \frac{1}{2}\right)(\pi - \theta) - \frac{3\pi}{4} + \frac{m\pi}{2}\right] \sin\left[(2l+1)\left(v + \frac{1}{2}\right)\pi\right] dp.$$

Using these formulas for transforming (2.19a), we can obtain the following asymptotic expression for the displacements in surface waves, going around the sphere l times and passing through the pole $\theta = 0$ and the antipodes $\theta = \pi$, g times:

$$u_R(t, R, \theta, \varphi) = \frac{1}{\sqrt{2\pi R_0 \sin \theta}} \operatorname{Re} \int_p^{\infty} \frac{1}{p} \exp\left[i\left(pt + \frac{\pi}{4}(2g-1)\right)\right] \times \\ \times \left\{ \sum_{k=1}^{K_S(p)} \frac{U_{kS}(p, \varphi)}{I_{kvS} C_{kS}} \sqrt{\frac{v_{kS}}{R_0}} \tilde{V}_{kv}^{(1)} \exp\left[-i\left(v_{kS} + \frac{1}{2}\right)\tilde{\theta}\right] \right\} dp, \quad (2.20)$$

$$u_{\theta}(t, R, \theta, \varphi) = \frac{1}{\sqrt{2\pi R_0 \sin \theta}} \operatorname{Re} \int_p^{\infty} \frac{1}{p} \exp\left[i\left(pt - \frac{\pi}{4}(2g+3)\right)\right] \times \\ \times \left\{ \sum_{k=1}^{K_S(p)} \frac{U_{kS}(p, \varphi)}{I_{kvS} C_{kS}} \sqrt{\frac{v_{kS}}{R_0}} \tilde{V}_{kv}^{(2)} \exp\left[-i\left(v_{kS} + \frac{1}{2}\right)\tilde{\theta}\right] \right\} dp,$$

$$u_{\varphi}(t, R, \theta, \varphi) = \frac{1}{\sqrt{2\pi R_0 \sin \theta}} \operatorname{Re} \int_p^{\infty} \frac{1}{p} \exp\left[i\left(pt - \frac{\pi}{4}(2g-1)\right)\right] \times \\ \times \left\{ \sum_{k=1}^{K_T(p)} \frac{U_{kT}(p, \varphi)}{I_{kvT} C_{kT}} \tilde{V}_{kv}^{(3)} \exp\left[-i\left(v_{kT} + \frac{1}{2}\right)\tilde{\theta}\right] \right\} dp.$$

Here for $Q = S, T$, $K_Q(p)$ is the maximum number of harmonics in Rayleigh waves ($Q = S$) and Love waves ($Q = T$) for a given p ;

$$U_{kQ} = \sum_{m=-\infty}^{\infty} v_{kQ}^{m-1} D_{kmn}^Q \exp [im(\phi + (-1)^s \pi/2)], \quad (2.21)$$

where v_{kQ} is the analog of the wave number, a root of the equation $P_{kQ}^2(v) - p^2 = 0$.

The phase and group velocities of the k -th harmonic v_{kQ} and C_{kQ} along the surface of the sphere are expressed through v_{kQ} :

$$v_{kQ} = \frac{pR_0}{v_{kQ} + i\gamma_a}, \quad C_{kQ} = R_0 \frac{dp_{kQ}}{dv} (v = v_{kQ}). \quad (2.22)$$

The multiple passages of the waves through the poles $\tilde{\theta}(0 < \tilde{\theta} < \infty)$ are related to the coordinate of the observation point $\theta(0 \leq \theta < \pi)$

$$\tilde{\theta} = (-1)^s \theta + 2\pi(g - l).$$

The components u_F and u_θ form the Rayleigh wave; u_ϕ , the Love wave.

Expressions for P_{kQ} , C_{kQ} in terms of integrals of the eigenfunctions.

The method of variational analysis can be used (25) for $Q = S, T$ to obtain

$$P_{k\nu Q}^2 = \frac{N^2 G_{1k}^Q + 2NG_{2k}^Q + G_{3k}^Q}{R_0^2 I_{k\nu Q}}, \quad C_{kQ} = \frac{NG_{1k}^Q + G_{2k}^Q}{R_0 p_{kQ} I_{k\nu Q}}. \quad (2.23)$$

Here the integrals G_{1k}^Q – G_{3k}^Q have the form

$$\begin{aligned} G_{1k}^S &= \int_0^{R_0} [(\lambda + 2\mu)(\tilde{V}_{kv}^{(2)})^2 + \mu(\tilde{V}_{kv}^{(1)})^2] dR, \\ G_{2k}^S &= \int_0^{R_0} \left[\left(\mu \tilde{V}_{kv}^{(1)} \frac{d\tilde{V}_{kv}^{(2)}}{dR} - \lambda \tilde{V}_{kv}^{(2)} \frac{d\tilde{V}_{kv}^{(1)}}{dR} \right) R - (2\lambda + 3\mu) \tilde{V}_{kv}^{(1)} \tilde{V}_{kv}^{(2)} \right] dR, \\ G_{3k}^S &= \int_0^{R_0} \left\{ \left[(\lambda + 2\mu) \left(\frac{d\tilde{V}_{kv}^{(1)}}{dR} \right)^2 + \mu \left(\frac{d\tilde{V}_{kv}^{(2)}}{dR} \right)^2 \right] R^2 + \right. \\ &\quad \left. + 2 \left(2\lambda \tilde{V}_{kv}^{(1)} \frac{d\tilde{V}_{kv}^{(2)}}{dR} - \mu \tilde{V}_{kv}^{(2)} \frac{d\tilde{V}_{kv}^{(1)}}{dR} \right) R + 4(\lambda + \mu)(\tilde{V}_{kv}^{(1)})^2 - \mu(\tilde{V}_{kv}^{(2)})^2 \right\} dR, \\ G_{1k}^T &= \int_0^{R_0} \mu(\tilde{V}_{kv}^{(3)})^2 dR, \\ G_{2k}^T &\equiv 0, \end{aligned} \quad (2.24)$$

$$G_{3k}^T = \int_0^{R_0} \mu \left[\left(\frac{d\tilde{V}_{kv}^{(3)}}{dR} \right)^2 R^2 - 2\tilde{V}_{kv}^{(3)} \frac{d\tilde{V}_{kv}^{(3)}}{dR} R - (\tilde{V}_{kv}^{(3)})^2 \right] dR,$$

$$N = \sqrt{v(v+1)}.$$

Methods for calculation of $\tilde{V}_{kv}^{(i)} - P_{kvQ}^i$ are given in (3, 4, 14).

SURFACE WAVES FROM ELEMENTARY SOURCES. The source type influences only the terms U_{k0} in (2.20). We will consider expressions for U_{k0} for several elementary impacts, noting that $v \gg |m|$.

I. Radial axially symmetric force. Let

$$\mathbf{F} = F_R(t, R, \theta) \mathbf{a}_R.$$

From (2.7) we obtain:

$$\begin{aligned} f_{0v}^{(1)} &= (v_{ks} + 1/2) \int_{-\infty}^{+\infty} e^{-ipt} \int_0^\pi F_R P_v(\cos \theta) \sin \theta d\theta dt, \\ f_{mv}^{(1)} &= 0 \quad \text{при } m \neq 0, \quad f_{mv}^{(2)} \equiv f_{mv}^{(3)} \equiv 0, \\ U_{ks} &= \frac{1}{v_{ks}} \int_0^{R_0} f_{0v}^{(1)} \tilde{V}_{kv}^{(1)}(R) R^2 dR, \\ U_{kT} &= 0. \end{aligned} \tag{2.25}$$

In particular, for an ideally concentrated radial force at the point $R = H, \theta = 0$

$$\mathbf{F} = \delta(R - H) \frac{\delta(\theta)}{H^2 \sin \theta} \varphi(t) \mathbf{a}_R$$

$$U_{ks} = \tilde{V}_{kv}^{(1)}(H) \times S(p), \tag{2.26}$$

where $S(p) = \int_{-\infty}^{+\infty} e^{-ipt} \varphi(t) dt$ - time spectra of the source.

II. Axially symmetric force directed along a meridian. Let

$$\mathbf{F} = F_\theta(t, R, \theta) \mathbf{a}_\theta.$$

In this case we have from (2.7):

$$\begin{aligned} f_{mv}^{(1)} &\equiv f_{mv}^{(3)} \equiv 0; \quad f_{mv}^{(2)} = 0 \quad \text{при } m \neq 0; \\ f_{0v}^{(2)} &= \int_{-\infty}^{+\infty} e^{-ipt} \int_0^\pi F_\theta \frac{dP_v(\cos \theta)}{d\theta} \sin \theta d\theta dt; \\ U_{ks} &= \frac{1}{v_{ks}} \int_0^\infty f_{0v}^{(2)} \tilde{V}_{kv}^{(2)} R^2 dR; \\ U_{kT} &= 0. \end{aligned} \tag{2.27}$$

For an ideally concentrated meridional source acting at $R = H$, $\theta = 0$

$$\mathbf{F} = \frac{2\delta(R-H)\delta(\theta)\varphi(t)}{H^3 \sin^2 \theta} \mathbf{a}_\theta, \quad U_{kS} = -\frac{v_{kS}}{H} V_{kv}^{(2)}(H) S(p). \quad (2.28)$$

III. Rotational impact. Let

$$\mathbf{F} = F_\varphi(t, R, \theta) \mathbf{a}_\varphi.$$

From (2.7) we obtain

$$\begin{aligned} f_{mv}^{(1)} &\equiv f_{mv}^{(2)} \equiv 0; \quad f_{mv}^{(3)} = 0 \quad \text{при } m \neq 0; \\ f_{0v}^{(3)} &= - \int_{-\infty}^{+\infty} e^{-ipt} \int_0^\pi F_\varphi \frac{dP_v(\cos \theta)}{d\theta} \sin \theta d\theta dt; \\ U_{kS} &= 0; \quad U_{kT} = -\frac{1}{v_{kT}} \int_0^{R_0} f_{0v}^{(3)} V_{kv}^{(3)}(H) R^2 dR. \end{aligned} \quad (2.29)$$

In the special case of an ideally concentrated rotational impact acting at the point $R = H$, $\theta = 0$,

$$\begin{aligned} \mathbf{F} &= 2\delta(R-H) \frac{\delta(\theta)\varphi(t)}{H^3 \sin^2 \theta} \mathbf{a}_\varphi, \\ U_{kT} &= \frac{v_{kT}}{H} V_{kv}^{(3)}(H) S(p). \end{aligned} \quad (2.30)$$

IV. Tangential force field with fixed azimuth. We will consider only a particular case of this force field:

$$\mathbf{F} = F_T(t, R, \theta) \mathbf{a}_T,$$

where $\hat{\mathbf{a}}_T$ is a unit horizontal vector of fixed azimuth:

$$\begin{aligned} (\mathbf{a}_T, \mathbf{a}_R) &= 0, \quad (\mathbf{a}_T, \mathbf{a}_\theta) = \cos(\delta - \varphi), \\ (\mathbf{a}_T, \mathbf{a}_\varphi) &= \sin(\delta - \varphi). \end{aligned}$$

Then,

$$\mathbf{F} = F_T [\mathbf{a}_\theta \cos(\delta - \varphi) + \mathbf{a}_\varphi \sin(\delta - \varphi)]$$

and we obtain from (2.7):

$$\begin{aligned}
 f_{mv}^{(1)} &\equiv 0; \quad f_{mv}^{(2)} = f_{mv}^{(3)} = 0 \quad \text{при } m \neq \pm 1; \\
 f_{1v}^{(2)} &= \frac{e^{-i\delta} f_{Tv}}{2v(v+1)}; \quad f_{-1v}^{(2)} = \frac{-e^{-i\delta} f_{Tv}}{2}; \\
 f_{1v}^{(3)} &= \frac{e^{-i\delta} f_{Tv}}{2iv(v+1)}; \quad f_{-1v}^{(3)} = \frac{e^{i\delta} f_{Tv}}{2i}; \\
 f_{Tv} &= \int_{-\infty}^{+\infty} e^{-ipt} \int_0^\pi F_T \left(\frac{dP_v^1}{d\theta} + \frac{P_v^1}{\sin \theta} \right) \sin \theta d\theta dt.
 \end{aligned}$$

summing with respect to m ,

$$\begin{aligned}
 U_{kS} &= \frac{(-1)^g}{v_{kS}^2} i \cos(\delta - \varphi) \int_0^{R_0} f_{Tv} \tilde{V}_{kv}^{(2)} R^2 dR, \\
 U_{kT} &= -\frac{(-1)^g}{v_{kT}^2} i \sin(\delta - \varphi) \int_0^{R_0} f_{Tv} \tilde{V}_{kv}^{(3)} R^2 dR.
 \end{aligned} \tag{2.34}$$

In the case of a concentrated force at $R = H$, $\theta = 0$,

$$F = \frac{\delta(R-H)\delta(\theta)\varphi(t)a_T}{H^2 \sin \theta}$$

$$\begin{aligned}
 U_{kS} &= (-1)^g i \cos(\delta - \varphi) \tilde{V}_{kv}^{(2)}(H) S(p), \\
 U_{kT} &= -(-1)^g i \sin(\delta - \varphi) \tilde{V}_{kv}^{(3)}(H) S(p).
 \end{aligned} \tag{2.32}$$

V - VIII. Other elementary sources. The formulas for the remaining elementary sources considered in Section 1 can be obtained from equations (1.33)-(1.36) upon replacing U_{kR} by U_{kS} , U_{kL} by U_{kT} , ξ_{kR} by v_{kS}/H , ξ_{kT} by v_{kT}/H , $\tilde{V}_k^{(1)}$ by $\tilde{V}_{kv}^{(1)}$, $\tilde{V}_k^{(2)}$ by $(-1)^g \cdot \tilde{V}_{kv}^{(2)}$, and $\tilde{V}_k^{(3)}$ by $(-1)^g \tilde{V}_{kv}^{(3)}$.

Asymptotic formulae for (2.19) - (2.32) as $pR_0/b(R_0) \rightarrow \infty$. Let $pR_0/b(R_0) \rightarrow \infty$, and let the quantity $p(R_0 - R)/b(R_0) = pz/b(R_0)$ under these conditions. Then it is not difficult to show that equations (2.19) - (2.32) tend toward the corresponding formulas (1.19) - (1.32), if it is noted that as $pR_0/b(R_0) \rightarrow \infty$ and

$$\begin{aligned}
 g &= 0, \quad l = 0, \quad H = R_0 - h; \quad R_0 \theta \rightarrow r; \\
 U_{kR} &\rightarrow -U_{kz}; \quad U_{k0} \rightarrow U_{kr}; \quad \frac{v_{kQ}}{R} \approx \frac{v_{kQ}}{H} \rightarrow \xi_{kQ};
 \end{aligned}$$

$$\tilde{V}_{kv}^{(i)}(R_1) \tilde{V}_{kv}^{(i)}(R_2) \rightarrow \tilde{V}_k^{(i)}(z_1) \tilde{V}_k^{(i)}(z_2), \tag{2.33}$$

where $z_1 = R_0 - R_1$, $z_2 = R_0 - R_2$, $i = 1, 2, 3$;

$$\tilde{V}_{kv}^{(1)}(R_1) \tilde{V}_{kv}^{(2)}(R_2) \rightarrow -\tilde{V}_k^{(1)}(z_1) \tilde{V}_k^{(2)}(z_2).$$

3. Some properties of the derived solutions

We will consider some properties of the derived equations (1.20) and (2.20), which are essential for understanding the excitation and propagation of surface waves. Let us consider that the force field F , describing the seismic source, is localized in some zone, situated in the halfspace near the initial coordinates $z = 0$, $r = 0$, and for the sphere near the pole $R = R_0$, $\theta = 0$. We will locate a non-disturbing receiver at the point z, r, ϕ (R, θ, ϕ), which registers the q 'th component of displacement at this point ($q = z, r, \phi$ in a halfspace and $q = R, \theta, \phi$ for a sphere). The quantity r (or θ) has the meaning of epicentral distance, ϕ - azimuth at the epicenter to the station, z (or R) - depth of receiver, measured from the free surface (or center of sphere).

The theoretical seismogram $u_q(t)$ of the surface wave, recorded by such a receiver, is described by equations (1.20) or (2.20) for sufficiently large r (or θ). The Rayleigh surface wave (index $Q = R$ or S) will be recorded only for $q = z, r$, or R, θ : it is polarized in a vertical plane (section of great circle), passing through the epicenter and receiver. The Love surface wave (wave index $Q = L$ or T) will be recorded only for $q = \phi$: it is linearly polarized, the displacement vector normal to the polarization plane of Rayleigh waves.

Each seismogram can be considered as a superposition of an infinite number of harmonics (in other terminology - normal waves, overtones, or modes): the index k ($k = 1, 2, \dots, \infty$) designates the number of the harmonic. Let $u_{kq}(t)$ be the contribution of the k 'th harmonic to the q 'th component of the seismogram. Then

$$u_k(t) = \sum_{k=1}^{\infty} u_{kq}(t). \quad (3.1)$$

We will consider further not the seismogram itself, $u_q(t)$ or $u_{kq}(t)$ but their spectra, i.e., the Fourier transform with respect to time of the form $\int_0^{\infty} \Phi(p)e^{ipt} dt$; we designate these spectra as $\Phi_q(p)$ and $\Phi_{kq}(p)$ respectively. It is evident that

$$\Phi_q = \sum_{k=1}^{\infty} \Phi_{kq}, \quad (3.2)$$

$$u_{kq}(t) = \frac{1}{\pi} \operatorname{Re} \int_0^{\infty} \Phi_{kq}(p) e^{ipt} dp. \quad (3.3)$$

Surface wave spectra. The spectral density Φ_{kq} differs from zero only for $p > p_{kQ}^*$, where p_{kQ}^* is the limiting frequency in the spectra of the k 'th harmonic of the kQ 'th wave. The boundary frequencies are related as

$$\bar{p}_{k+1,Q} > \bar{p}_{kQ} > \bar{p}_{k-1,Q} > \dots > \bar{p}_{1Q}.$$

In a halfspace $p_1 L = p_1 R = 0$, in a sphere $p_1 S > 0$, $p_1 T > 0$. However since we are not interested in the low frequency part of the spectra of the disturbance with $p < p_*$, which is impossible to express in the form of a propagating wave, it will be accepted that for surface waves $\Phi_{kq} = 0$ for $p < \max(p_*, p_{kQ}^*)$.

The spectral density can be expressed in the form of a product

$$\Phi_{kq} = \prod_{i=0}^4 B_{kq}^{(i)}. \quad (3.4)$$

For the terms $B_{kq}^{(i)}$ we have the following formulae:

a) in a halfspace (where $Q = R$ for $q = z, r$; $Q = L$ for $q = \phi$):

$$B_{kq}^{(0)} = B^{(0)} = \sqrt{\frac{\pi}{2}} e^{-i\frac{\pi}{4}}, \quad (3.5)$$

$$B_{kq}^{(1)} = [v_{kQ}(p) C_{kQ}(p) I_{kQ}(p)]^{-1}, \quad (3.6)$$

$$B_{kq}^{(2)} = \frac{\exp[-ir\xi_{kQ}(p)]}{\sqrt{r\xi_{kQ}(p)}}, \quad (3.7)$$

$$B_{kq}^{(3)} = U_{kQ}(p, \varphi), \quad (3.8)$$

$$B_{kq}^{(4)} = \alpha_q \tilde{V}_k^{(j_q)}(p), \quad j_z = 1, \quad j_r = 2, \quad j_\phi = 3, \quad (3.9)$$

$$\alpha_z = 1, \quad \alpha_r = -i, \quad \alpha_\phi = i;$$

b) in a sphere (where $Q = S$ for $q = R, \theta$, $Q = T$ for $q = \phi$):

$$B_{kq}^{(0)} = B^{(0)} = \sqrt{\frac{\pi}{2}} e^{-i\frac{\pi}{4}}, \quad (3.5a)$$

$$B_{kq}^{(1)} = [v_{kQ}(p) C_{kQ}(p) I_{kQ}(p)]^{-1}, \quad (3.6a)$$

$$B_{kq}^{(2)} = \frac{\exp\left\{-i\left[\left(v_{kQ}(p) + \frac{1}{2}\right)\tilde{\theta} - \frac{\pi}{2}g\right]\right\}}{\sqrt{v_{kQ}(p)\sin\theta}}, \quad (3.7a)$$

$$B_{kq}^{(3)} = U_{kQ}(p, \varphi), \quad (3.8a)$$

$$B_{kq}^{(4)} = \alpha_q \tilde{V}_k^{(j_q)}(R), \quad j_R = 1, \quad j_\theta = 2, \quad j_\phi = 3,$$

$$\alpha_R = 1, \quad \alpha_\theta = i(-1)^{\theta+1}, \quad \alpha_\phi = i(-1)^\phi. \quad (3.9a)$$

g is the number of passages through the epicenter and anti-epicenter.

The multiplier $B_{kq}^{(1)}$ depends only on the properties of the medium (variation of velocity and density with depth or radius) and the type of waves recorded.

(2) The multiplier $B_{kq}^{(2)}$ describes the propagation effect of the waves: the numerator defines the phase delay of the oscillations for a given frequency, the denominator--the decrease of amplitude due to geometric spreading along the path r (or θ). The additional phase contribution of $\pi g/2$ in a sphere arises on account of g -passages through the epicenter and anti-epicenter.

(3) The multiplier $B_{kq}^{(3)}$ for a given wave depends only on the source mechanism, and also on the azimuth of the recording station with respect to the source. For an axially symmetric source $B_{kq}^{(3)}$ does not depend on ϕ .

(4) The multiplier $B_{kq}^{(4)}$ depends on the depth of burial of the receiver and its orientation (i.e., for which component of displacement it is set up to measure).

Polarization of Rayleigh waves. The ratio of spectral densities Φ_{kr}/Φ_{kz} or $\Phi_{k\theta}/\Phi_{kR}$ defines the polarization of Rayleigh waves. Since $B_{kr}^{(j)} = qB_{kz}^{(j)}$ and $B_{kR}^{(j)} = B_{k\theta}^{(j)}$ for $j < 4$, then

$$\frac{\Phi_{kr}}{\Phi_{kz}} = \frac{B_{kr}^{(4)}}{B_{kz}^{(4)}} = -i \left[\frac{\tilde{V}_k^{(2)}(p, z)}{\tilde{V}_k^{(1)}(p, z)} \right], \quad (3.10)$$

$$\frac{\Phi_{k\theta}}{\Phi_{kR}} = \frac{B_{k\theta}^{(4)}}{B_{kR}^{(4)}} = i(-1)^{g+1} \left[\frac{\tilde{V}_{kv}^{(2)}(R)}{\tilde{V}_{kv}^{(1)}(R)} \right]. \quad (3.10a)$$

Thus the eigenfunctions $\tilde{V}_k^{(i)}$, $\tilde{V}_k^{(i)}$ define the Rayleigh wave polarization ellipse; the direction of rotation and ratio of the ellipse axes depends on the depth of burial of the receiver; The apparent change in the direction of rotation for a sphere for changes in g is associated with the fixed choice of direction in designating u_θ (+ for motion away from the epicenter).

Reciprocity principle. For simple sources of the concentrated force type, it is easy to establish the principle of reciprocity--the invariance of spectral density Φ_{kz} for change in the placement of the source and receiver with preservation of the source orientation with respect to the receiver and vice versa. In fact for a simple vertical force at depth we have according to (1.26):

$$\begin{aligned}
B_{kq}^{(3)} &= \tilde{V}_k^{(1)}(p, h) S(p), \\
\Phi_{kz}^{\downarrow}(h, z) &= \left[\prod_{i=0}^3 B_{kz}^{(i)} \right] S(p) \tilde{V}_k^{(1)}(p, h) \tilde{V}_k^{(1)}(p, z), \\
\Phi_{kr}^{\downarrow}(h, z) &= \left[\prod_{i=0}^3 B_{kr}^{(i)} \right] S(p) \tilde{V}_k^{(1)}(p, h) (-i) \tilde{V}_k^{(2)}(p, z).
\end{aligned} \quad (3.11)$$

For a simple horizontal force at depth h , oriented toward the receiver according to (2.32) we have

$$\begin{aligned}
B_{kq}^{(3)} &= i \tilde{V}_k^{(2)}(p, h) S(p), \\
\Phi_{kz}^{\rightarrow}(h, z) &= \left[\prod_{i=0}^3 B_{kz}^{(i)} \right] S(p) i \tilde{V}_k^{(2)}(p, h) \tilde{V}_k^{(1)}(p, z), \\
\Phi_{kr}^{\rightarrow}(h, z) &= \left[\prod_{i=0}^3 B_{kr}^{(i)} \right] S(p) \tilde{V}_k^{(2)}(p, h) \tilde{V}_k^{(2)}(p, z).
\end{aligned} \quad (3.12)$$

For a simple horizontal force at depth h perpendicular to the epicenter-receiver direction, from (1.32)

$$\begin{aligned}
B_{kq}^{(3)} &= -i \tilde{V}_k^{(3)}(p, h) S(p), \\
\Phi_{k\varphi}^{\rightarrow}(h, z) &= \left[\prod_{i=0}^3 B_{k\varphi}^{(i)} \right] S(p) \tilde{V}_k^{(3)}(p, h) \tilde{V}_k^{(3)}(p, z).
\end{aligned} \quad (3.13)$$

From this the reciprocity relations follow:

$$\begin{aligned}
\Phi_{kz}^{\downarrow}(h, z) &= \Phi_{kz}^{\downarrow}(z, h), \\
\Phi_{kr}^{\rightarrow}(h, z) &= \Phi_{kr}^{\rightarrow}(z, h), \\
\Phi_{k\varphi}^{\rightarrow}(h, z) &= \Phi_{k\varphi}^{\rightarrow}(z, h), \\
\Phi_{kz}^{\rightarrow}(h, z) &= -\Phi_{kz}^{\downarrow}(z, h), \\
\Phi_{kr}^{\downarrow}(h, z) &= -\Phi_{kr}^{\downarrow}(z, h).
\end{aligned} \quad (3.14)$$

Similar relations for a sphere are not difficult to obtain.

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Additional force systems. These forces have been normalized to represent unit forces of couples acting at the source. They are suitable for direct substitution in to (1.20). The result (1.20) agrees with the theoretical equations of Saito and Tsai and Aki (1970).

(a) Arbitrarily directed force. (δ, β). The force is represented as

$$\vec{F} = \frac{\delta(z-h)\delta(r)}{2\pi r} \left[\cos(\delta-\phi)\sin\beta \vec{a}_r + \sin(\delta-\phi)\sin\beta \vec{a}_\phi + \cos\beta \vec{a}_z \right] s(t)$$

where δ is the azimuth of the force from north and β is the angle that the force makes with the positive z axis, which is taken to be downward. Note that $0 \leq \delta \leq 2\pi$ and $0 \leq \beta \leq \pi$. The \vec{a}_q are the unit vectors in the q'th direction. $s(t)$ is the source time function. The U_{kQ} become

$$U_{kR} = \left[\cos\beta V_k^{(1)}(h,p) - i \sin\beta \cos(\delta-\phi) V_k^{(2)}(h,p) \right] . S(p)/2\pi$$

$$U_{kL} = -i \sin\beta \sin(\delta-\phi) \int_{-\infty}^{\infty} \exp(-ipt) s(t) dt; V_k^{(3)}(h,p) . S(p)/2\pi$$

$$\text{where } S(p) = \int_{-\infty}^{\infty} \exp(-ipt) s(t) dt;$$

(b) Dipole without moment. (δ, β).

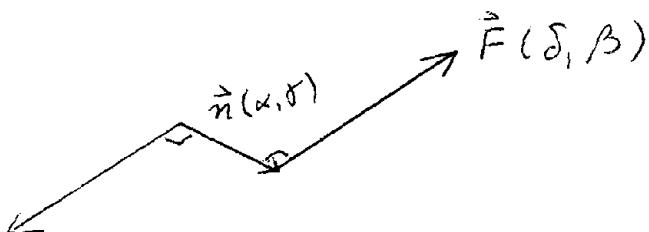
$$U_{kR} = \left[\cos^2\beta \frac{dV_k^{(1)}}{dh}(h,p) - k_R^2 \sin^2\beta \cos^2(\delta-\phi) V_k^{(2)}(h,p) \right. \\ \left. + \frac{i}{2} \sin 2\beta \cos(\delta-\phi) (k_R V_k^{(1)}(h,p) + \frac{dV_k^{(2)}}{dh}(h,p)) \right] . S(p)/2\pi$$

$$U_{kL} = i \sin\beta \sin(\delta-\phi) \left[\cos\beta \frac{dV_k^{(3)}}{dh}(h,p) + i k_L \cos(\delta-\phi) \sin\beta V_k^{(3)}(h,p) \right] . S(p)/2\pi$$

(c) Couple (δ, β), (α, γ). The positive force in the direction (δ, β) is converted into a couple by the application of an operator to (a):

$$\left[\cos\gamma \frac{d}{dh} - ik_Q \cos(\alpha-\phi) \sin\alpha \right], Q = L \text{ or } R;$$

graphically, the couple appears as



¹ This differs from section V normalization to a unit force.

by factor $(2\pi)^{-1}$ which is needed for
A-31

The u_{kQ} become

$$U_{kR} = \{ \cos\gamma \cos\beta \frac{dV_k^{(1)}}{dh}(h, p) - k_R \sin\beta \sin\gamma \cos(\delta-\phi) \cos(\alpha-\phi) V_k^{(2)}(h, p) \\ + ik_R \sin\gamma \cos\beta \cos(\alpha-\phi) V_k^{(1)}(h, p) \\ + ik_R \cos\gamma \sin\beta \cos(\delta-\phi) \frac{dV_k^{(2)}}{dh}(h, p) \} \cdot \frac{S(p)}{2\pi}$$

$$U_{kL} = -i \sin\beta \sin(\delta-\phi) \left[\cos\gamma \frac{dV_k^{(3)}}{dh}(h, p) \right. \\ \left. + ik_L \sin\gamma \cos(\alpha-\phi) V_k^{(3)}(h, p) \right] \cdot S(p)/2\pi$$

where (δ, β) and (α, γ) are related by the condition of perpendicularity

$$\operatorname{ctg}\gamma \operatorname{ctg}\beta = -\cos(\delta-\alpha).$$

(d) Two perpendicular dipoles without moment of equal magnitude and opposite sign (δ, β) , (α, γ) . For this set of dipoles, the angles define the directions of the tensional T and compressional P axes by

$$(\delta, \beta) = T \quad (\alpha, \gamma) = P.$$

$$U_{kR} = \{ (\cos^2\beta - \cos^2\gamma) \frac{dV_k^{(1)}}{dh}(h, p) \\ - k_R \left[\sin^2\beta \cos^2(\delta-\phi) - \sin^2\gamma \cos^2(\alpha-\phi) \right] V_k^{(2)}(h, p) \\ + \frac{i}{2} \left[\sin 2\beta \cos(\delta-\phi) - \sin 2\gamma \cos(\alpha-\phi) \right] \cdot \\ \cdot \left[k_R V_k^{(1)}(h, p) - \frac{dV_k^{(2)}}{dh}(h, p) \right] \} \frac{S(p)}{2\pi}$$

$$U_{kL} = \{ -\frac{i}{2} \left[\sin 2\beta \sin(\delta-\phi) - \sin 2\gamma \sin(\alpha-\phi) \right] \frac{dV_k^{(3)}}{dh}(h, p) \\ + \frac{1}{2} k_L V_k^{(3)}(h, p) \left[\sin^2\beta \sin(2\delta-2\phi) - \sin^2\gamma \sin(2\alpha-2\phi) \right] \} \frac{S(p)}{2\pi}$$

with the necessary perpendicularity relation

$$\operatorname{ctg}\gamma \operatorname{ctg}\beta = -\cos(\delta-\alpha).$$

(e) Double couple without moment (δ, β) , (α, γ) . The angles are defined as in (c).

$$U_{kR} = \{ 2 \cos\gamma \cos\beta \frac{dV_k^{(1)}}{dh}(h, p) - 2k_R V_k^{(2)} \sin\beta \sin\gamma \cos(\delta-\phi) \cos(\alpha-\phi) \\ + i \left[k_R V_k^{(1)}(h, p) + \frac{dV_k^{(2)}}{dh} \right](h, p) . \\ . \left[\sin\gamma \cos\beta \cos(\alpha-\phi) + \sin\beta \cos\gamma \cos(\delta-\phi) \right] \cdot S(p)/2\pi \\ U_{kL} = \{ -i \left[\sin\beta \sin(\delta-\phi) \cos\gamma + \sin\gamma \cos\beta \sin(\alpha-\phi) \right] \frac{dV_k^{(3)}}{dh}(h, p) \\ - k_L V_K^{(1)}(h, p) \sin\beta \sin\gamma \sin(\alpha+\delta - 2\phi) \} S(p) 2\pi$$

and

$$\operatorname{ctg}\gamma \operatorname{ctg}\beta = - \cos(\delta-\alpha)$$

(f) Explosive source.

$$U_{kR} = \left[\frac{dV_k^{(1)}}{dh}(h, p) - k_R V_k^{(2)}(h, p) \right] S(p)/2\pi$$

$$U_{kL} = 0.$$

(g) Double couple mechanism in terms of strike, dip, and slip of motion on the fault plane. The orthogonality conditions on the couples or dipoles of (d) or (e) imply that only three quantities are required to define motion on the fault plane. These quantities are the fault plane dip d , the strike θ , and the slip s . These parameters are shown in Figure A.1. The dip is measured from the horizontal in a downward direction and varies between 0° and 90° . The strike is measured clockwise from north and varies from 0° through 360° . The slip angle gives the direction of motion of the hanging wall with respect to the foot wall; the slip angle varies from 0° to 360° and is measured counterclockwise from the strike.

With this convention, the U_{kQ} become

$$U_{kR} = \frac{S(p)}{2\pi} \left\{ \sin s \sin 2d \left[\frac{dv_k^{(1)}}{dh} (h, p) + \frac{1}{2} \xi_{kR} v_k^{(2)} (h, p) \right] \right. \\ \left. + \xi_{kR} v_k^{(1)} (h, p) \left[-\cos s \sin d \sin 2(\phi-\theta) \right. \right. \\ \left. \left. - \frac{1}{2} \sin s \sin 2d \cos 2(\phi-\theta) \right] \right. \\ \left. + i \left[\xi_{kR} v_k^{(1)} (h, p) + \frac{dv_k^{(2)}}{dh} (h, p) \right] \cdot \right. \\ \left. \left[-\cos d \cos s \cos(\phi-\theta) \right. \right. \\ \left. \left. + \cos 2d \sin s \sin(\phi-\theta) \right] \right\}$$

$$U_{kL} = \frac{S(p)}{2\pi} \left\{ -i \left[\sin s \cos 2d \cos(\phi-\theta) \right. \right. \\ \left. \left. + \cos s \cos d \sin(\phi-\theta) \right] \frac{dv_k^{(3)}}{dh} (h, p) \right. \\ \left. + \xi_{kL} v_k^{(3)} (h, p) \left[\cos s \sin d \cos 2(\phi-\theta) \right. \right. \\ \left. \left. - \frac{1}{2} \sin s \sin 2d \sin 2(\phi-\theta) \right] \right\}$$

A reverse fault striking north and dipping 45° to the east is given by $\theta = 0^\circ$, $d = 45^\circ$, and $s = 90^\circ$. A right lateral vertical strike slip fault which strikes north is given by $\theta = 0^\circ$, $d = 90^\circ$, and $s = 180^\circ$.

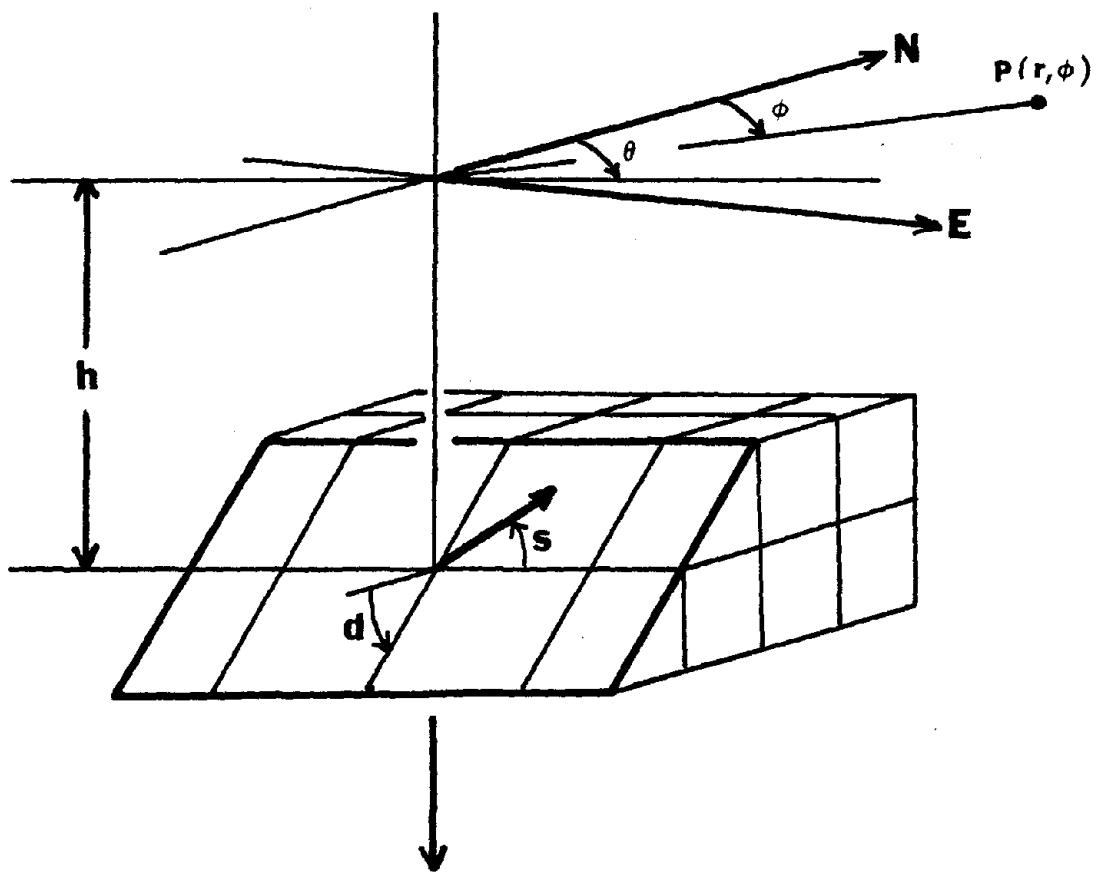


Fig. A. 1. Coordinate and fault plane geometry

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