

# NATIONAL CENTER FOR EARTHQUAKE ENGINEERING RESEARCH

State University of New York at Buffalo

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# Low-Level Dynamic Characteristics of Four Tall Flat-Plate Buildings in New York City

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by

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conducted at

Department of Civil Engineering and Operations Research School of Engineering and Applied Sciences Princeton University Princeton, New Jersey 08544

and

Lamont-Doherty Earth Observatory of Columbia University Palisades, New York 10964

Technical Report NCEER-92-0034

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### **Low-Level Dynamic Characteristics of Four Tall Flat-Plate Buildings in New York City**

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### PREFACE

The National Center for Earthquake Engineering Research (NCEER) was established to expand and disseminate knowledge about earthquakes, improve earthquake-resistant design, and implement seismic hazard mitigation procedures to minimize loss of lives and property. The emphasis is on structures in the eastern and central United States and lifelines throughout the country that are found in zones of low, moderate, and high seismicity.

NCEER's research and implementation plan in years six through ten (1991-1996) comprises four interlocked elements, as shown in the figure below. Element I, Basic Research, is carried out to support projects in the Applied Research area. Element II, Applied Research, is the major focus of work for years six through ten. Element III, Demonstration Projects, have been planned to support Applied Research projects, and will be either case studies or regional studies. Element IV, Implementation, will result from activity in the four Applied Research projects, and from Demonstration Projects.



Research in the Building Project focuses on the evaluation and retrofit of buildings in regions of moderate seismicity. Emphasis is on lightly reinforced concrete buildings, steel semi-rigid frames, and masonry walls or infills. The research involves small- and medium-scale shake table tests and full-scale component tests at several institutions. In a parallel effort, analytical models and computer programs are being developed to aid in the prediction of the response of these buildings to various types of ground motion.

I Two of the short-term products of the **Building Project** will be a monograph on the evaluation of lightly reinforced concrete buildings and a state-of-the-art report on unreinforced masonry.

The **structures and systems program** constitutes one of the important areas of research in the **Building Project.** Current tasks include the following:

- 1. Continued testing of lightly reinforced concrete external joints.
- 2. Continued development of analytical tools, such as system identification, idealization, and computer programs.
- 3. Perform parametric studies of building response.
- 4. Retrofit of lightly reinforced concrete frames, flat plates and unreinforced masonry.
- 5. Enhancement of the IDARC (inelastic damage analysis of reinforced concrete) computer program.
- 6. Research infilled frames, including the development of an experimental program, development of analytical models and response simulation.
- 7. Investigate the torsional response of symmetrical buildings.

*The evaluation of existing buildings for seismic effects has been one of the main goals of the Building Project. The primary emphasis has been on concrete structures that were designed only for gravity loads. One problem in the evaluation and design of structures for dynamic loads is the accurate analytical representation of the stiffness and damping properties. This report presents the results offield tests on four flat-plate structures, which are a very common type of structural system in the East. Such tests on real buildings include the effects of the interaction of walls with the other elements. The low-level ambient wind vibration tests help the calibration of analytical methods, especially in the determination ofthe effective width offloor slabs.*

#### **ABSTRACT**

Many reinforced concrete structures utilize flat plates to provide lateral resistance when architectural constraints prevent wide-spread use of shear-walls. Understanding the resistance of flat-plate frames to large lateral loads is important for serviceability as well as seismic vulnerability assessment of hundreds of buildings on the East Coast. In February of 1991, the authors collected ambient vibration measurements to study the behavior of four flat plate, reinforced concrete structures in Manhattan ranging from 27 to 52 stories. All but one structure rests on a rock supported foundation. This report presents the development of ambient vibration analysis software based on the fast Fourier transform. Long ambient vibration records allow accurate power spectrum estimation. A robust peak picking method, also tailored for ambient vibration data, facilitates the interpretation of autopower and phase spectra. Damping estimates based on spectral peak band-widths are very sensitive to bias and leakage errors in the spectrum estimate. However, root-mean-square statistics of acceleration as well as of velocity and displacement can reasonably be estimated from the auto-power spectrum of response acceleration. Measured fundamental periods are shorter than those calculated from the 1982 UBC formula, but are longer than those calculated from the 1988 UBC Code or the proposed NYC Seismic Code. The results presented herein provide a base-line set of periods, deflections, and damping ratios to be compared to results expected to be obtained during strong winds. Periods estimated from ambient, mostly windinduced vibration measurements are used to calibrate finite element models of the structures, and are compared to values calculated from code design rules. A future goal of this on-going research is to relate effective flat plate width to lateral load level at service load levles, i.e., between ambient conditions and strong winds. Estimating dynamic parameters under various loading conditions coupled with ultimate load tests on model structures is intended to reveal the load-dependent stiffness of these structures.

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#### **ACKNOWLEDGEMENTS**

The structural engineering firm of Rosenwasser/Grossman, P.C., initiated this study, designed the buildings, and arranged access to them. The researchers thank the owners of the four buildings tested, The DeMatteis Organization, Grand Palais Associates, and Rose Associates, for their cooperation in this study. Financial support for this project was provided by Rosenwasser/Grossman, P.C., and The National Center for Earthquake Engineering Research under grant No's. NCEER 91-1011 and 91-1031.

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# **SECTION 1 INTRODUCTION**

Since the late 1960's, building codes for tall structures in New York City have increasingly emphasized structural resistance to wind and other lateral loads. Forthcoming New York City Building Code revisions will include seismic provisions and will increase the lateral design loads for many tall structures. Many high-rise structures in New York City are residential, especially in the upper stories. Architectural constraints in high-rise apartment buildings preclude the prevalent usage of shear walls to resist lateral forces. Consequently, in order to enhance their lateral resistance, several hundred buildings have incorporated, and rely upon, flat plates. In the structures studied herein, the flat-plate system is designed to resist a considerable portion of the lateral loads, the remainder of which is carried by shear walls interacting with the frame. In existing design codes for seismically active regions of the country, the base shear resistance of the structure is prescribed as a function of the structure's first natural period. Codes provide simplified formulas for estimating a building's fundamental period at its ultimate capacity. However, the actual dynamic properties of the structure at service load levels and at limit state levels may differ considerably from values estimated using codified formulas. Periods computed in the design process, which correspond to a 100-year wind producing roof deflection of H/400, are longer than the measured ambient periods, which correspond to deflections as small as  $H/2,000,000$ . In addressing wind response, serviceability considerations, and occupant comfort, extensive dynamic measurements of some of the tallest structures in New York City, such as the World Trade Center (aspect ratio  $= 6$ ), and several other slender concrete structures (aspect ratio  $\geq$ 10), have been undertaken in the past [Grossman 1990]. However, many moderately sized buildings, between 30 and 60 stories, have, as of yet, been overlooked by such dynamic analyses. The current trend toward more slender structures and reduced foot-print areas could result in greater sway in, albeit, shorter structures, and present an engineering challenge not found in squatter, although taller buildings. The combined conditions of the reduction in shear wall areas, more slender structures, consideration of occupant comfort, and more demanding codes with respect to lateral loads call for a method to assess the strength of flat plates in resisting lateral loads at ultimate conditions, as well as the behavior of flat-plate structures at service load levels. This report presents work in progress to meet these challenges. The ambient vibration measurements of this study are being compared to partially completed detailed computer modeling by the designers of the four "as built" reinforced concrete flat-plate framed structures. The design methodology used in the comparison [Grossman, 1987] was fine tuned in a parallel study [Grossman, in progress] using ultimate laboratory load tests of flat-plate sub-structures [Moehle, 1990].

Of key interest is the contribution of several structural and non-structural elements to the structure's lateral resistance, and the degradation of the lateral resistance at greater deflections. The dependency of parameters such as the effective flat-plate width, the effective shear wall stiffness, and the effective coupling-beam moment of inertia, on the deflection of the structure is uncertain, and needs to be determined reliably.

For purposes of comparing lateral sway among different buildings, roof deflections can be expressed as a fraction of the total height of the structure, H. For example, typical roof deflections under ambient conditions are less than H/800, the 100-year wind deflections are on the order of H/400, and seismic deflections should not exceed H/200.

The scale of these structures virtually precludes forced response experiments and traditional modal analysis. Nonetheless, ambient vibrations are very informative and have certain advantages, namely:

1. The length of time for data recording can be arbitrarily large, allowing for long-duration spectral averages and, hence, improved spectral resolution.

- 2. Structures vibrating in ambient conditions (at small amplitudes) behave more linearly than structures responding to earthquakes. Structural parameters, estimated from response records alone, are time-invariant if the excitation is stationary and the structure exhibits linear elastic behavior over the course of a single measurement record.
- 3. The variance of the response acceleration is nearly independent of the section of the data record being analyzed. Adjusting the recorder gains to the known signal level before recording begins, maximizes the signal-to-noise ratio of the recording, without the risk of clipping the data.
- 4. Comparing the motion of the structure subjected to forces during strong winds to predetermined thresholds of human comfort or perception is useful in assessing serviceability.

The immediate objectives of this study are to:

- I. Develop parameter estimation software for the analysis of ambient vibration measurements from large scale structural systems. The software package computes auto-power spectra and phase spectra between a "measurement" or "response" sensor and a "reference" sensor. It also picks spectral peaks, a feature which facilitates the estimation of the following dynamic parameters: natural frequency, amplitude, phase, damping ratio, and root-mean-square displacement, velocity, and acceleration. The programs require discrete building-response (output) time-records only.
- 2. Verify the results from the programs developed in step 1 with a parallel but independently developed routine. Verify the root-mean-square computations with small-scale tests on laboratory models.
- 3. Compare experimentally estimated periods with periods computed (in a parallel study) by finite-element modeling of the "as-built" structure.
- 4. Establish a set of base-line parameters for comparison with future tests during different

wind conditions.

The <u>long-range objectives</u> of this study are to:

- 1. Assess the actual contribution of flat-plates, shear walls, and coupling beams to the overall lateral stiffness in four similar, reinforced concrete structures.
- 2. Compare the buildings' compliance with current and proposed New York City design codes for wind and seismic forces respectively.
- 3. Compare the natural periods of structures under construction to their periods after some wind storms. Model the measured buildings in their "as built" condition using state of the art software not available to the design team at the time the structures were designed. Investigate the influence of partitions and cladding on the building's response to lateral loads.
- 4. Investigate the influence of foundations and soil conditions on the measured building response at different periods.

**In** conducting the measurements and analysis, the researchers enlisted equipment and expertise from the Lamont-Doherty Geological Observatory of Columbia University (L-DGO), and from the Department of Civil Engineering and Operations Research of Princeton University. The buildings measured were designed by the firm of Rosenwasser/Grossman, P.C. L-DGO has a wide interest in the measurement and analysis of earthquake-induced ground motions. **It** provided the geophones, accelerometers, recording instruments, and some of the preprocessing and analysis hardware used in this study. Professor Shi Yuan's experience in ambient vibration analysis expedited the development of a systematic and robust software package. The spectral analysis routines were written at Princeton University's Department of Civil Engineering and Operations Research, and verified by methods independently developed at L-DGO.

Following this introduction, Section 2 contains a description of the four buildings chosen for the study. Section 3 reviews the standard code procedures of estimating a building's period. Sections 4 and 5 present the measurement and experimental data analysis methods respectively. Section 6 compares the finite-element results to the experimental results. Section 7 discusses the research underway to evaluate the resistance of flat plates to lateral loads, and Section 8 contains the conclusions.

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### SECTION 2 DESCRIPTION OF THE BUILDINGS

The four reinforced-concrete frame structures each represent a common type of east-coast construction; the design philosophy for all the buildings requires using flat plates to improve resistance to lateral forces. Buildings A and B were under construction during the measurements. Building A is a new building and was clad up to two-thirds its height at the time of the measurements; it was not occupied. It had interior partitions up to the  $45<sup>th</sup>$  floor. The partitions' mass was positioned (stored) in the upper stories, although the partitions were not yet installed. Building B was a new, bare, unclad, and unoccupied concrete frame at the time of the measurements. Buildings C and D were completed and occupied during the measurements; they both had brick cladding. Building D was constructed during the mid 1970's, was topped-out but remained partially un-clad for several years during the 1970's recession. It was finished several years later and is now fully occupied. Figures 2-1 to 2-4 describe the typical structural floor plans, the elevation views of the four buildings, and other pertinent information. Table 2-1 summarizes some of the building features.



Building A is the only building of the four utilizing a pile foundation. Twenty-five feet of fill

overlays 90 to 130 feet of poorly graded, saturated sands with little silt, having standard penetration test blow counts (N values) ranging between medium compact ( $10 < N < 30$ ) at the top 50 feet of the sand strata to compact ( $N > 30$ ) at the deeper strata ( $> 75$ ) where the piles were fetched. These strata occupy a portion of the southern end of Manhattan. They have a long history of settlement of existing buildings when, at adjacent lots, displacement piles are driven (by impact hammers) to support new structures. Augured piles were therefore recommended for Building A when in close proximity to adjacent existing structures, in order to reduce settlement that would have been caused by the vibration from impact hammers. However, the augured piles were difficult to control during the construction period because of loss of ground into the augured shaft.

The water table is approximately 10 feet below the fill (35' to 40' below street level) and liquefaction of the top strata, while a remote possibility, needed not to be considered for this structure in accordance with the existing or newly proposed NYC Code. Piles are long (90 feet minimum for the augured piles and about 110 feet for the driven pipe piles). Even if some liquefaction were to occur, it is unlikely that it would cause the loss of building support.  $S_3 =$ 1.5 would have been the appropriate site category and coefficient for this structure had it been designed according to the now newly proposed NYC Seismic Code. The building's design preceded the proposed NYC Seismic Code.

Buildings Band D are supported by spread footings, bearing on rock of class 3-65 and 2-65 [20-40 ton/sq.ft. capacity] and the proposed NYC Seismic Code assigns a site coefficient of So  $=0.67$  for such sites. Building C is supported on spread footings and short piles to rock [class] 3-65, 20 ton/sq.ft.] and its site coefficient will require further study (but is likely to be  $S1 =$ 1.0) .

All four buildings are considered by the proposed NYC Seismic Code to be Dual Systems

having concrete shear walls with concrete Ordinary Moment Resisting Frame [OMRF (Rw = 5)].

Concrete ordinary moment resisting frames (OMRF's) are, in these cases, flat plates without spandrel beams. This system became popular in eastern U.S. cities after the Second World War, initially with spandrel beams, but since the 1970's often without them. The flat plate is therefore an important lateral load-resisting structural member.

Shear-wall frame interaction is present in all four structures. In building A, for example, about 50% of the total lateral overturning moment is resisted by shear deformation of the columns and slabs and the remainder by flexural action of the coupled structural walls. Recent laboratory testing [Hwang and Moehle 1990, 1991a, 1991b] and a sensitivity review [Grossman 1991] allow determination of the flat-plate effectiveness at different lateral load levels. Other studies [Grossman 1987, Moehle and Wallace 1990] provide recommendations to estimate the structural wall stiffness acting under varied gravity and lateral loads. A followup study will attempt to verify these earlier studies by comparing the measured dynamic properties of the structures with the computed dynamic properties obtained using the methodologies proposed in Section 7 to estimate stiffnesses.

The flat plates were designed to support a considerable portion of the lateral loads in addition to gravity loads. The column-slab joints were designed to transfer large moments and continuity of bottom mat reinforcing was required to account for reversal of moments. In these structures the integrity requirements of the 1989 ACI Code were generally observed, even though they were constructed before such requirements became mandatory.

2-3







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FIGURE 2-3 BUILDING C - Typical Floor Plan and Elevation

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#### FIGURE 2-4 BUILDING D - Typical Floor Plan and Elevation

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# **SECTION** 3 **CODE-PRESCRIBED ESTIMATION OF FUNDAMENTAL PERIODS**

In order to estimate the lateral forces acting on a structure it is necessary first to estimate its fundamental period. Several factors contribute to the uncertainty of the estimation of fundamental periods from design plans, such as: variations in the modulus of elasticity, section dimensions, stripping operations, creep, and shrinkage. Over the life of the structure, cracks in structural members and disengagement of the partitions reduce the stiffness of the structure and lead to greater uncertainty of the periods. It is therefore difficult during the design stage to predict the dynamic properties of a structure at an arbitrary point in its service life. Without accurate estimates of fundamental periods, results from wind-tunnel experiments lead to inaccuracies in wind-related design requirements. Furthermore, the design base-shear for earthquake resistant design is a function of the structure's period. The applicable national or local Codes must be consulted to determine the empirical lower bound fundamental periods in order to establish the minimum lateral load requirements. It is apparent at this time that such Codes have not settled on a uniform method to determine the periods. The West Coast experience has been influencing this empirical approach which more accurately describes the stiffer structures designed to withstand large seismic forces in California. An exception is the earlier 1982 UBC Code.

The Uniform Building Code of 1982 specified that the period of a multi-story framed building can be estimated by dividing the total number of stories by ten [Blume, 1979]. This method results in longer periods than those measured in this study.

$$
1982 \text{ UBC} \t\t T = N/10 \t\t (3.1)
$$

 $N =$  number of stories of the building

The 1988 version of the Uniform Building Code allows for one of two methods to be used in computing the natural period. The first method (Method A) takes the material of the building into account.

1988 UBC  $T = C_t (h_n)^{3/4}$  (3.2)  $C_t$  = 0.020 for dual moment resisting frames and eccentric braced frames;  $C_t = 0.1 / \sqrt{A_c}$  for structures with concrete or masonry shear walls;  $A_c = \sum A_d [0.2 + (D_e/h_n)^2]$  and depends upon the dimensions of the shear walls;

Ae is the minimum cross sectional shear area in any horizontal plane in the first story, in square feet, of a shear wall. D<sub>e</sub> is the length in feet of shear wall in the first story in the direction parallel to the applied forces.  $A_c$  is the combined effective area, in square feet, of the shear walls of the first story of a structure.  $h_n$  is the effective height in feet above the base of the building [UBC 1988]. The effective height of the building,  $h_n$ , accounts for reduction in floor area at high floor levels, roof equipment, and elevator housings on the roof. The second method (Method B) specifies a Rayleigh-Ritz approach to determine the period.

East Coast experience has led to less stiff structures because of moderate lateral load requirements as well as the tendency of eastern construction to be taller and more slender. The proposed New York City seismic provisions modify the 1988 UBC Code by assigning  $C_t$  = 0.035 for frame concrete structures and by adjusting  $C_t$  for dual systems as follows:

 $C_t = 0.020$  when  $h_n \leq 160$  ft.

 $C_t = 0.030$  when  $h_n \ge 400$  ft.

Linear interpolation is permitted for intermediate heights.

This method results in shorter periods than those measured in this study. Table 3-I summarizes these results.

BLDG.	Height <u>հռ</u>	1982 <b>UBC</b>	1988 UBC (A)	Proposed $\overline{\mathbf{N}\mathbf{Y}}\overline{\mathbf{C}}$	N-S	<b>Measured Ambient</b> E-W
$\mathbf A$	486'	5.3 s	2.07 s	$3.11$ s	4.47 s	4.53 s
B	423'	4.1 s	1.87s	2.80 s	3.98 s	2.79 s
$\mathbf C$	249'	2.8 s	$1.25$ s	1.41 s	1.94 s	1.69 s
D	460'	4.9 s	1.99 s	2.98 s	3.76 s	3.88s

TABLE 3-1 Comparison of Code-based Estimation of Fundamental Periods to the Measured Values in Ambient Conditions.

In comparing periods from code equations with measured ambient periods, it should be noted that the equations in the design code are meant to describe the structure in its ultimate state, for example, story drift of H/200, whereas the measured ambient periods correspond to story drifts of much less than H/1000. Stiffness degradation at larger story drifts will result in longer periods than those measured when the building is in ambient conditions.

Note that the code equations are really not meant to provide accurate estimates of the fundamental period, but for finding the proper design force level, as prescribed by a design spectrum. To be conservative, the periods from code formulas should therefore be less than the measured periods under ambient conditions, which applied in the 1988 UBC Code and the proposed NYC Seismic Code. (See Table 3-1.)

 $\frac{1}{2} \sum_{i=1}^n \frac{1}{2} \sum_{j=1}^n \frac{1}{2} \sum_{j=$  $\mathbb{F}$  $\mathbb{L}$  $\bar{1}$  $\mathbb{F}$  $\mathcal{A}^{\mathcal{A}}$  $\mathcal{L}^{\text{max}}_{\text{max}}$  and  $\mathcal{L}^{\text{max}}_{\text{max}}$  $\mathbb{L}$  $\mathsf{I}$  $\label{eq:2.1} \frac{1}{\sqrt{2}}\left(\frac{1}{\sqrt{2}}\right)^{2} \left(\frac{1}{\sqrt{2}}\right)^{2} \left(\$  $\bar{\mathbb{L}}$  $\mathbf{I}$  $\mathbf{L}$  $\|$  $\mathbb{F}$  $\|$  $\mathbb{F}$  $\mathbf{I}$  $\|$ 

## **SECTION 4 INSTRUMENTATION AND MEASUREMENT PROCEDURES**

In all four buildings, Terra Technology SSA-302 triaxial force-balance accelerometers were the sensors used to measure horizontal and vertical accelerations at various floor levels. Kinemetrics 5-second seismometers were used to measure the velocity motion of the foundation. The Terra Tech accelerometers exhibit excellent linearity by virtue of the feed-back control of the seismic mass. In addition, the fluid damped design makes them very rugged. Triaxial accelerometers were placed at different floor levels in the buildings. Two triaxial accelerometers on the same floor near the top of the building allowed the identification of torsional modes. In buildings A and B, the deployed accelerometers measured to  $\pm 0.2g$  full scale. In buildings C and D, the accelerometers measured to  $\pm 2.0g$  full scale. PRS-4 synchronous, digital, 12-bit, auto-scaling, battery powered recorders sampled, digitized, and recorded the data. Each recorder has 1 Mega-byte of RAM for data storage. There was one PRS-4 recorder per triaxial accelerometer. The PRS-4's are especially well suited for temporary instrumentation of large scale structures because they have synchronized  $(\pm 1$  millisecond) internal clocks, therefore they can operate independently from one another, do not require any interconnecting cables, and are ruggedized. The PRS-4 units were preprogrammed to sample simultaneously at a pre-specified time of day, sample rate, and duration. The wind was very light during the measurements of buildings A and B but was 35 knots during the measurements of buildings C and D. This is evidenced by much heavier participation of the lower frequencies for buildings C and D. The following table summarizes the instrumentation location and the data acquisition on the four structures. In Table 4-1, B stands for basement, E and W for east and west positions of the floor plan. Sensor locations not in either the east or west positions were placed near the center of the floor plan.

BLDG.			SENSOR LOCATIONS - FLOOR NUMBER		SAMPLE RATE TIME				
$\mathbf{A}$	B	10	-20	30 <b>SEP</b>	40	52E		$52W$ 50/sec	$20 \text{ min.}$
B	B.	10	<b>20</b>	30	36E	36W	40	$25/\text{sec}$	40 min.
$\mathbf C$	B	6.	28	29E	29W			$25$ / sec	30 min.
D	B.	10	51	51E 51W				$25/\text{sec}$	30 min.

**TABLE 4-1 Measurement Locations and Data Acquisition**

Measurements from Building A were taken on February 4, 1991; those from Building B were taken on February 7, 1991. Buildings C and D were measured on February 22, 1991. Buildings A and B were measured a second time on August 15, 1991. Appendix B contains the results from the first set of measurements. For completeness, the results from the second set of measurements are included in Appendix C.

After each test, the data was transferred via a Compaq portable PC to a network of Sun Spark stations for analysis. After formatting the data into the proper ASCII format, auto-power and phase spectra were estimated. Peak picking software facilitated the spectral analysis.

# **SECTION 5 DATA ANALYSIS AND SPECTRAL PARAMETER ESTIMATION**

Ideally, estimation of the parameters of a system is accomplished by fitting a model of the system to simultaneously sampled input and output data records. While the costs associated with artificial, measurable excitation of civil engineering structures are high, parameters estimated from forced-vibration tests are accurate and indicative of the nature of the structure itself [Ho 1989]. The force-levels required for uniform excitation of high-rise structures requires expensive and elaborate test arrangements, including reaction-mass servo-hydraulic actuators in many cases. Furthermore, forcing the structural response to levels significantly greater than ambient levels may risk non-linear behavior and subsequent damage to the structure, if measurement feedback is not used to limit the actuator excitation. Because of these experimental difficulties, and in light of the relative ease of collecting ambient vibration measurements as described in Section 4, the four high-rise structures were measured in ambient conditions.

If certain assumptions regarding the characteristics of the excitation to the structure, the structural behavior, and the response measurements hold, then the analysis of ambient vibration measurements will yield more meaningful results [Luz, 1992]. These assumptions are:

- 1. The excitation is sufficiently broad band and stochastic to uniformly excite at least the lowest 20 to 30 modes of the three-dimensional structure. **In** each of these structures the lower resonant frequencies range from 0.2 Hertz to 5 Hertz. Spatially uncorrelated excitations also helps reduce the chances of disproportionate energy concentration in certain modes.
- 2. The measured acceleration responses are at least weakly stationary. This permits meaningful time averages of the auto-power and cross-power spectra. The response

will be stationary if the excitation (wind pressure) is stationary, and the structure behaves linearly.

- 3. The resonant modes are lightly damped and the resonant frequencies are distinct and well separated. **In** this condition, spectral peak amplitudes have negligible contribution from out-of-band resonances. The figures in Appendices Band C verify that these assumptions hold.
- 4. Samples of the vibration measurements are simultaneous in time at all measurement locations. This condition is essential for proper cross-power spectral estimation. The test hardware guarantees nearly simultaneous sampling, as is indicated in Section 4.

If the above assumptions are warranted, frequencies with maximum spectral power (peak frequencies) can be interpreted as structural natural frequencies. **In** ambient conditions, tall buildings in urban areas are excited by wind turbulence and mechanical equipment. Wind pressures are spatially and temporally less correlated in turbulent flow than in laminar flow, but the spectra of turbulent velocity fluctuations decreases logarithmically with frequency [Simiu and Scanlan 1986]. Wind pressures are not uniformly distributed along the height of these high-rise structures. And because many of these buildings rise above the surrounding structures, the spectrum of wind velocites near the tops of the buildings can be qualitatively different from those in the lower stories. Since the auto-power spectra can not be normalized by the (unmeasured) wind forces, the spatial distribution of structural vibrations corresponding to a particular modal frequency are termed "operational deflection shapes." An operational deflection shape (ODS) is defined as the spatial distribution of peak spectral amplitudes at a resonant frequency.

Parameter estimation methods that implement an input-output ARMA model of the system were tested for ambient vibration applications [Ghanem, 1991]. However, the parameters resulting from such methods often depend upon assumptions regarding initial conditions, excitation, and
measurement noise. **In** addition, these methods are more computationally intensive than methods based on the Fast Fourier Transform (FFT). Frequency domain analysis is especially useful for ambient vibration data, since the long durations of data allow for a considerable amount of averaging while maintaining a very high resolution of frequencies. Small frequency increments,  $\Delta f$ , are necessary to characterize low frequency peaks. Computer programs tailored for spectral estimation from multi-measurement ambient data were written to first estimate power and phase spectra, then use these spectra to pick the peak frequencies, amplitudes, and to estimate damping, RMS acceleration, RMS velocity, and RMS displacement. To verify the operation of the programs described above, an almost identical method was implemented independently at Columbia University's Lamont-Doherty Geological Observatory. Consistent results confirmed the accuracy of the programs.

#### **5.1 Spectrum Estimation**

Acceleration data are digitized and sampled with a sample interval of  $\Delta t$  seconds.  $a(t) = a(i\Delta t)$ , i=1...P, where P is the total number of points in the digitized data record. The data record is divided into K overlapping sample functions,  $a_k$ , k=1...K, of N (a power of 2) points each. The sample functions overlap each other by 50% of their length.

$$
a_{k}(i\Delta t) = a((j+i)\Delta t), \quad j = \frac{N(k-1)}{2}, \quad k = 1 \cdots K, \quad i = 1 \cdots N
$$
\n(5.1)

Each sample function is multiplied by the Hanning window, w .

$$
\widetilde{a}_k(i\Delta t) = a_k(i\Delta t) w_i \tag{5.2}
$$

$$
w_i = \frac{1}{2} - \frac{1}{2}\cos\left(\frac{2\pi(i-1)}{N-1}\right)
$$
\n(5.3)

The complex valued discrete Fourier transform (OFT) of each segment is computed using the Cooley-Tukey (decimation in time) algorithm with the Danielson-Lanczos lemma to calculate the transforms [Oppenheim 1979, Press 1988]. This radix-2 method of computing the discrete Fourier transform is one of a variety of Fast Fourier Transform (FFT) methods. The Fourier transform is defined for both positive and negative frequencies. K FFT's are computed from each of two time records: a response acceleration record,  $y(i\Delta t)$ , at an elevated floor; and a reference acceleration record,  $x(i\Delta t)$  near the base of the building. For proper phase spectrum estimation, it is essential that the samples for the x and y time records be simultaneous. Each record is segmented and windowed as the FFT's,  $\tilde{X}_k$  and  $\tilde{Y}_k$ , are computed. Windowing the data reduces the leakage error in the FFT.

$$
\widetilde{X}_{k}(j\Delta f) = \sum_{i=1}^{N} \widetilde{x}_{k}(i\Delta t)e^{-i2\pi i j/N}
$$
\n(5.4)

$$
\widetilde{Y}_{k}(j\Delta f) = \sum_{i=1}^{N} \widetilde{y}_{k}(i\Delta t)e^{-i2\pi i j/N}
$$
\n(5.5)

Where the italic " $i$ " in the exponent is the imaginary number and should not be confused with the non-italic index, "i." The discrete frequency interval is  $\Delta f = 1/(N\Delta t)$ , and  $j=1...N/2$ . The discrete, one-sided, auto power spectrum,  $G_{yy}(j\Delta f)$ , and cross power spectrum,  $G_{xy}(j\Delta f)$ , are computed by averaging the auto-power and cross power spectra of the K overlapping sample functions.

$$
G_{yy}(j\Delta f) = \frac{1}{NK\sum_{i=1}^{N} w_i^2} \sum_{k=1}^{K} \left[ \left| \widetilde{Y}_k(j\Delta f) \right|^2 + \left| \widetilde{Y}_k(-j\Delta f) \right|^2 \right]
$$
(5.6)

$$
G_{xy}(j\Delta f) = \frac{1}{NK\sum_{i=1}^{N} w_i^2} \sum_{k=1}^{K} \left[ \widetilde{X}_k(j\Delta f) \widetilde{Y}_k^*(j\Delta f) + \widetilde{X}_k^*(-j\Delta f) \widetilde{Y}_k(-j\Delta f) \right]
$$
(5.7)

The asterisk (\*) indicates the complex conjugate of the variable. The factor in the denominator of equations (5.6) and (5.7) results in a normalization such that

$$
\sum_{j=1}^{N/2} G_{yy}(j\Delta f) = \frac{1}{P} \sum_{i=1}^{P} [y(i\Delta t)]^2
$$
\n(5.8)

The phase spectrum is calculated using the cross power spectrum with respect to the lowest measured floor (usually the tenth in these buildings). The discrete phase spectrum,  $\theta(j\Delta f)$  is defined by

$$
\theta(j\Delta f) = \arctan\left(\frac{\text{Re}\left(G_{xy}(j\Delta f)\right)}{\text{Im}\left(G_{xy}(j\Delta f)\right)}\right)
$$
\n(5.9)

The phase spectrum evaluated at peak frequencies of the auto power spectrum is almost always within  $5^{\circ}$  of either  $0^{\circ}$  or  $180^{\circ}$ . The power and phase spectra are given in Appendix B. The results of an analysis of a second set of measurements for Buildings A and B are included in Appendix C. Program AMB, listed in Appendix D, computes the auto-power spectrum and the phase spectrum given a pair of vibration records.

### **5.2 Parameter Estimation**

Structural resonant frequencies are estimated in each auto spectrum by curve-fitting a quadratic to the three largest values in the regions of peak spectral power, and calculating the coordinates of the peak of the quadratic. The curve-fit is exact since three points uniquely determine a quadratic; no least squares or maximum likelihood methods are needed for this estimation. Since the power and phase spectra are not transfer or transmittance functions, traditional modal analysis curve-fitting methods [Ho 1989, Richardson 1982, Richardson 1985, VoId 1990] cannot be employed, nor can many other methods which rely on an input-output model of the structure [DiPasquale 1987, Ewins 1984, Ghanem 1991, Ho 1989]. Nevertheless, since points in the power spectrum are usually asynunetrically spaced around the peak frequency, the frequency of the true largest power spectrum value usually does not coincide with the sampled peak frequency; the true peak frequency is usually within one frequency interval of the largest data value (see Figure 5-1.) By fitting a quadratic to the three largest points in the spectrum, the peak coordinates may be estimated easily.



**Figure 5-1 Quadratic Curve-fit to the Auto-power Spectral Peak**  $G_{yy}(f) = c_0 + c_1f + c_2f^2$ Peak Frequency:  $f_p = c_1 / (2c_2)$ , Peak Amplitude:  $G_{yy}(f_p) = c_0 + c_1^2 / (4c_2)$ . The variables  $a^2$  and  $b^2$  are used in Appendix A.

Given the coordinates of the largest sampled peak,  $(f_2, G_{yy}(f_2))$ , and the two adjacent coordinates,  $(f_1, G_{yy}(f_1))$  and  $(f_3, G_{yy}(f_3))$ , a quadratic,  $G_{yy}(f) = c_0 + c_1f + c_2f^2$ , passing through all three points is computed by solving

$$
\begin{pmatrix} G_{yy}(f_1) \\ G_{yy}(f_2) \\ G_{yy}(f_3) \end{pmatrix} = \begin{pmatrix} 1 & f_1 & f_1^2 \\ 1 & f_2 & f_2^2 \\ 1 & f_3 & f_3^2 \end{pmatrix} \begin{pmatrix} c_0 \\ c_1 \\ c_2 \end{pmatrix}
$$
\n(5.10)

for  $c_0$ ,  $c_1$ , and  $c_2$  using Gauss-Jordan inversion with partial pivoting. The estimated peak frequency is  $f_p = c_1 / (2c_2)$ . The estimated peak amplitude is  $G_{yy}(f_p) = c_0 - c_1^2 / (4c_2)$ .

The structural damping ratio is calculated using the band-width of the auto power spectrum. The peak coordinate estimated from the curve-fit,  $G_{yy}(f_p)$ , and the auto power spectrum data at a point near the half-peak level,  $G_{yy}(f_b)$ , are used to compute the damping ratio (refer to Figure 5-1). Band-width formulae for computing damping ratios are derived from the analytic expression of the dynamic amplification factor of a SDOF structure under steady-state harmonic motion (for details see Appendix A). The amplification factor is defined in terms of displacements and assumes viscous damping [Beards 1983, Paz 1985]. Since our power spectra are computed from acceleration data, and the normalization is given by equation (5.8), dividing each point by  $(2\pi f)^4$  results in displacement spectra. Assuming that the structure behaves linearly, has viscous damping, and has a stationary, steady-state, narrow-band response, that acceleration power spectra are normalized as in equation (5.8), and that the computed auto-power spectrum has no bias or leakage errors, then the following expression results in an approximate damping estimate.

$$
\zeta = \frac{|f_{b} - f_{p}|}{f_{p} \sqrt{\left(\frac{f_{b}}{f_{p}}\right)^{4} \frac{G_{yy}(f_{p})}{G_{yy}(f_{b})} - 1}}
$$
\n(5.11)

The curve-fit peak coordinate is  $(f_p, G_{yy}(f_p))$ ; the coordinate at the band-width level is  $(f_p)$ ,  $G_{yy}(f_b)$ ). The band-width frequency coordinate,  $f_b$ , is always taken to be greater than the peak frequency,  $f_p$ ; and the band-width amplitude is no greater than 70% of the peak amplitude. This method of frequency estimation and damping calculation allows for the frequency data to be asymmetrically spaced around the peak and uses exact data values when appropriate. Equation (5.11) is a conservative approximation which ignores higher order terms (see Appendix A). Recently, FFT-based averaged spectrum estimation methods have been applied to the problem of damping estimation of long span suspension bridges [Jones 1990, Littler 1991].

For finite lengths of data, there is a trade-off between the error in the spectral amplitudes and the spectral resolution of any FFT-based spectrum estimation routine [Gade 1988]. The variance of a spectrum with fine spectral resolution is larger than the variance of an averaged spectrum with coarse spectral resolution computed from a data set of a fixed length [Press 1988]. The normalized standard error of a spectral estimate is  $K^{-1/2}$ , where K is the number spectral averages [Bendat 1986]. However, spectral bias errors result from averaging spectra and have the effect of widening spectral peaks. An estimate of the bias error of a spectral estimate is  $E_b = -\Delta f/(6\zeta f)^2$  where  $\Delta f$  is the frequency resolution of the spectrum, and f and  $\zeta$ are the frequency and damping ratio associated with a peak in the spectrum [Bendat 1986]. The bias can be somewhat controlled by the shape of the window function. In choosing a window function, there is a trade off between the narrowness of the peak and side-lobe amplitudes; so windows that produce narrow peaks are not always advantageous. Other spectrum estimation methods, employing multiple orthogonal window functions, have been shown to reduce bias errors more effectively [Thompson 1982, Park 1987].

In an effort to find the combination of the window function and number of averages which gives a spectrum with an appropriate band-width, several approaches were employed. All approaches used the same simulated acceleration response record of a linear SDOF system of a known natural frequency and damping ratio. The linear acceleration method was used to simulate the acceleration response record [Paz 1985]. Spectra of the resulting simulated acceleration response record were computed using various window functions and numbers of averages. Method 1 simply compared the known damping ratio to the damping ratio computed using equation (5.11). Methods 2 and 3 compared the bandwidth factor, q, computed from a closed form solution to the band-width factors computed from spectral moments, equation

(5.13), and from envelope statistics, equation (5.14). Three expressions of the band-width factor, q, are [Vanmarcke 1972]

$$
q^{2} = 1 - \frac{1}{1 - \zeta^{2}} \left[ 1 - \frac{1}{\pi} \tan^{-1} \left( \frac{2\zeta \sqrt{1 - \zeta^{2}}}{1 - 2\zeta^{2}} \right) \right]^{2}
$$
  
q<sup>2</sup> = 1 -  $\frac{\lambda_{1}^{2}}{1 - \zeta^{2}}$  (5.12)

$$
A^2 = 1 - \frac{1}{\lambda_0 \lambda_2} \tag{5.13}
$$

$$
q = \frac{\sigma_f}{\sigma_x} \tag{5.14}
$$

where  $\zeta$  is the damping ratio, and  $\lambda_i$  is the i<sup>th</sup> spectral moment of the auto-power spectrum.

$$
\lambda_{i} = \sum_{j=1}^{N/2} (j\Delta f)^{i} G_{yy}(j\Delta f)
$$
\n(5.15)

 $\sigma_{\rm r}$  and  $\sigma_{\rm x}$  are computed using envelope statistics.  $\sigma_{\rm r}$  and  $\sigma_{\rm x}$  are the standard deviations of the time rates of change of the random processes x and r respectively. The envelope of the random process is r, and x is computed as the modulous of the complex quantity  $x + i\hat{x}$ , where  $\hat{x}$  is the Hilbert transform of x and can be computed in the frequency domain or in a finite impulse response filter form in the time domain [Bendat 1986, Boashash 1987, Deutsch 1969, Rao 1990]. The impulse response function of the Hilbert transform is

$$
h(t) = \frac{1}{\pi t} \tag{5.16}
$$

The Fourier transform of  $\hat{x}(t)$  is  $\hat{X}(f)$  and the Fourier transform of  $x(t)$  is  $X(f)$ .

$$
\widehat{X}(f) = \begin{cases}\ni X(f) & \text{if } f > 0 \\
0 & \text{if } f = 0 \\
-i X(f) & \text{if } f > 0\n\end{cases}
$$
\n(5.17)

After computing  $\hat{X}(f)$  from equation (5.17),  $\hat{x}(t)$  is computed using the inverse Fourier transform of  $\hat{X}(f)$ . The time derivatives of the processes are computed using the central difference rule.

$$
\dot{x}_j = \frac{x_{j+1} - x_{j-1}}{2\Delta t}
$$
\n(5.18)

Figure 5-2 shows a portion of the simulated acceleration response record, its envelope function, and the power spectrum of the simulated acceleration record.

No combination of window function and frequency interval resulted in a very good match for all three methods of band-width comparison simultaneously. The final decision on what level of averaging and window function to use was based largely on the results from Method 1. Since the frequency values are of much greater interest in this study, a finer spectral resolution was favored. From these considerations, a frequency interval of approximately 0.0061 Hz and the Hanning window function were chosen to be suitable. This resulted in 13 spectral averages for Building A, 28 spectral averages for Building B, 20 averages for Building C, and 20 for Building D. A more rigorous analysis would take advantage of multi-taper methods [Park 1987, Thompson 1982] and would be an interesting topic for further study. These methods result in spectral estimates which have much smaller bias errors.

In summary, spectral band-widths can be determined to a large extent by computational errors in peak regions of the spectrum. If the auto-power spectrum is calculated with little regard to the sources of these errors, damping estimates from band-width measures may have little physical significance. In addition, damping estimates from band-width measures of ambient vibration data are contingent upon implicit assumptions, are subject to data processing effects, and therefore are used only as comparative measures within this analysis.

#### 5.3 **Root-Mean-Square Computations**

Since the auto power spectra are normalized as in equation (5.8), the square root of the sum of the power spectrum points is the root-mean-square acceleration of the record, provided that the

time history was demeaned or high-pass filtered prior to power spectrum computation. More importantly, however, root-mean-square velocity or displacement can be obtained if each spectral data point is divided by  $(2\pi f)^2$  or  $(2\pi f)^4$ , respectively. Root mean squared displacement computations in the frequency domain are contingent upon the following assumptions:

- 1. The original discrete acceleration record can be represented by a Fourier series expansion.
- 2. Each power spectrum coordinate equals the amplitude squared of the corresponding Fourier series term.
- 3. Term by term integration of the Fourier series representing the acceleration record results in a Fourier series representing the velocity record. And term by term integration of the Fourier series representing the velocity record represents the displacement record.

Root-mean-squared (RMS) displacement computations are most sensitive to low-frequency accelerations. Power spectra often contain low frequency biases due to low-frequency noise in the measured acceleration records. Since low-frequency biases in the power spectrum result in gross over-estimates of RMS velocity and especially RMS displacement, all power spectrum values below 1 rad/sec (0.159 Hz) are divided by 1 instead of  $(2\pi f)^2$  or  $(2\pi f)^4$ , thereby attenuating the very low-frequency noise. (For small values of f, the function  $1/(2\pi f)^4$ ) becomes very large.) 1 rad/sec (Period = 6.28 sec) is a reasonable lower bound for structural oscillations of the buildings in this study. The equations for frequency domain calculation of the root mean square equations are:

RMS-acc = 
$$
\left[\sum_{j=1}^{N/2} G_{yy}(j\Delta f)\right]^{\frac{1}{2}}
$$
 (5.18)

RMS-vel = 
$$
\left[\sum_{j=1}^{N/2} \frac{G_{yy}(j\Delta f)}{\omega^2}\right]^{\frac{1}{2}}
$$
 (5.19)

$$
RMS\text{-}dsp = \left[\sum_{j=1}^{N/2} \frac{G_{yy}(j\Delta f)}{\omega^4}\right]^{\frac{1}{2}}
$$
(5.20)

where

$$
\omega = \begin{cases}\n2\pi j \Delta f & \text{if } j\Delta f \ge 0.159 \\
2\pi 0.159 & \text{if } j\Delta f < 0.159\n\end{cases}
$$
\n(5.21)

Alternately, integrating acceleration time histories numerically results in velocity and displacement time histories. Numerical integration can be unstable if the record to be integrated contains frequencies greater than  $1/(4\Delta t)$  Hz [Hamming 1989]. To estimate RMS displacement from recorded acceleration time records, the authors found that the following procedure was sufficiently accurate:

- 1. Demean the acceleration time record by subtracting the average acceleration from all data points.
- 2. Low-pass filter the acceleration time record with a cut-off frequency at  $1/(5\Delta t)$  Hz. A non-recursive Kaiser filter was implemented in these analyses [Hamming 1989].
- 3. Integrate the demeaned, filtered, acceleration data once using Tick's rule [Hamming 1989]. Tick's rule is accurate and stable up to  $1/(4\Delta t)$  Hz.
- 4. Band-pass filter the integrated acceleration record between 0.15 Hz and  $1/(5\Delta t)$  Hz.
- 5. Integrate again using Tick's rule and find the RMS value of the 'displacement' time record.

These methods were verified in the laboratory by measuring and recording simultaneous acceleration and displacement time histories of a SDOF system excited randomly by base accelerations. A linear voltage displacement transducer (LVDT) measured the mass's deflections while a force balance accelerometer measured the mass's accelerations. RMS displacement was computed in three ways: directly from the recorded displacement time records, in the frequency domain using the acceleration power spectrum and equations (5.20) and (5.21), and in the time domain from the five-step process outlined above. The resulting RMS values were within 2% of each other. Comparable results between the time and frequency domain methods using the ambient vibration data collected from the measured structures in New York City provided further confirmation of the accuracy of the frequency domain approach. Table 5-1 compares the measured RMS displacement and the computed RMS displacements for the SDOF structure tested in the laboratory. RMS displacement values computed from the 52nd floor of Building A are indicated in the second row of Table 5-1.

**TABLE S-I Validation of the Frequency Domain RMS Computation Method**

	Measured Displacement	Frequency Domain Time Domain	
<b>SDOF-RMS</b>	$0.591$ in.	$0.588$ in.	$0.578$ in.
$NYC - RMS$		$0.112$ cm	$0.130 \text{ cm}$

To facilitate the analysis of the power spectra and phase spectra computed by program AMB, program PEAK was written. This program passes through the spectral data several times to compute RMS quantities, pick peaks, sort the peaks by increasing frequency, fit a quadratic to the peaks, and estimate values for the peak frequencies, amplitudes, and damping ratios. To qualify as a peak, the power spectrum data must satisfy all of the following conditions:

- 1. The power spectrum frequency coordinate must lie within a user-specified frequency band.
- 2. The amplitude must be larger than any other point not already identified as a peak.
- 3. The derivative of the spectrum with respect to frequency must change signs at a peak.
- 4. The frequency at a peak must be outside of a user-specified frequency buffer around any previously identified adjacent peak. This helps in eliminating picking large sidelobes as legitimate peaks.

In estimating the damping ratio, the spectral coordinate used for the band-width estimation is no greater than 70% of the peak amplitude. The source code for program PEAK is included in Appendix E.

**FIGURE 5-2** Simulated Acceleration Response, Computed Envelope, and Estimated Power Spectrum



Simulated Acceleration & Envelope

Power Spectrum

 $\mathbf{I}$  $\overline{1}$  $\overline{1}$  $\mathbf{I}$  $\mathbf{I}$ 

# **SECTION** 6 **DYNAMIC PARAMETERS**

Table 6-1 compares the periods and amplitudes of the ambient measurements resulting from an analysis using the methods described in Section 5 with those from finite-element dynamic analysis, and with code-formula estimates. The finite-element analysis is preliminary, reflecting engineering judgment as to the flat-slab width, effective shear wall stiffness, and effective coupling-beam moment of inertia.

The wind was very light during the measurements of Buildings A and B but was 35 knots during the measurements of Buildings C and D. This is evidenced by heavier modal participation in the lower frequencies for Buildings C and D. The concentration of spectral energy in the lower frequencies in Buildings C and D made the higher frequencies in Buildings C and D more difficult to identify.

Lengthy ambient vibration records improved power spectrum estimation by allowing a fine frequency resolution while averaging several spectra together. Most spectra were 2048 points long and were averaged up to 28 times. This enhanced the estimation of the natural periods, and damping ratios. The output files from program PEAK, which analyzed the power spectrum and phase spectrum data files, are shown with their corresponding spectra in Appendix B.

Figures 6-1 to 6-10 illustrate the "operational deflected shapes" (ODS) of the four buildings. Note that these are not normalized mode shapes since the spectral amplitude of ambient vibration measurements cannot be normalized by the spectral excitation amplitudes. "Operational deflected shapes" is the term used by the modal analysis community to describe the spatial description of the amplitude of vibration at resonant frequencies under operating conditions. Operational deflection shapes corresponded well to engineering judgment. That is, for the most part, the number of inflection points in the ODS is one less than the number of the associated resonant frequency for that ODS. Furthermore, ODS amplitudes are larger in the lower resonant frequencies.

The root mean squared (RMS) displacement values presented in Table 6-1 are not the first mode RMS values, but are computed using equation (5.20). RMS velocity and acceleration values, along with plots of acceleration response amplitude and phase spectra are given in Appendices Band C. Measured RMS accelerations were much less than the threshold of perception  $(0.005 \text{ g or } 5 \text{ cm/sec}^2 \text{ [Similar and Scalar 1986]}).$ 



### **TABLE 6-1 Dynamic Parameters Estimated by Ambient Vibration Measurements, by Finite Element Analysis, and by Code.**

\* Computed periods were obtained using  $K_d=2.0$  in equation (7.1) for flat slabs and uncracked sections for columns, walls and beams [Grossman, in progress].

\*\* N/A indicates that these values are not available, this work is in progress by the designers.

Figure 6-1 **Building A; North-South Accelerations** 60 50 4.47 s 40  $1.32 s$ **FLOOR**  $0.71 s$  $30\,$  $0.61~\mathrm{s}$  $0.49\ {\rm s}$ 20 ÌЦ.  $0.40\;{\rm s}$  $0.37~\mathrm{s}$  $10\,$  $0<sub>1</sub>$  $0.00$  $0.01$  $-0.01$ 0.02 **Dynamic Floor Acceleration (cm/s/s)** Figure 6-2 **Building A; East-West Accelerations** 60 50 40 4.53 s **FLOOR**  $1.23 s$ 30  $0.64 s$  $0.40 s$  $20\,$  $0.28s$ 10  $\theta$  $-0.006$   $-0.004$   $-0.002$   $0.000$   $0.002$   $0.004$   $0.006$   $0.008$ **Dynamic Floor Acceleration (cm/s/s)** 





6-5



 $6-6$ 









# **SECTION 7 RESISTANCE OF FLAT PLATE SYSTEMS TO LATERAL FORCES**

In high-rise buildings that are constrained architecturally from the prevalent usage of windowless structural walls, buildings that have columns which do not fall in a regular grid pattern, and buildings that have a structural system determined by interior floor-plans, flat plates have been designed to resist lateral forces. Forming costs for flat plates are also less those for beam-drops. However, the lateral resistance that flat plates provide is not clearly understood. The approach engineers should choose in the design of flat plate systems, when the flat plate is intended to provide lateral resistance, is not precisely specified in ACI 318-83 [ACI-318 Code 1983]. The Equivalent Frame Method was evaluated by the late Professor Vanderbilt [Vanderbilt 1981]. But disagreement in the 1983 ACI Code Committee 318 regarding its suitability resulted in the inclusion of both Equivalent Frame and Effective Width methods in the ACI 318-83 code. (See commentary [ACI-318 Commentary 1983] Section 13.3.1.2.) A lack of pertinent test data caused great uncertainty in the detailing of connections between plates and supporting columns in order for the plate to develop its full moment capacity. The discussions during the development of the 1983 Code concerning flat plate design for lateral forces led the Reinforced Concrete Research Council (RCRC) to assign researchers to tasks aimed at better understanding flat plate behavior. One of the tasks was motivated by the recent inclusion of moderate seismic requirements in many states East of the Rocky Mountains.

A methodology describing the effective width,  $\alpha l_2$ , of flat plates at interior supports [Grossman, 1987] was verified and adjusted [Grossman, in progress] to results of laboratory tests [Moehle 1990] is as follows (for notation consult ACI 318):

$$
\alpha l_2 = K_1 \left[ 0.3 l_1 + C_1 \frac{l_2}{l_1} + \frac{(C_2 - C_1)}{2} \right] \frac{d}{0.9 h}
$$
 (7.1)

with limits

$$
0.2 K_d l_2 \le \alpha l_2 \le 0.5 K_d l_2 \tag{7.2}
$$

where  $K_d$  depends on the load level and the computed drift index. Additional adjustments at the edge of the plate, and at side and corner supports were also recommended but not duplicated here. The  $K_d$  values were calibrated from the flat plate substructure tests [Moehle, 1990] and *tentatively* (pending additional field measurements) may be taken equal to:



A future goal of the current measurements and analyses of wind excited response of the four full-scale buildings will be to verify the appropriate values for  $K_d$  at smaller drift ratios (< H/800) and to confirm the values of  $K_d$  interpolated from the substructure tests when possible. Computed dynamic properties (using  $K_d=2.0$ ) for two measured structures [Grossman, in progress] are summarized in Table 6-1.

More intimate knowledge of the fundamental periods and damping in actual flat plate structures is needed to better estimate seismic lateral loads, soil-structure interaction, and serviceability of tall structures. These issues will bear more weight when the need to retrofit existing structures (not designed for seismic loads), becomes pertinent. The research presented herein is the initial phase of work toward the task of correlating the design methods developed for flat-plate structures and shear walls with measured ambient dynamic properties. This correlation will implement ETABS models of the "as-built" condition of the measured models. Douglas and Ried [1982] calibrated similar SAP models to the measured bridge dynamic parameters, and their methods will be implemented here.

The approach starts by deriving a quadratic relationship between M parameters  $(P_i, i=1...M)$ and N significant structural variables  $(X_k, k=1...N)$ . In this analysis, the parameters,  $P_i$ , are the natural periods. The structural variables,  $X_k$ , are the effective plate width, the effective shear wall area, and the effective coupling beam moment of inertia. M must be at least as large as N. Analytic parameters, obtained from an ETABS computation as a function of the structural variables, are denoted  $P_i$ ; while measured parameters, obtained from ambient dynamic measurements are denoted  $PM_i$ . The fitting of the ETABS model to the measured parameters,  $PM_i$ , by adjusting the structural variables,  $X_k$ , follows the following steps.

- 1. Specify a base-line value,  $X^b_k$ , an upper bound,  $X^u_k$ , and a lower bound,  $X^l_k$  for each structural variable.  $X^l_k < X^b_k < X^u_k$ .
- 2. Run ETABS 2N+1 times. For each run, each  $X_k$  changes from  $X_{k}^{b}$  to  $X_{k}^{l}$  to  $X_{k}^{u}$ one at a time. Each run computes different set of M parameters (periods). Parameters computed by varying the structural variables are denoted  $PQ<sup>i</sup>$ . There are three sets of parameters corresponding to each structural variable, corresponding to  $X^b_k$ ,  $X^l_k$ , and  $X^u$ <sub>k</sub> respectively.

$$
PQ_1^i = P_i(X_1^b \cdots X_N^b)
$$
  
\n
$$
PQ_2^i = P_i(X_1^1 \cdots X_N^b)
$$
  
\n
$$
PQ_3^i = P_i(X_1^u \cdots X_N^b)
$$
  
\n
$$
\vdots
$$
  
\n
$$
PQ_{2N+1}^i = P_i(X_1^b \cdots X_N^u)
$$
  
\n(7.3)

3. Derive a quadratic equation describing how each parameter,  $PQ<sup>i</sup>$ , varies with changes in the N structural variables. The quadratic form of  $PQ^i$  is given by

$$
PQ^{i} = C^{i} + \sum_{k=1}^{N} (X_{k}A_{k}^{i} + X_{k}^{2}B_{k}^{i})
$$
\n(7.4)

The constant polynomial coefficients,  $A_k^i$ ,  $B_k^i$ , and  $C^i$ , of the quadratic can be found by solving

$$
\begin{pmatrix}\nPQ_1^i \\
PQ_2^i \\
PQ_3^i \\
\vdots \\
PQ_{2N+1}^i\n\end{pmatrix} = \begin{bmatrix}\n1 & x_1^b \cdots x_N^b & x_1^{b^2} \cdots x_N^{b^2} \\
1 & x_1^1 \cdots x_N^b & x_1^{1^2} \cdots x_N^{b^2} \\
1 & x_1^u \cdots x_N^b & x_1^{u^2} \cdots x_N^{h^2} \\
\vdots & \vdots & \vdots \\
1 & x_1^b \cdots x_N^u & x_1^{b^2} \cdots x_N^{u^2}\n\end{bmatrix} \begin{bmatrix}\nC^i \\
A_1^i \\
\vdots \\
A_N^i \\
B_1^i \\
\vdots \\
B_N^i\n\end{bmatrix} (7.5)
$$

using, for example, singular value decomposition.

4. If each PQi can be put in terms of a quadratic function of the structural variables, the structural variables corresponding to the measured parameters can be found in a least squares sense. Using the polynomial expressions for  $PQ^i$  given by equations (7.4) and (7.5), an error function can be written

$$
E^2 = \sum_{i=1}^{M} (PM_i - PQ^i)^2
$$
 (7.6)

Equation (7.6) is minimized by setting the partial derivatives of  $E^2$  with respect to  $X_k$ equal to zero and obtaining N equations to solve for the N different structural variables,  $X_k$ . There are a variety of numerical minimization procedures which can accomplish this task.

5. Once a set of the structural variables has been identified, the procedure can be repeatedly refined by using the identified values of  $X_k$  as  $X_k$ <sup>b</sup> and diminishing the difference between  $X_k^u$  and  $X_k^l$ . This will refine and reduce the uncertainty in the true structural parameters,  $X_k$ .

By correlating ETABS models to natural periods measured under various wind conditions, including very strong winds, and by including recent static tests on flat-plate/column models, it is anticipated that criteria for degradation of ambient stiffness to stiffness at load levels anticipated during seismic events will be established. Top story displacements during strong winds may be on the order of H/800 or H/1000. Recent static tests on scaled flat plate substructures indicates the degradation of lateral resistance for drifts corresponding to HJ800, H/400, H/200 and larger. The laboratory study fine tuned an empirical formula, equation (7.1), relating the effective width at interior supports of the flat plate to the load level, causing a defined top story drift. The purpose of the research at hand is to confirm the laboratory results at lower load levels by identifying the effective plate width of full scale structures subjected to various load levels. Once verified, this stiffness degradation can be used in seismic vulnerability assessment of many flat-plate and shear-wall structures East of the Rocky Mountains.

 $\bar{\rm t}$ 

# **SECTION 8 CONCLUSIONS**

- 1. Vibration measurement analysis of high-rise structures in ambient conditions is useful in revealing the differences between the as-built characteristics of the structure and the design representation of the building. The analyses of measurements from four reinforced concrete flat plate structures in Manhattan ranging from 27 stories to 52 stories are presented in this report. The aspect ratios of the structures ranged from 2.1 to 5.3. The four structures were in various phases of construction at the time of the measurements.
- 2. Using modular, synchronized (to within 1 millisecond), digital data recorders simplifies the instrumentation of large scale structures. These recorders do away with the task of running cables between remote sensors and a central recording unit. Hence, measurements can be taken unobtrusively and with no interruption in construction. All the structures were measured with triaxial force-balance accelerometers in at least four locations. Data was collected at 25 or 50 samples per second.
- 3. Ambient vibration measurements were analyzed using software tailored for this purpose. The auto-power spectrum and phase spectrum are estimated using fast Fourier transforms and a windowing method. Interpretation of the spectral estimates was enhanced by peak-picking software, also tailored for ambient structural vibration analysis.
- 4. Recording long lengths of ambient vibration data allows for several spectral averages which enhances spectral resolution. Curve-fitting the spectral peaks with a quadratic further improves the precision of the estimated periods. Measured ambient fundamental periods from all four buildings are shorter than those given by the 1982 Uniform Building Code, but are longer than fundamental periods estimated using formulas in the 1988 Uniform Building Code, and the proposed NYC Seismic Code. Fundamental

periods range from 1.69 seconds in the shortest building to 4.53 seconds in the tallest building.

- 5. Ambient vibration analysis using spectral analysis methods is robust. However results derived from spectral band-width measures, such as damping ratios, should be viewed with extreme caution. Bias and leakage errors in the spectrum estimates affects bandwidth measures strongly. Damping estimates in ambient conditions reported in this study range from 0.2% in the higher modes to 5.6% in the fundamental mode.
- 6. Three methods for assessing the artificial widening of spectral peaks due to averaging and windowing finite lengths of discrete data are reviewed. These methods can be used as explicit guidelines in choosing target frequency intervals, and required data record lengths for a specific application. Comparing damping ratios computed from the spectra of simulated acceleration records with the known damping ratios is the most robust method for choosing the proper level of frequency resolution and spectral averages. Future analyses using multi-taper spectral estimation methods, in which the data is repeatedly windowed with orthogonal weighting functions, should result in spectra with less bias.
- 7. The computation of root mean square (RMS) acceleration and RMS displacement from normalized power spectra of acceleration data was confirmed using a small-scale laboratory test in which acceleration and displacement were measured and recorded simultaneously. Root mean square displacements range from 0.007 em under very light wind to 0.08 cm under stronger winds. The tallest buildings were measured only under light winds.
- 8. The first mode alone cannot account for the total motion of tall structures in ambient conditions. Indeed, the higher modes often participate more strongly in the lower floors. Dynamic measurements are very useful in evaluating the higher mode response of the structure. The participation of lower modes is heavier during strong winds. So

the total displacement response to strong wind can be estimated using the first mode alone within acceptable error margins.

- 9. Differences in periods estimated from ambient vibration measurements and preliminary finite element analysis can be reconciled by fitting uncertain quantities, such as effective flat-plate widths, in the finite element model. In so doing, the effective flat plate width can be related to lateral load level under service load conditions.
- 10. Fitting flat plate widths in the finite element model to measured periods can give a better understanding of the contribution of flat plates to the lateral resistance in tall buildings, and how this resistance changes with higher load levels.
- 11. The data presented herein is part of an on-going study of the behavior of flat-plate reinforced concrete structures. Fundamental periods from these tests will be compared to periods from future tests under different wind conditions or different stages of construction. Combining the results of these studies with the results of destructive tests of flat-plate substructures will lead to a method for evaluating the behavior of high-rise flat-plate structures at large drift ratios (H/200).

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 $\sim 60$  km s  $^{-1}$ 

 $\langle \hat{u}_\mathrm{max} \rangle$ 

 $\sim$ 

 $\sim 10^{-10}$ 

 $\label{eq:2.1} \frac{1}{2} \int_{\mathbb{R}^3} \frac{1}{\sqrt{2}} \, \frac{1}{\sqrt{2}} \,$
#### **APPENDIX** A

### **DERIVATION OF TWO DAMPING ESTIMATORS**

Textbook formulae for damping estimation using impulse response functions or spectral bandwidths should be used with caution when applied to experimentally obtained discrete data. This appendix contains the derivation of these two common damping estimators and assesses their suitability to experimental data from large-scale structures [Beards 1983, Paz 1985].

## **A.I Time Domain**

In the time domain, damping is often estimated using the logarithmic decrement. Consider the SDOF displacement response to an impulse.

$$
\mathbf{x}(t) = A_0 e^{-\omega_n \zeta t} \sin \left( \omega_n \sqrt{1 - \zeta^2} t - \varphi \right)
$$
 (A.1)

This is a sinusoidal oscillation with exponentially decreasing amplitude. The response displacement reaches adjacent relative maxima at times  $t_1$  and  $t_2$ , with displacements  $A_1$  and A<sub>2</sub>.  $t_2 - t_1 = \tau$ . The variable  $\tau$  is the damped natural period.

$$
\frac{A_1}{A_2} = \frac{e^{-\omega_n \zeta t_1}}{e^{-\omega_n \zeta (t_1 + \tau)}} = e^{-\omega_n \tau}
$$
\n(A.2)

By definition,

$$
\omega_{\rm n} \sqrt{1 - \zeta^2} \tau = 2\pi \tag{A.3}
$$

Defining the logarithmic decrement,  $\delta$ ,

$$
\delta = \ln \left( \frac{A_1}{A_2} \right) \tag{A.4}
$$

taking the natural log of both sides of equation (A.2), and substituting equation (A.3) results in an expression for damping.

$$
\zeta = \frac{\delta}{\sqrt{4\pi^2 + \delta^2}}\tag{A.5}
$$

If the damping is purely viscous then several oscillations may be used to compute the logarithmic decrement

$$
\delta = \frac{1}{n} \ln \left( \frac{A_1}{A_{n+1}} \right) \tag{A.6}
$$

Equation (A.5) is exact. In its derivation, no expressions were truncated and no terms were ignored. If the damping is viscous, then  $\delta$  will not depend upon the strain, deflection, temperature, or strain rate. In this case,  $\delta$  will be constant with respect to the section of the data record chosen for analysis. In addition, the logarithmic decrement may be computed using acceleration amplitudes directly if the response is purely single mode. However, several damping mechanisms usually contribute to the total damping and each mechanism's contribution usually is *not* independent of strain, deflection, temperature, and strain rate. Often, damping ratios computed with the initial, large motion, response are larger than those computed using the subsequent, smaller amplitude, motion. Further experimental complications arise from the fact that impulse decay measurements are often polluted with higher mode response, and that low frequency noise can seriously distort the results. Additionally, large, purely impulsive forces are difficult to generate in high-rise structures and damping ratios for higher modes are very difficult to obtain using this technique.

## **A.2 Frequency Domain**

Frequency domain damping evaluation, although simpler experimentally, requires many more mathematical assumptions than does evaluation of damping in the time domain. Consider the dynamic amplification of a SDOF system responding to steady state harmonic excitation of circular frequency  $\omega$ . If the response quantity is acceleration, the dynamic amplification factor is

$$
\frac{\ddot{x}}{x_{st}} = \frac{\omega^2}{\sqrt{\left(1 - \Omega^2\right)^2 + \left(2\zeta\Omega\right)^2}}
$$
\n(A.7)

where  $\Omega = \omega/\omega_n$ . The peak of the function described by equation (A.7) occurs when

$$
\Omega = \sqrt{1 - 2\zeta^2} \tag{A.9}
$$

and has the value of

$$
a = \frac{\ddot{x}}{x_{st}} = \frac{\omega^2}{2\zeta\sqrt{1 - \zeta^2}}
$$
(A.10)

The value of the acceleration frequency response function, equation (A.7), at the frequency  $\Omega = 1 + \Delta\Omega/2$  is  $\left(1 + \Delta\Omega\right)^2 \omega^2$ 

$$
b = \frac{1}{\sqrt{\left[1 - \left(1 + \frac{\Delta\Omega}{2}\right)^{2}\right]^{2} + \left[2\zeta\left(1 + \frac{\Delta\Omega}{2}\right)^{2}\right]^{2}}}
$$
(A.11)

Values for a2 and b2 are indicated graphically in Figure 5-1. Since the eventual aim is to compute damping estimates from power spectra, the ratio of the squares of equations (A.10) and  $(A.11)$  is

$$
\frac{a^2}{b^2} = \frac{\left(1 - 2\zeta^2\right) \left\{ \left[1 - \left(1 + \frac{\Delta\Omega}{2}\right)^2\right]^2 + \left[2\zeta\left(1 + \frac{\Delta\Omega}{2}\right)\right]^2 \right\}}{4\zeta^2 \left(1 - \zeta^2\right) \left(1 + \frac{\Delta\Omega}{2}\right)^4}
$$
\n(A.12)

Solving for  $\zeta^2$ , we obtain

 $\bar{z}$ 

$$
\zeta^2 = \frac{\left(1 - \varepsilon_2\right) \left(\Delta\Omega + \frac{\Delta\Omega^2}{4}\right)^2 + \frac{a^2}{b^2}\varepsilon_1\varepsilon_4}{4\left(\frac{a^2}{b^2}\varepsilon_4 + \varepsilon_3 - 1\right)}
$$
\n(A.13)

where

$$
\varepsilon_1 = 4\zeta^4 \tag{A.14}
$$

$$
\varepsilon_2 = 4 \left( \zeta - \zeta \right) \tag{A.15}
$$
\n
$$
\varepsilon_3 = \Delta \Omega + \frac{\Delta \Omega^2}{4} \tag{A.16}
$$

and

$$
\varepsilon_4 = \begin{cases} \left(1 + \frac{\Delta \Omega}{2}\right)^4 & \text{if the data is acceleration} \\ 1 & \text{if the data is displacement} \end{cases}
$$
 (A.17)

For small values of damping,  $\epsilon_1 < \epsilon_2 < \epsilon_3 < \epsilon_4$ . Neglecting higher order terms, such as  $\Delta\Omega^2$ with respect to  $\Delta\Omega$ , and both  $\Delta\Omega^2$  and  $\Delta\Omega$  compared to 1, and assuming the power spectrum data is from an acceleration record,  $\varepsilon_1$ ,  $\varepsilon_2$ , and  $\varepsilon_3$  are set to zero, and

$$
\zeta = \frac{\Delta\Omega}{2\sqrt{\frac{a^2}{b^2}\left(1 + \frac{\Delta\Omega}{2}\right)^4 - 1}}
$$
\n(A.18)

which is equivalent to equation  $(5.11)$ .

The truncated formula,  $(5.11)$  or  $(A.18)$ , results in lower damping estimates than using the exact formula,  $(A.13) - (A.17)$ . This under-estimate of the damping is recovered (in part if not entirely) by the artificial widening of the spectral peaks caused by windowing the time domain data before computing the FFf's [Oppenheim 1979, Press 1988]. Using equation  $(5.11)$  or  $(A.18)$  for damping estimation from power spectra of acceleration data is contingent upon the following assumptions:

1. The damping is small  $\ll 10\%$  and the damping mechanism is viscous. For damping  $\lt$ 2%, the truncation effects are insignificant.

- 2. The original discrete acceleration record can be represented by a Fourier series.
- 3. Each power spectrum coordinate equals the amplitude squared of the corresponding Fourier series term.
- 4. The response is narrow band, stationary, and steady-state and the structure behaves linearly.
- 5. Power spectrum computation does not appreciably widen the spectral peaks. This is often the over-riding assumption.

In summary, damping estimates from impulse responses are the most accurate, if an adequately strong and brief impulse can be generated. The absence of forced response experiments requires frequency domain estimates. These estimates are often upper bounds due to artificial peak widening in the process of computing averaged windowed power spectra.

# **APPENDIX B**

# **POWER SPECTRA, PHASE SPECTRA, ESTIMATED FREQUENCIES, AND ESTIMATED DAMPING RATIOS: FIRST SET OF MEASUREMENTS**

Parameters reported in Section 6 are estimated from the spectra shown in Appendix B. This data will serve as a base-line set of measurements for comparison to parameters obtained from measurements in stronger winds. The measurement dates are as follows:





 $(m/s/s)**2$ 

rJ)

3.752e-03

RMS-dsp 1.620e-03

5.441e-02









 $(s/s)^{**2}$ 

B-4

6.473e-03

8.571e-03

6.615e-02



RMS-ve1 1.721e-02 RMS-dsp 1.272e-02

8.417e-02





 $5.674e-02$ 

 $(m/s/s) * *2$ 



RMS-acc RMS-vel RMS-dsp<br>4.156e-02 2.618e-03 1.233e-03 4.156e-02 2.618e-03 1.233e-03





 $\hat{\boldsymbol{\beta}}$ 



3.5207  $5.6267e-06$  $5.6326e-06$ 5 4.4550 1.1168e-05 1.1206e-05 177  $RMS$ -acc  $RMS-vel$  $RMS-dsp$  $4.503e-02$  $7.520e-03$  $5.459e-03$ 

 $0.3$ 

 $0.2$ 

4.4556

























 $R - 12$ 







 $(m/s/s)**2$ 

degrees







## Peaks In G10.51n













Peaks In G51.10e





#### Peaks In G51s.51ne



 $\bar{\mathbb{L}}$  $\mathbb T$  $\bar{\rm T}$  $\bar{1}$  $\bar{1}$  $\overline{1}$  $\mathbb{R}^n$  $\bar{V}$  $\label{eq:2} \frac{1}{2} \int_{\mathbb{R}^3} \frac{1}{\sqrt{2}} \, \mathrm{d} \mu \, \mathrm$  $\bar{+}$  $\bar{1}$  $\bar{\mathcal{A}}$  $\bar{\rm T}$  $\bar{+}$  $\parallel$  $\mathcal{L}^{\text{max}}_{\text{max}}$  ,  $\mathcal{L}^{\text{max}}_{\text{max}}$  $\mathcal{F}$  $\bar{\rm E}$  $\bar{V}$  $\bar{\mathbb{F}}$  $\mathcal{F}$  $\bar{1}$  $\bar{1}$  $\bar{\mathcal{A}}$  $\ensuremath{\mathop{\rule{0pt}{0.5ex}\hbox{}}\mathop{\rule{0pt}{0.5ex}\hbox{}}}}$  $\bar{\bar{1}}$  $\overline{1}$
#### **APPENDIX** C

# **POWER SPECTRA, PHASE SPECTRA, ESTIMATED FREQUENCIES, AND ESTIMATED DAMPING RATIOS: SECOND SET OF MEASUREMENTS**

A second set of measurements from Buildings A and B were collected at later dates, when the buildings were at a different stage of construction. Although parameters estimated from these measurements are not presented in Section 6, the spectra are included in Appendix C for completeness and for reference in future reports. The measurement dates are as follows:

Building A: Building B: August 15, 1991 August 15, 1991











C-5



C-6



 $C-7$ 











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 $\overline{1}$  $\pm$  $\bar{\Gamma}$  $\mathbf{L}$  $\bar{V}$  $\bar{1}$  $\bar{\rm E}$  $\bar{1}$  $\mathcal{F}$  $\label{eq:2.1} \frac{1}{\sqrt{2}}\int_{0}^{\infty}\frac{1}{\sqrt{2\pi}}\left(\frac{1}{\sqrt{2\pi}}\right)^{2\alpha} \frac{1}{\sqrt{2\pi}}\int_{0}^{\infty}\frac{1}{\sqrt{2\pi}}\left(\frac{1}{\sqrt{2\pi}}\right)^{\alpha} \frac{1}{\sqrt{2\pi}}\frac{1}{\sqrt{2\pi}}\int_{0}^{\infty}\frac{1}{\sqrt{2\pi}}\frac{1}{\sqrt{2\pi}}\frac{1}{\sqrt{2\pi}}\frac{1}{\sqrt{2\pi}}\frac{1}{\sqrt{2\pi}}\frac{1}{\sqrt{2\pi}}$  $\mathcal{L}^{\text{max}}_{\text{max}}$  $\bar{1}$  $\bar{1}$  $\label{eq:2.1} \mathcal{L}(\mathcal{L}^{\text{max}}_{\mathcal{L}}(\mathcal{L}^{\text{max}}_{\mathcal{L}}),\mathcal{L}^{\text{max}}_{\mathcal{L}^{\text{max}}_{\mathcal{L}}(\mathcal{L}^{\text{max}}_{\mathcal{L}^{\text{max}}_{\mathcal{L}^{\text{max}}_{\mathcal{L}^{\text{max}}_{\mathcal{L}^{\text{max}}_{\mathcal{L}^{\text{max}}_{\mathcal{L}^{\text{max}}_{\mathcal{L}^{\text{max}}_{\mathcal{L}^{\text{max}}_{\mathcal{L}^{\text{max}}_{\mathcal{$ 

 $\bar{1}$ 

 $\bar{\rm I}$ 

# **APPENDIX D SOURCE CODE LISTING FOR COMPUTER PROGRAM AMB**

The computer program AMB make extensive use of the numerical analysis libraries from Numerical Recipes in C: The Art of Scientific Computing [Press et. al 1988]. These routines are copyright (C) 1987,1988,1992 Numerical Recipes Software and are reproduced by permission, from the book Numerical Recipes: The Art of Scientific Computing, published by Cambridge University Press. Comments in the code indicate page numbers from the book where the routines came from. Some routines from this book were taken verbatim, whereas others required some modification.

#### **D.I Input**

Program AMB is interactive with the user, but requires no special graphics interface. The compiled version of the program has no command line arguments. Control information, such as sample rate and desired frequency resolution are supplied interactively.

#### **D.l.I Example of Terminal** Session

```
reeltime 15% amb
 Reference data file name: duane.52e.n
 Response data file name: duane.52c.n
 Number of data points: 60000
 Sample rate (Hz): 50
Maximum frequency in the spectrum (Hz): 25.00
 Target frequency interval (Hz): 0.01
 Actual frequency interval (Hz): 0.012207
 Points per spectrum: 2048
 Number of averages: 28
 Amplitude / Phase file name: G52cen
reeltime 16%
```
The input files must be in ASCII format. The entries in the input files must be separated by tabs, blank space, or carriage returns (new lines). The time axis of the data must be omitted. The data files require no special headers. The time histories from the measurement location sensor and the reference location sensor must be in separate files. (The UNIX command "more" displays a file page by page.)

#### D.1.2 Example of **Input** Files

```
reeltime 16% more duane.52e.n
-2.889475e-02
-1.002818e-02
1. 2067 81e-02
```
.., following data deleted ...

reeltime 17% more duane.52c.n 4.691148e-02 4.844121e-02 2.226596e-02

... following data deleted ...

### D.2 **Output**

Program AMB creates a single output file of three-column, tab delimited, ASCII data. The first column contains the frequency axis data. The second column contains the power spectrum data as computed by equation (5.6). The third column contains the phase spectrum data as computed by equation (5.9).

#### D.2.1 Example of **Output** File



... following data deleted ...

```
/* Program amb.c - Ambient Vibration Spectral Analysis
                                                                        \star /
/* calculates the power spectra Gx, Gy, and Gxy using data sets x & y
                                                                        ^{\star}/
/* Averaging and Windowing are Incorporated.
                                                                        \star /
/* W.H. Press et al. Numerical Recipes In C. Cambridge Press - Ch. 12
                                                                        \star /
/*
                                                                        \star /
/* to compile: cc -0 -o amb amb.c fft.c nrutil.c
                                                                        ^{\star}/
/* (c) H.P. Gavin, Dept. of Civil Eng., Princeton University, 2-91
                                                                        \star /
#define MXLN
               85
#include <math.h>
#include <stdio.h>
main()\left\{ \right.char
                x_filename[MXLN],
                y_filename[MXLN],
                G_filename[MXLN],
                \mathbf{C}:
        FILE
                *fpx, *fpy,
                                /* file pointers to x & y data files
                                                                        \star /
                * fpG ;
                                /* file pointer to the spectra file
                                                                        ^{\star}/
        float
                *gx, *gy,
                                /* power spectra of x & y data files
                                                                        \star/
                *gxyr, *gxyi,
                                /* real & imaginary parts of cross spct */
                *Coh,
                                /* coherence spectrum
                                                                        \star /
                *Pha,
                                                                        \star /
                                /* phase
                                          spectrum
                                /* sample rate and frequency interval
                sr, df,\starw2,
                                /* circular frequency in rad/sec squared*/
                data,
                *vector();
        int
                m = 2,
                               /* number of frequency values / segment */
                tg_pts,
                                /* target number of freq vals / segment */
                                /* number of segments to be averaged
                k,
                                                                        */
                                /* counter
                                                                        \star /
                j,
                \text{ovrlap} = 1,
                                \frac{1}{2} 1: overlap segments; 0: no overlap
                                                                        \star/points = 0;
                                /* number of data points (10,000)
                                                                        \star /
        void
                spectra(),
                               /* returns spectra and cross spectra
                                                                        \star /
                \coh(),
                                /* returns coherence and phase
                                                                        \star/
                free\_vector();
        printf (" Reference data file name: ");
        scant ("ss", x_filename);
        if (( fpx=fopen(x_filename, "r")>=NULL) {
                printf (" error: cannot open %s\n", x_filename);
                exit (0);
        \mathcal{I}printf (" Response data file name: ");
        scanf ("ss",y_ffilename);
```

```
if (( fpy=fopen(y_filename,'r") ) ==NULL) {
        printf (" error: cannot open %s\n", y_filename);
        exit (0);
\mathcal{Y}\frac{1}{2} Size The Data \frac{1}{2} \frac{1}{2}printf (" Number of data points: ");
scanf ("%d", &points);
if (points == 0) {
        while ( (c=fscanf(fpx, "f", \&data)) != EOF ) ++points;
        printf (" %d X data points\n",points);
        points = 0;
        while ( (c=fscanf(fpy, "f", &data)  != EOF ) ++points;
        printf (" %d Y data points\n",points);
        rewind (fpx) j
        rewind (fpy) j
\mathcal{F}/* Choose Frequency Resolution */
printf (" Sample rate (Hz): ");
scanf ("%f", &sr);
printf (" Maximum frequency in the spectrum (Hz): %.2f\n", sr/2)j
printf (" Target frequency interval (Hz): ");
scanf ("*f", \&df);
tg_pts = sr/df;while (m < tg_{pts}) m <<= 1; m >>= 2;
k = \text{points/m} - 1;
df = sr / (2<sup>*</sup>m);
printf (" Actual frequency interval (Hz): f\f\n", df);
printf (" Points per spectrum: %d\n", m);
.<br>printf (" Number of averages: %d\n", k);
if (k < 1) {
         printf (" error: %d is too small\n", k);
        exit(0);\mathcal{F}gx = vector(1,m);gy = vector(1,m);gxyr = vector(1,m);gxyi = vector(1,m);spectra (fpx, fpy, gx, gy, gxyr, gxyi, m, k, ovrlap);
fclose (fpx) j
fclose (fpy) j
Coh = vector(1,m);Pha = vector(1,m);
coh (gx,gy,gxyr,gxyi,Coh,Pha,m);
```

```
0-4
```

```
/* Print the Spectra */
       printf (" Amplitude / Phase file name: ");
       scanf ("%s", G_filename);
       if (( fpG=fopen (G_filename, "W") ) ==NULL) {
               printf (" error: cannot open %s\n",G_filename);
               exit (1);
       }
       for (j=2;j<=m;j++)fprintf (fpg, "Re\t%\\t@7.2f\n', (j-1)*df, gy[j], Pha[j]);
       fclose (fpG) i
       free_vector(Coh, 1, m);
       free\_vector(Pha,1,m);free\_vector(gx,1,m);free\_vector(gy,1,m);free\_vector(gxyr,1,m):
       free_vector(gxyi,1,m);
}
       /* Function spectra() */
       /* Power Spectra Using The Fast Fourier Transform */
       /* Adapted from Numerical Recipes in C - Ch 12 */
       /* normalization: SUM(gx[i]) i=1..m = variance(x) */
static float sqrarg;
#define SQR(a) (sqrarg=(a), sqrarg*sqrarg)
#define tpi 6.28318530717959
#define WINDOW(j,c,b) (0.5-0.5*cos(tpi*((j)-1)*(c)))
/* #define WINDOW(j,a,b) (1.0-fabs((((j)-1)-(a))*(b)))
/* Parzen
/* #define WINDOW(j,a,b) (1.0-SQR(((j)-1)-(a))*(b)))/* #define WINDOW(j,a,b) 1.0
void spectra(fpx,fpy,gx,gy,gxyr,gxyi,m,k,ovrlap)
FILE *fpx, *fpy;
float gX[],gy[],gxyr[],gxyi[]j
int m, k, ovrlap;
{
                                                      /* Hanning
                                                      /* Welch
                                                       /* Square
                                                                      */
                                                                       */
                                                                       */
                                                                       */
        int n, tn, tn3, tn4, kk, j,tj;float a, b, c, sumw=0.0, den=0.0;
        float *x1,*x2,*y1,*y2i
        float *SX,*SYi
        float *vector() i
        void segments(),twofft(),free_vector();
        tn4 = (tn3 = (tn = n+(n=m+m))+3)+1;a = m - 0.5;
        b = 1.0/(m+0.5);
        c = 1.0/(n-1.0);/* Useful Quantities
                                               /* Operating vectors
                                               /* FFT'd vectors
                                                                       */
                                                                       */
                                                                       */
```

```
0-5
```

```
x1 = vector(1,m):
x2 = vector(1, n);y1 = vector(1, m);
y2 = vector(1, n);Sx = vector(1, tn);Sy = vector(1, tn);for (j=1; j<=n; j++) sumw += SQR(WINDOW(j,c,b));
for (j=1; j<=m; j++) gx[j]=gy[j]=gxyr[j]=gxyi[j]=0.0;if (ovrlap) /* Initialize the "save" half-buffer'. */
        for (j=1; j<=m; j++) (
                 fscar(fpx, "if", \&x1[j]);fscanf(fpy, "%f",&yl[j]);
        }
for (kk=l; kk<=k; kk++) {
        segments(fpx, x1, x2, a, b, c, m, ovrlap);
        segments(fpy,y1,y2,a,b,c,m,ovrlap);
                         /* Fourier transform the windowed data */
        twofft(x2,y2,Sx,Sy,n,1);
                 \prime* Sum the results into the previous segments */
        gx[1] += (SQR(Sx[1]) + SQR(Sx[2]));
        gy[1] += (SQR(Sy[1]) + SQR(Sy[2]));for (j=2; j<=m; j++) (
                tj = j+j;gx[j] += (SQR(Sx[tj-1]) + SQR(Sx[tj])+ \text{SQR}(\text{Sx}[t n 3-t j]) + \text{SQR}(\text{Sx}[t n 4-t j]));
                 gy[j] += (SQR(Sy[tj]) + SQR(Sy[tj-1])+ \text{SQR(Sy[tn3-tj]) + \text{SQR(Sy[tn4-tj])});gxyr[j] += (Sx[tj-1]*Sy[tj-1] + Sx[tj]*Sy[tj]+ Sx[tn3-tj]*Sy[tn3-tj]
                                  + Sx[tn4-tj] *Sy[tn4-tj]);
                 gxyi[j] += (Sx[tj]*Sy[tj-1] - Sx[tj-1]*Sy[tj]- Sx[tn4-tj] *Sy[tn3-tj]
                                  +, Sx[tn3-tj]*Sy[tn4-tj];}
         den += sumw:}
den * = n;
                                  /* Correct normalization. */for (j=1; j<=m; j++) {
        gx[j] /= den;
         gy[j] /= den;
         gxyr[j] /= den;
         gxyi[j] /= den;
\mathcal{F}free\_vector(x1,1,m);free\_vector(x2,1,n);free\_vector(y1,1,m);free\_vector(y2,1,n);
free_vector(Sx,l,tn) ;
free\_vector(Sy, 1, tn);
return;
```

```
0-6
```
 $\mathcal{E}$ 

```
/*
                                                                   \star/
                         Function segments()
        /* Gets two
complete segments into the work space
                                                                   \star/
        /*
                        Numerical Recipes in C p446
                                                                   */
void
segments(fp,zl,z2,a,b,c,m,ovrlap)
FILE
        *f_p:
float
        *zl,*z2;
float
        a,b,ci
int
        m, ovrlap;
\left\{ \right.int j;
        if (ovrlap) {
                 for (j=1; j<=m; j++) z2[j]=z1[j];for (j=1; j<=m; j++) fscanf(fp, "f", z1[j]);
                 for (j=1; j<=m; j++) z2[m+j]=z1[j];} else {
                 for (j=1; j<=m+m; j++)f scanf (f p, "f", \& z2[j]) ;
        }
        for (j=1; j<=m+m; j++) /* Apply the window to the data.*/
                 z2[j] *= WINDOW(j, c, b);
        returni
}
        \frac{1}{\sqrt{2}} Function coh () \frac{1}{\sqrt{2}} /
        /* Computes the coherence and the phase functions using */
        \prime^* the power spectra gx and gy and the cross spectrum gxy*/
void coh(gx,gy,gxyr,gxyi,Coh,Pha,m)
float *gx,*gy,*gxyr,*gxyi,
        *Coh,*Phai
int m;
\{int
                 j ;
         float
denl, den2;
         for (j=1; j<=m; j++) {
                 denl = SQR(gxyr[j]) + SQR(gxyi[j]);if (den2 = (gx[j]*gy[j]) != 0.0)Coh[j] = den1/den2;else Coh[j] = O.Oi
                 if (gxyi[j] := 0.0)Pha[j] = fabs(atan2(gxyi[j],gxyr[j])) * 57.29578;
                 else Pha[j] = Pha[j-1];
         \mathcal{F}return;
\mathcal{E}
```

```
D-7
```

```
\sqrt{*} \sqrt{*} \sqrt{*} \sqrt{*}/* Fast Fourier Transform routines from Numerical Recipes in C, by Press*/
/* et al. Cambridge University Press, 1988 */
#include <math.h>
        \frac{1}{\sqrt{2}} Function four 1() \frac{1}{\sqrt{2}} /
       /* Fast Fourier Transform - Numerical Recipes in C p411.*/
#define SWAP(a,b) tempr=(a);(a)=(b);(b)=tempr
void four1(data,nn,isign)
float data[];
int nn, isign;
€
        int
               n,mmax,m,j,istep,i;
        double wtemp,wr,wpr,wpi,wi,theta;
       float
               tempr, tempi i
       n=nn « 1i
        j = 1;/* Bit Reversal */
        for (i=1; i < n; i += 2) {
               if (j > i) {
                       SWAP(data[j], data[i]);
                       SWAP(data[j+1], data[i+1]);
               \mathcal{F}m=n \gg 1;
               while (m \geq 2 \&\& j > m) (
                       j -= mi
                       m »= 1i
                }
               j += mi
        \mathcal{Y}mmax=2;while (n > mmax) {
               istep = 2*mmax;theta = 6.28318530717959 / (isign*mmax);
               wtemp = sin(0.5*theta);
               wpr = -2.0*wtemp*wtemp;wpi = sin(theta);
               wr = 1.0i
               wi = O.Oi
                for (m=1; m<mmax; m+=2) {
                        for (i=mi i<=ni i+=istep) {
                               j = i + mmax; / Danielson-Lanczo
                               tempr = wr*data[j] - wi*data[j+1];tempi = wr*data[j+1] + wi*data[j];data[j] = data[i] - tempr;data[j+1] = data[i+1] - tempi;data[i] += tempr;data[i+1] += tempi;\mathcal{Y}/* Trigonometric Recurrence
```

```
wr = (wtemp=wr)*wpr - wi*wpi + wri
                            wi = wi*wpr + wtemp*wpi + wi;\mathcal{F}mmax = istep;\mathcal{Y}returni
\mathcal{E}/* function twofft() *<br>/* The simultaneous FFT of two real valued functions */
         /* The simultaneous FFT of two real valued functions */<br>/* from Numerical Recipes In C p 415 */
                   from Numerical Recipes In C p 415
void twofft(data1,data2,fft1,fft2,n)
float data1[],data2[],fft1[],fft2[];
int ni
\left(int nn3,nn2,jj,j;float rep, rem, aip, aim;
         void four1() i
         nn3=1+(nn2=2+n+n)i
         for (j=1,jj=2;j<=n;j++,jj+=2) (
                   fftl[jj-1]=datalf[j];f[t1[j] = data2[j];}
         four1(fft1, n, 1);fft2[1]=fft1[2];
          fft1[2]=fft2[2]=0.0;for (j=3ij<=n+1ij+=2) {
                   rep=0.5*(fft1[j]+fft1(nn2-j));rem=O.5*(fft1[j]-fft1[nn2-j]);
                   aip=0.5*(fft1[j+1]+fft1[nn3-j]);aim=0.5*(fft1[j+1]-fft1[nn3-j]);
                   fftl[j]=repi
                   fft1[j+1]=aim;fft1[nn2-j]=rep;fft1[m3-j] = -aim;fft2[j]=aip;fft2[j+1] = -remi
                   ftt2[nn2-j]=aip;fft2[nn3-j]=rem;}
          return;
}
          /* function realft() *<br>/* The FFT of 2n real valued discrete function pointS */
          /* The FFT of 2n real valued discrete function pointS */<br>/*       from  Numerical Recipes In C  p 417             */
          /* from Numerical Recipes In C p 417
void realft(data,n,isign)
```

```
D-9
```

```
float data[];
int n,isign;
        int i,i1,i2,i3,i4,n2p3;
        float cl=0.5, c2, h1r, h1i, h2r, h2i;double wr,wi,wpr,wpi,wtemp,theta;
        void four1();
        theta=3.141592653589793/(doub1e) n;
        if (isign == 1) {
                c2 = -0.5;four1(data,n,l);
        } else (
                c2=0.5;theta = -theta;
        }
        wtemp=sin(O.5*theta);
        wpr = -2.0*wtemp*wtemp;wpi=sin (theta) ;
        wr=1.0+wpr;wi=wpi;
        n2p3=2*n+3;for (i=2; i<=n/2; i++) {
                 i4=1+(i3=n2p3-(i2=1+(i1=i+i-1)));
                h1r=cl*(data[i1]+data[i3]) ;
                h1i = c1*(data[i2] - data[i4]);h2r = -c2*(data[i2]+data[i4]);h2i=c2*(data[i1]-data[i3]);data[i1]=h1r+wr*h2r-wi*h2i;
                 data[i2]=h1i+wr*h2i+wi*h2r;
                 data[i3]=h1r-wr*h2r+wi*h2i;
                 data[i4] = -h1i+wr*h2i+wi*h2r;wr=(wtemp=wr)*wpr-wi*wpi+wr;
                 wi=wi*wpr+wtemp*wpi+wi;
        }
        if (isign == 1) {
                 data[1] = (h1r=data[1]) + data[2];data[2] = h1r - data[2];} else {
                 data[1]=c1*(-th1r=data[1])+data[2]);data[2] = c1*(h1r - data[2]) ;
                 four1(data, n, -1);
         \mathcal{E}
```
{

 $\mathcal{E}$ 

```
0-10
```

```
/* FILE nrutil.c<br>/* Memory allocation functions from Numerical Recipes in C, by Press, */
/* Memory allocation functions from Numerical Recipes in C, by Press, */
/* Cambridge University Press, 1988
\frac{1}{2} #include <malloc.h> */
#include <stdio.h>
typedef struct FCOMPLEX {float r,i;} fcomplex; /* */
void nrerror(error_text)
/* print error message to stderr */
char error_text[];
(
        void exit();
        fprintf(stderr,"Numerical Recipes run-time error... \n");
        fprintf (stderr, \sqrt{8} \n\pi, error_text);
        fprintf(stderr,"... now exiting to system...\n\cdot \n\cdot);
        exit(1);\mathcal{Y}float *vector(nl,nh)
/* allocate storage for a vector */
int nl,nh;
{
        float *v;
        v=(float * ){malloc}({(unsigned) (nh-nl+1) * sizeof(float)});if (!v) nrerror("allocation failure in vector()");
        return v-nl;
\mathbf{1}fcomplex *Cvector(nl,nh)
int nl,nh; /* allocate storage for a complex vector
                                                                             */
- (
         fcomplex *v;
         v=(fcomplex *) \text{malloc}((unsigned) (nh-nl+1)*sizeof(fcomplex));if (!v) nrerror("allocation failure in Cvector()");
         return v-nl;
}
double *dvector(nl,nh)
1* allocate storage for a vector */
int nl,nh;
\left\{ \cdot \right\}double *v;
         v=(double *) \text{malloc}((unsigned) (nh-nl+1)*sizeof(double));if (!v) nrerror("allocation failure in dvector()");
         return v-nl;
\mathcal{F}*/int *ivector(nl,nh) /* allocate storage for a vector
```

```
int nl,nh;
€
        int *v;
        v=(int *)malloc((unsigned) (nh-nl+1)*sizeof(int));
        if (!v) nrerror("allocation failure in ivector()");
        return v-nl;
\mathcal{E}float **matrix(nrl,nrh,ncl,nch) /* allocate storage for a matrix
                                                                              */
int nrl,nrh,ncl,nch;
\left\{ \right.int i;
        float **m;
        m= (float **) malloc ((unsigned) (nrh-nrl+l) *sizeof (float*));
        if (!m) nrerror("allocation failure 1 in matrix()");
        m = nr1;for (i=nr1; i<=nrh; i++) {
                 m[i]=(float *) malloc((unsigned) (nch-ncl+l)*sizeof(float))
                 if (\{\text{Im}[i]\}) nrerror("allocation failure 2 in matrix()");
                 m[i] -= nc1;
        \mathbf{a}return m;
\mathcal{F}fcomplex **Cmatrix(nrl,nrh,ncl,nch)
int nrl,nrh,ncl,nch; /* allocate storage for a Complex matrix
                                                                              */
\left\{ \right.int i;
         fcomplex **m;
        m=(fcomplex **)malloc((unsigned) (nrh-nrl+l)*sizeof(fcomplex*));
         if (!m) nrerror ("allocation failure 1 in Cmatrix()");
        m = nr1;for (i=nr1;i<=nrh;i++) {
                 m[i] = (fcomplex *)malloc((unsigned)(nch-ncl+1)*sizeof(fcompl
                 if (\lfloor m[i] \rfloor) nrerror("allocation failure 2 in Cmatrix()");
                 m[i] -= ncl;\mathcal{E}return m;
\rightarrowdouble **dmatrix(nrl,nrh,ncl,nch) /* allocate storage for a matrix */
int nrl,nrh,ncl,nch;
\{int i;
         double **m;
         m=(double **) malloc((unsigned) (nrh-nrl+l)*sizeof(double*));
         if (!m) nrerror("allocation failure 1 in dmatrix()");
         m = nr1;
```

```
for (i=nr1; i<=nrh; i++) {
                 m[i]=(double *) malloc((unsigned) (nch-ncl+l)*sizeof(double
                 if (\{m[i]\}) nrerror("allocation failure 2 in dmatrix()");
                 m[i] -= ncl;
        \mathcal{F}return m;
\mathbf{I}int **imatrix(nrl,nrh,ncl,nch) /* allocate storage for a matrix
                                                                               */
int nrl,nrh,ncl,nch;
\left\{ \right.int i;
         int **m;
        m=(int **) malloc((unsigned) (nrh-nrl+1)*sizeof(int*));
         if (|m\rangle nrerror("allocation failure 1 in imatrix()");
        m = nr1;for (i=nr1; i<=nrh; i++) (
                 m[i] = (int * ) malloc((unsigned) (nch-ncl+1)*sizeof(int));
                  if (\text{Im}[i]) nrerror("allocation failure 2 in imatrix()");
                 m[i] -= ncl;\mathcal{Y}return m;
\mathcal{Y}float **submatrix(a,oldrl,oldrh,oldcl,oldch,newrl,newcl)
float **a;
int oldrl,oldrh,oldcl,oldch,newrl,newcl;
€
         int i,j;
         float **m;
         m=(fload **) malloc((unsigned) (oldrh-oldrl+1)*sizeof(float*));
         if (|m\rangle nrerror("allocation failure in submatrix()");
         m -= new1;
         for(i=oldrl,j=newrl;i<=oldrh;i++,j++) m[j]=a[i]+oldcl-newcl;
         return m;
\mathcal{L}void free_vector(v,nl,nh)
float *v;
int nl,nh;
{
         free((char*) (v+n1));}
void free_Cvector(v,nl,nh)
fcomplex *v;
int nl,nh;
```

```
free ( (char<sup>*</sup>) (v+n1));}
                                                                                                              e Artista (m. 1938)<br>1903: Johann Barnett, filosof filosof (m. 1904)<br>1903: Johann Barnett, filosof filosof (m. 1904)
void free_dvector(v,nl,nh)
                                                                                                                                                       \mathcal{L}_{\rm{max}} = \mathcal{L}^{\rm{max}}double *v;
int nl,nh;
{
                    fl;'ee ( (char*) (v+nl» j
                                                                                                                                                      Special Control of Super
}
void free ivector(v,nl,nh)
int *v;
int nl,nh;
                    free ( (char*) (v+nl) ) ;<br>free ( (char*) (v+nl) ) ;
\left\{ \right.}
void free_matrix(m,nrl,nrh,ncl,nch)
float **m; \ldots \ldotsint nrl,nrh,ncl,nch;
 \left\{ \right.\mathcal{L}_{\text{max}} , and \mathcal{L}_{\text{max}}int i;
                     for(i=nrh;i>=nrl;i--) free((char*) (m[i]+ncl));
                                                                                                                                                  \label{eq:2.1} \mathcal{L}_{\mathcal{A}}(\mathcal{A})=\mathcal{L}_{\mathcal{A}}(\mathcal{A})\mathcal{L}_{\mathcal{A}}(\mathcal{A})=\mathcal{L}_{\mathcal{A}}(\mathcal{A})\mathcal{L}_{\mathcal{A}}(\mathcal{A})\mathcal{A}free((char*) (m+nr1));
                                                                                                                                                              The State State
 }
                                                                                     \label{eq:2.1} \mathcal{L}_{\mathcal{A}}(\mathcal{A})=\mathcal{L}_{\mathcal{A}}(\mathcal{A})\otimes\mathcal{L}_{\mathcal{A}}(\mathcal{A})\otimes\mathcal{L}_{\mathcal{A}}(\mathcal{A})\otimes\mathcal{L}_{\mathcal{A}}(\mathcal{A})\otimes\mathcal{L}_{\mathcal{A}}(\mathcal{A})\otimes\mathcal{L}_{\mathcal{A}}(\mathcal{A})\otimes\mathcal{L}_{\mathcal{A}}(\mathcal{A})\otimes\mathcal{L}_{\mathcal{A}}(\mathcal{A})\otimes\mathcal{L}_{\mathcal{A}}(\mathcal{A})\otimes\mathcal{void free_Cmatrix(m,nrl,nrh,ncl,nch)
                                                                                                                                                         \label{eq:2} \frac{1}{\sqrt{2\pi}}\int_{\mathbb{R}^3}\frac{1}{\sqrt{2\pi}}\left(\frac{1}{\sqrt{2\pi}}\right)^2\frac{dx}{dx}dx.fcomplex **m;
 int nrl,nrh,ncl,nch;
 \left\{ \right.\text{int}_{\mathbb{R}^d} \left\{ \frac{1}{2} \sum_{i=1}^d \left\{ \frac{1}{2} \right\} \right\} = \left\{ \frac{1}{2} \sum_{i=1}^d \left\{ \frac{1}{2} \sum_{i=1}^d \left\{ \frac{1}{2} \right\} \right\} \right\} = \left\{ \frac{1}{2} \sum_{i=1}^d \left\{ \frac{1}{2} \sum_{i=1}^d \left\{ \frac{1}{2} \right\} \right\} \right\} = \left\{ \frac{1}{2} \sum_{i=1}^d \left\{ \frac{1}{2} \right\} \right\for(i=nrh;i>=nrl;i--) free((char*) (m[i]+ncl));
                      free((char*) (m+nrl));\mathcal{P}_{\text{int}}(x,y) = \int_{\mathcal{P}_{\text{int}}(x,y)} \mathcal{P}_{\text{int}}(x,y) \mathcal{P}_{\text{int}}(x,y) \mathcal{P}_{\text{int}}(x,y) \mathcal{P}_{\text{int}}(x,y) \mathcal{P}_{\text{int}}(x,y) \mathcal{P}_{\text{int}}(x,y)\mathcal{F}stage of the stage of the stage
                                                                                                                                                                1.1211void free_dmatrix(m,nrl,nrh,ncl,nch)
 double **m;
 int nrl, nrh, ncl, nch;
 {
                                                                                                                             \mathcal{L}_{\text{max}} and \mathcal{L}_{\text{max}} and \mathcal{L}_{\text{max}} are the set of the set of the set of the \mathcal{L}_{\text{max}}int i;
                                                                                                                                                                                  Contract
                                                                                                                                                                     \mathcal{L}(\mathcal{L}^{\mathcal{L}}) .
                      for(i=nrh;i=nr1;i--) free((char*) (m[i]+ncl));free((char*) (m+nr1));Contract Bank Contract Dealer
 }
 void free_imatrix(m,nrl,nrh,ncl,nch)
                                                                                                                              the group of the state and considerable
 int **mj
                                                                                                                                                                          \label{eq:1} \mathcal{L}(\mathcal{F}) = \mathcal{L}(\mathcal{F}) \mathcal{L}(\mathcal{F})int nrl, nrh, ncl, nch;
                                                                                                                                                                         \label{eq:2.1} \mathcal{A}(\mathbf{x}) = \frac{1}{2} \mathcal{A}(\mathbf{x}) \mathcal{A}(\mathbf{x})\left\{ \right.
```

```
int i;
        for(i=nrh;i>=nr1;i--) free({{char}}^\star) (m[i]+ncl)) ;
         free((char*) (m+nr1));\mathcal{Y}void free_submatrix(b,nrl,nrh,ncl,nch)
float **b;
int nrl,nrh,ncl,nch;
\left\{ \right.free((char*) (b+nr1));\mathcal{I}float **convert_matrix(a,nrl,nrh,ncl,n~h)
float *a;
int nrl,nrh,ncl,nch;
\left\{ \right.int i,j,nrow,ncol;
         float **m;
         nrow=nrh-nrl+l;
         ncol=nch-ncl+l;
         m = (float **) malloc((unsigned) (nrow)*sizeof(float*));if (!m) nrerror("allocation failure in convert_matrix()");
         m = nr1;
         for(i=0,j=nrl;i<=nrow-1;i++,j++) m[j]=a+ncol*i-ncl;return m;
}
void free_convert_matrix(b,nrl,nrh,ncl,nch)
float **b;
int nrl,nrh,ncl,nch;
{
         free((char*) (b+nr1));\mathcal{F}
```
#### APPENDIX E

# SOURCE CODE LISTING FOR COMPUTER PROGRAM PEAK

Some routines used by program PEAK are copyright (C) 1987,1988,1992 Numerical Recipes Software and are reproduced by permission, from the book Numerical Recipes: The Art of Scientific Computing, published by Cambridge University Press.

## E.l Input

Input to program PEAK is interactive. The program is invoked with a command line argument which is the name of the file created by program AMB. The rest of the input, such as the frequency band-width for peak searching, the number of peaks searched for, and the name of the output file, is entered interactively by the user.

## E.2 Output

The output of the file is in the form of a table. Examples of these tables are given in

Appendices B and C.

## E.3 Example of Terminal Session

```
reeltime 8% peak G52cen
 Number of Peaks: 4
 Minimum frequency for peak-picking: 0.1
 Maximum frequency for peak-picking: 2
 Closest peak spacing (Hz): 0.25
 Results file name: r52cen
reeltime 9%
reeltime 9% more r52cen
                         Peaks In G52cen
Hz-data Hz-fit Amp^2-data Amp^2-fit Phase-Damping %
 z-data Hz-fit Amp^2-data Amp^2-fit PhaseDampi<br>0.2441   0.2426   8.4451e-05   8.5265e-05   1   2.8<br>0.9888   0.9891   4.7639e-06   4.7664e-06   2   1.3
 0.9888  0.9891  4.7639e-06  4.7664e-06  2  1.3<br>1.5259  1.5243  1.4874e-07  1.4911e-07  158  1.8
 1.5259   1.5243   1.4874e-07   1.4911e-07   158   1.8<br>1.8066   1.8102   1.3160e-05   1.3314e-05   3   2.2
                         1.3160e-05 1.3314e-05 3
            RMS-acc RMS-vel RMS-dsp
            3.057e-02 9.156e-03 6.181e-03
```
reeltime 10%

/\* Program peak.c - Spectral Peak Picking For Ambient Data  $\star$  / /\* Use data from the output of program amb.c for ambient parameter est. \*/ /\* Finds amplitudes and phases using peak picking methods.  $\star$  /  $\star$ /\* Estimates frequencies and damping using a quadratic curve-fit and /\* a band-width method combining raw data with the estimated frequency. \*/ /\* RMS-displacement computation de-emphasizes low frequencies: line 71  $\star$  /  $/$ \*  $\star$  / /\* to compile: cc -0 -o peak peak.c gaussj.c nrutil.c  $\star$  / /\* to run: peak 'amplitude / phase filename's and the property  $1 - 10 = 10$  $\star$ /\* (c) H.P. Gavin, Dept. of Civil Eng., Princeton University, 2-91  $\star$  / #include <stdio.h> and the control of the cont #include <math.h>  $\mathcal{L}_{\text{max}}$  , we prove a split in the constraint of the split  $\mathcal{L}_{\text{max}}$ main(argc, argv) int  $\arg c$ ; char  $*$ argv $[]$ ;  $\left\{ \right.$ المتعاطي والمرا filename [85], /\* output file name char المريدان والموركين  $c$ ; \*fp;  $\longrightarrow$  /\* file pointer  $\longrightarrow$   $\longrightarrow$   $\longrightarrow$ FILE 1 frq, amp, pha, /\* frequency, amplitude, and phase float  $\star$  /  $frq$  old, /\* previous frequency  $*$  / /\* previous amplitude and the series amp\_old,  $\star$  /  $\gamma^*$  previous phase  $\star$  / pha\_old, /\* closest spacing of 2 peaks<br>/\* minimum frequency for peak picking  $\star$ /  $f\_tol$ ,  $\star$  / f\_min, /\* maximum frequency for peak picking  $^{\star}/$ f\_max,  $\star$  / df, /\* frequency interval a sa tanàna amin'ny faritr'i Nord-Aquitaine, ao Frantsa.<br>Ny INSEE dia mampiasa ny kaodim-paositra 2008–2014. Ilay kaominina dia kaominina mpikambana amin'ny fivondrona  $\sqrt{\star}$  rms acceleration  $\star$ /  $rms_a = 0.0,$  $\star$  /  $rms_v = 0.0$ , /\* rms velocity and the set  $rms_d = 0.0,$ /\* rms displacement  $\star$ /\* data in a high peak region  $\star$ 7 \*\*pk\_data,  $***C.$ /\* power polynomial coefficients  $*$ \*\*matrix(); /\* allocates matrix storage  $\star$ /  $n\_pk$ , /\* number of peaks to pick  $\star$  / int slope,  $\frac{1}{2}$  sign of [ dG(f) / df ] (+/- 1)  $\star$  /  $\frac{1}{2}$  /\* previous slope  $\star$  / slope\_old,  $\frac{1}{2}$  1: change in slope, 0: same slope slp\_cond,  $*$ /\* 1: next largest amplitude, 0: not  $*$ amp\_cond,  $\frac{1}{2}$  1: far from a peak, 0: close to pk  $\star$  / frq\_cond,  $\star$  / /\* current peak number  $i, j, k, n;$ struct peak {float f,a,p,d,F,A;} pk[20]; /\* general least squares polynomial fit \*/ void  $poly\_fit()$ ,  $\star$ / /\* deterministic curve fitting  $\texttt{dter\_fit}()$ ,

```
free_matrix(); /* deallocates matrix storage
                                                                                           */
                                                                       \mathcal{C}=\mathcal{C}_{\mathcal{A},\mathcal{A}} , \mathcal{C}_{\mathcal{A}}if (( fp = fopen(argv[1], "r")) == NULL ) {
           printf (" error: cannot open %s\n", argv[1]);
           printf (" usage: peak 'amplitude/phase filename'\n");
           exit (0);
\mathcal{Y}/* Find The RMS Values */
/* Change These Lines If The Data Is Velocity */
while ((c = fscan f(fp, "f f * f * f", \& frq, \& nq), \& pha) ) != EOF ) {
           if(frq < 0.15) frq=0.16;/* filter low frequency noise */<br>frq *= 2.0*acos(-1.0): /* w = 2.pi.f */
           frq *= 2.0*acos(-1.0); /* w = 2.pi.f<br>frq *= frq; /* w^2 = w.w */
           frq *= frq; \sqrt{2} w^2 = w.w
           rms_a += amp; \frac{1}{2} /* a^2 = a^2 or a^2 = v^2*w^2 */<br>rms_v += amp/frq; \frac{1}{2} /* v^2 = a^2/w^2 or v^2 = v^2 */
                                         r^* v<sup>2</sup> = a<sup>2</sup>/w<sup>2</sup> or v<sup>2</sup> = v<sup>2</sup> */
           rms_d += amp/(frq*frq); /* d^2 = a^2/w^4 or d^2 = v^2/w^2*/
\mathcal{Y}rms_a = sqrt(rms_a);rms_v = sqrt(rms_v);rms_d = sqrt(rms_d);rewind (fp) ;
printf (" Nwnber of Peaks: ");
scanf ("%d", &n_pk);
printf (" Minimum frequency for peak-picking: ");
scanf ("If", & f=min);printf (" Maximum frequency for peak-picking: ");
scanf ("f'', f'', f'' f'', f'''printf (" Closest peak spacing (Hz): ");
scanf ("f'', f(t_0); \ldotspk\_data = matrix(1, 20, 1, 2);C = matrix(1, 3, 1, 1);Service State
            /* Find The Peaks By Amplitude */
                       2. 法国际管理人员
i=1;\label{eq:2.1} \mathcal{L}_{\mathcal{A}}=\frac{1}{2}\left(\frac{1}{2}\left(\frac{1}{2}\right)^{2}+\frac{1}{2}\left(\frac{1}{2}\right)^{2}+\frac{1}{2}\left(\frac{1}{2}\right)^{2}+\frac{1}{2}\left(\frac{1}{2}\right)^{2}+\frac{1}{2}\left(\frac{1}{2}\right)^{2}+\frac{1}{2}\left(\frac{1}{2}\right)^{2}+\frac{1}{2}\left(\frac{1}{2}\right)^{2}+\frac{1}{2}\left(\frac{1}{2}\right)^{2}+\frac{1}{2}\left(\frac{1}{2}\right)^{2}+\frac{1}{2}\left(\frac{pk[1].a = 0.0;while ((c = fscanf(fp, "\$f \\$f \$f", &frq, &amp, &pha)) != EOF ) {
            if (frq > f_max) break;
            if (frq > f=min)if (\text{pk}[1].a < amp) {
                            pk[1].f = frq; ( ) and ( ) where \mathbb{P}^1pk[1].q = amp;pk[1] . p = pha;}
            if (i) (i)df = frq;TAR TERMINARY
                       i=0;
```

```
E-3
```

```
}
}
rewind(fp);
for (i=2; i \le n_{pk}; i++) (
        frq_old = amp_old = pk[i].a = slope = slope_old = 0;
        while ((c = fscanf(fp, "if %f *f", %frq, %amp, %pha)) := EOF )if (frq > f_max) break;
                 if (frq > f=min) (
                          slp_cond = Oi
                          if (amp> amp_old) slope = Ii
                          if (amp < amp_old) slope = -Ii
                          if (slope_old == 1 && slope == -1) slp_cond
                          amp_cond = Oi
                          if (\text{pk}[i].a < amp\_old &amp; \&amp; amp\_old < pk[i-1].amp_{cond} = 1;
                          frrq_{cond} = 1;
                          for(k=1; k< i; k++)if (fabs(frq_old - pk[k].f) < f_tolfrq_{\text{cond}} = 0;if (slp\_cond && amp_cond && frq_cond) (
                                  pk[i].f = frq_old;pk[i].a = amp\_old;pk[i].p = pha<sup>o</sup>ld;
                          }
                          frqold = frq;
                          amp\_old = amp;pha<sup>old = pha;</sup>
                          slope\_old = slope;}
         }
         rewind(fp);
\mathcal{E}/* Sort The Peaks By Frequency */
for (i=2; i<=n\_pk; i++) {
         frq = pk[i].f;amp = pk[i].a;pha = pk[i].p;k=i-1;while (k > 0 && pk[k].f > frq) {
                 pk[k+l].f = pk[k] .fj
                 pk[k+l].a = pk[kl .aj
                 pk[k+1] .p = pk[k] .pk--i
         }
         pk[k+1].f = frq;
```

```
pk[k+1].a = amp;pk[k+1].p = pha;
}
        /* Curve-fit The Peaks With A Quadratic */
amp\_cond=0; n=i=1;while ((c = fscan f(fp, "if 8f 8f", 6frq, 6amp, 6pha)) := EOF) (
        if (i > n_{pk} || frq > f_{max}) break;
        if (frq > pk[i].f - 1.1*df) n = amp_cond = 1;
        while (amp_cond) {
                pk\_data[n][1] = frq;pk\_data[n][2] = amp;if (frq > pk[i].f + 0.9*df) amp_cond = 0;
                fscant(fp, " if *f *f, kfq, kamp, kpha;
                n++;}
        n--;if (frq > pk[i].f & g_a n > 2) {
                 if (n == 3)dter_fit (pk_data,n,C);
                 else
                                 poly_fit (pk_data,n,2,C);
                pk[i].A = C[1][1] - C[2][1] *C[2][1] / (4.0 * C[3] [1])pk[i].F = -0.5*C[2][1]/C[3][1];i++;} else if (frq > pk[i] .f) (
                         pk[i].F = pk[i].f;pk[i].A = pk[i].a;i++;\mathcal{Y}}
        /* Estimate The Damping With Raw Data */
rewind (fp) ;
i=1:
while ((c = fscant(fp, "if %f *f", %frq, %amp, %pha)) := EOF) {
        if (i > n\_pk \mid \mid frq > f\_max) break;
        if (frq == pk[i].f) {
                 while \langle \text{amp} > 0.7* \text{pk}[i] .a \rangle {
                         fscanf(fp, *f *f *f *, *frq, *amp, *pha);
                         pk[i].d = 100.0 * fabs(frq - pk[i].F) /(pk[i].F * sqrt(pow((frq/pk[i].F), 4pk[i].A/amp - 1.0);
                 }
                 i++;}
}
fclose (fp);
```

```
的复数形式 计类型系统
                               /* Print The Results / */* in /printf (" Results file name: ");
               scanf ("\frac{1}{5}", filename);<br>if (( fp=fopen(filename, "w")) == NULL ) {
               scanf ("%s", filename);
                               printf (" error: cannot open s\ n", filename);
                 \text{exit}_s(0); \quad \dots}
                                              \frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\frac{1}{2}\left(\frac{1}{2}-\frac{1}{2}\right)\left(\frac{1}{2}-\frac{1}{2}\right)\left(\frac{1}{2}-\frac{1}{2}\right)\left(\frac{1}{2}-\frac{1}{2}\right)\left(\frac{1}{2}-\frac{1}{2}\right)\left(\frac{1}{2}-\frac{1}{2}\right)\left(\frac{1}{2}-\frac{1}{2}\right)\left(\frac{1}{2}-\frac{1}{2}\right)\left(\frac{1}{2}-\frac{1}{2}\right)\left(\frac{1}{2}-\frac{1}{2}\right)\left(\frac{1}{2}-\frac{1}{2}\right)\left(\fprintf (fp, "\t\t\t\tPeaks In s\n\n'\n', argv[1]);
               fprintf (fp, "Hz-data\t\tHz-fit\t\tAmp^2-data\tAmp^2-fit\t");
               fprintf (fp, "Phase\tDamping \ell \\n");
                                                                               (一) 2000公式
                               \begin{array}{ll} \texttt{i} <= \texttt{n\_pk}_t \texttt{i++} \\ \texttt{f} <= \texttt{f} \texttt{f} \texttt{f} + \texttt{f} \texttt{f} \texttt{f} \texttt{f} + \texttt{f} \texttt{f} \texttt{f} \texttt{f} \texttt{f} + \texttt{for (i=1; i \le n_p k; i++).
                                pk[i] .f,pk[i] .F,pk[i] .a,pk[i] .A,pk[i] .p,pk[i] .d);
                fprintf (fp, "n");
                fprintf (fp, "\t\tRMS-acc\t\tRMS-vel\t\tRMS-dsp\n");
                fprintf (fp, "\t\t%.3e\t%.3e\t%.3e\n", rms_a, rms_v, rms_d);
                fprintf (fp, "\n\n\cdot");
                                                                        \sim 10^{-11} and \sim 10^{-10} and \sim 10^{-10}\sim 10^{-1}Constitution
                free_matrix.(pk\_data,1,20,1,2);free_matrix(C,1,3,1,1);\int_{0}^{\infty} fclose (fp); \int_{\mathbb{R}^{3}}^{\infty} , \int_{\mathbb{R}^{3}} , \int\mathcal{Y}% Function dter_fit() *<br>
Function dter_fit() *<br>
/* Fit a power polynomial to data deterministically */
                /* Fit a power polynomial to data deterministically
void dter_fit (data,n,C)
float **data,
                                              /* the data to be fit
                                                                                                                                                 */
                                                 \mathsf{A}^* the power polynomial coeffs & [X]t \{y\}*/
                **Ci
 int n_i, n_i number of data points
                                                                                                                                                 */
(
                                **X, /* the power poly basis function matrix *1
                float
                                 **matrix() i
                int^{\sqrt{2}}\mathtt{i}, \mathtt{j} ; and if \mathtt{j} is a set of the set of the
                                                                                 \label{eq:2.1} \mathcal{F}(\mathcal{F}_{\mathcal{G}}) = \mathcal{F}(\mathcal{F}_{\mathcal{G}}) \mathcal{F}_{\mathcal{G}}(\mathcal{F}_{\mathcal{G}}) \mathcal{F}_{\mathcal{G}}(\mathcal{F}_{\mathcal{G}})void gaussj(), /* gauss-jordan elimination */
    , which are {\rm free\_matrix}() ; and \ldots , we have
x = \max_{x \in \mathcal{X}} (1_{\epsilon}n_{\epsilon}1,n);for (i=1; i<=n; i++) {
                                X[i][1] = 1.0;for (j=2; j<=n; j++)X[i][j] = X[i][j-1]*data[i][1];r i se la servizione di
                                 C[i][1] = \text{data}[i][2];
```


```
\label{eq:2.1} \mathcal{L}(\mathcal{A}) = \mathcal\mathcal{E}gaussj (X,n,C,l);
           free_matrix(X,1,n,1,n);
           return;
}
                                                                                          */
           /* Function poly_fit()
                                                                                          */
           /* Fit a power polynomial to a set of data using
           /* linear least squares & gauss-jordan inversion
                                                                                          */
void poly_fit (data, n, order, C)<br>float **data, /* the
float **data, /* the data to be fit<br>**C; /* the power polynomie
                                                                                                      */
                                  /* the power polynomial coeffs & [X]t \{y\}*/
int order, /* the order of the fit
                                                                                                      */
           n; \frac{1}{\sqrt{2}} /* the number of data points > order
                                                                                                      */
\left\{ \right.**X, /* the power poly basis function matrix */<br>**A. /* [X]t [X] */
            float
                                  \prime^* [X] t [X] \rightarrow'**matrix();
            int
                      i, j, k;void gaussj(), /* gauss-jordan elimination
                                                                                           */
                       free_matrix() ;
            A = matrix(1, order+1, 1, order+1);X = matrix(1, n, 1, order+1);printf ("ad\nu", n);for (i=1; i<=n; i++) printf ("Re\tke\n", data[i][1], data[i][2]);for (i=1; i<=n; i++) (
                       X[i][1] = 1.0;for (j=2; j<=order+1; j++)X[i][j] = X[i][j-1]*data[i][1];}
            for (k=1; k<=order+1; k++)for (j=k; j<=order+1; j++) {
                                  A[k][j] = 0.0;for (i=1; i<=n; i++)A[k][j] += X[i][k]*X[i][j];A[j][k] = A[k][j];}
            for (k=1; k<=order+1; k++) {
                       C[k][1] = 0.0;for (i=1; i<=n; i++)C[k][1] += X[i][k] *data[i] [2];
            \mathcal{F}
```

```
E-7
```
gaussj (A,order+l,C,l);

 $free_matrix(X,1,n,1,order+1);$ free\_rnatrix(A,l,order+l,l,order+l); return;

}
```
/*
                         FILE gaussj.c
                                                                            */
/*
                                                                            */
        Gauss-Jordan Elimination from Numerical Recipes in C
#include <math.h>
#define SWAP(a,b) {float temp=(a);(a)=(b);(b)=temp;}
void gaussj(a,n,b,m)
float **a,**bj
int n,mj
{
        int *indxc,*indxr,*ipiv,*ivector()j
        int i,icol,irow,j,k,l,llj
        float big, dum, pivinv;
        void nrerror(), free_ivector();
        index = \text{vector}(1, n);indxr=ivector(1,n);ipiv=ivector(l,n)j
        for (j=1; j<=n; j++) ipiv[j]=0;for (i=lii<=nji++) {
                 big=O.Oj
                 for (j=1; j<=n; j++)if (iDiv[j] != 1)for (k=lik<=nik++) {
                                          if (ipiv[k] == 0) {
                                                   if (fabs(a[j][k]) \geq big) {
                                                           big=fabs(a[j][k]);
                                                           irow=jj
                                                           icol=kj
                                                   }
                                           } else if (ipi[k] > 1)nrerror{"GAUSSJ: Singular Matrix-I"
                                  \mathcal{E}++(ipiv[icol]);
                 if (irow != icol) {
                          for (l=1; l<=n; l++) SWAP(a[irow][1], a[icol][1])
                          for (l=1; l<=m; l++) SWAP(b[irow][1],b[icol][1])
                 }
                 indxr[i]=irowj
                 index[i]=icol;if (a[icol][icol] == 0.0) nrerror("GAUSSJ: Singular Matrix-
                 pivinv=1.0/a[icol] [icol];
                 a[icol][icol]=1.0;for (l=1; l<=n; l++) a[icol] [1] *= pivinv;
                 for (l=1; l<=m; l++) b[icol][1] *= pivinv;
                 for (ll=1; ll<=n; ll++)if (11 := icol) (
                                  dum=a[11] [icol];
                                  a [11] [icol] =0.0 i
                                  for (l=1; l<=n; l++) a[ll][1] -= a[icol][1]*d
                                  for (l=1; l<=m; l++) b[ll][l] -= b[icol][l]*d
```

```
}<br>for (l=n;l>=1;l--) {
                                                                                                                                                 if (indxr[l] != indxc[l])
                                                                                                                                                                                                                    for (k=1; k<=n; k++)SWAP(a[k][indxr[1]],a[k][indxc[1]]);
                                                                                                                                                                                                                     \mathcal{L}^{\text{max}}_{\text{max}} and \mathcal{L}^{\text{max}}_{\text{max}}}
                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                                        \begin{split} \mathcal{A}^{\text{max}}_{\text{max}} & \mathcal{A}^{\text{max}}_{\text{max}} \geq \frac{1}{2} \mathcal{A}^{\text{max}}_{\text{max}} \\ & \mathcal{A}^{\text{max}}_{\text{max}} \geq \frac{1}{2} \mathcal{A}^{\text{max}}_{\text{max}} \end{split}free_ivector(ipiv,1,n);
                                                                        free_ivector(indxr, 1, n);
                                                                        free_ivector(indxc,1,n);
\mathcal{Y}\frac{1}{2} \sum_{i=1}^{n} \frac{1}{2} \sum_{j=1}^{n} \frac{1}{2} \sum_{j=1}^{n\sim \sim\label{eq:2.1} \frac{1}{\sqrt{2\pi}}\int_{0}^{\infty}\frac{1}{\sqrt{2\pi}}\left(\frac{1}{\sqrt{2\pi}}\right)^{2\pi} \frac{1}{\sqrt{2\pi}}\left(\frac{1}{\sqrt{2\pi}}\right)^{2\pi} \frac{1}{\sqrt{2\pi}}\left(\frac{1}{\sqrt{2\pi}}\right)^{2\pi} \frac{1}{\sqrt{2\pi}}\left(\frac{1}{\sqrt{2\pi}}\right)^{2\pi} \frac{1}{\sqrt{2\pi}}\left(\frac{1}{\sqrt{2\pi}}\right)^{2\pi} \frac{1}{\sqrt{2\pi}}\left(\frac{1}{\\label{eq:2.1} \mathcal{O}(\mathcal{O}(\log n)) \leq \mathcal{O}(\log n)\label{eq:2.1} \begin{split} \mathcal{L}_{\text{max}}(\mathbf{r},\mathbf{r})&=\mathcal{L}_{\text{max}}(\mathbf{r},\mathbf{r})\left(\mathbf{r},\mathbf{r}\right)\\ &\leq\mathcal{L}_{\text{max}}(\mathbf{r},\mathbf{r})\left(\mathbf{r},\mathbf{r}\right)\left(\mathbf{r},\mathbf{r}\right)\\ &\leq\mathcal{L}_{\text{max}}(\mathbf{r},\mathbf{r})\left(\mathbf{r},\mathbf{r}\right)\left(\mathbf{r},\mathbf{r}\right)\\ &\leq\mathcal{L}_{\text{max}}(\mathbf{r},\mathbf\label{eq:2.1} \mathcal{O}(\frac{1}{\sqrt{2}}\int_{\mathbb{R}^{2}}\left|\mathcal{O}(\frac{1}{\sqrt{2}}\right)\left|\mathcal{O}(\frac{1}{\sqrt{2}}\right|\right)=\frac{1}{2}\sum_{i=1}^{2}\left|\mathcal{O}(\frac{1}{\sqrt{2}}\right)\left|\mathcal{O}(\frac{1}{\sqrt{2}}\right|\right)=\frac{1}{2}\sum_{i=1}^{2}\left|\mathcal{O}(\frac{1}{\sqrt{2}}\right)\left|\mathcal{O}(\frac{1}{\sqrt{2}}\right|\right)=\frac{1}{2}\sum_{i=1}^{2}\left|\mathcal{O}(\\sim 10^{11} km
                                                                                                  \label{eq:2.1} \mathcal{L}(\mathcal{L}^{\mathcal{L}}_{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}_{\mathcal{L}})) = \mathcal{L}(\mathcal{L}^{\mathcal{L}}_{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}_{\mathcal{L}})) = \mathcal{L}(\mathcal{L}^{\mathcal{L}}_{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}_{\mathcal{L}})) = \mathcal{L}(\mathcal{L}^{\mathcal{L}}_{\mathcal{L}}(\mathcal{L}^{\mathcal{L}}_{\mathcal{L}}))\mathcal{A}_{\mathcal{A}}^{\mathcal{A}}\left(\mathcal{A}_{\mathcal{A}}^{\mathcal{A}}\right)=\mathcal{A}_{\mathcal{A}}^{\mathcal{A}}\left(\mathcal{A}_{\mathcal{A}}^{\mathcal{A}}\right)=\mathcal{A}_{\mathcal{A}}^{\mathcal{A}}\left(\mathcal{A}_{\mathcal{A}}^{\mathcal{A}}\right)=\mathcal{A}_{\mathcal{A}}^{\mathcal{A}}\left(\mathcal{A}_{\mathcal{A}}^{\mathcal{A}}\right)=\mathcal{A}_{\mathcal{A}}^{\mathcal{A}}\left(\mathcal{A}_{\mathcal{A}}^{\mathcal{A}}\right)=\mathcal{A\label{eq:2.1} \frac{1}{2} \int_{\mathbb{R}^3} \frac{1}{\sqrt{2}} \, \frac{1}{\sqrt{2}} \,\label{eq:2.1} \begin{split} &\mathcal{P}^{(1)}(\mathbf{y},\mathbf{y})\geq \mathcal{P}^{(1)}(\mathbf{y},\mathbf{y})=\mathcal{P}^{(1)}(\mathbf{y},\mathbf{y})\geq \mathcal{P}^{(1)}(\mathbf{y},\mathbf{y})=\mathcal{P}^{(1)}(\mathbf{y},\mathbf{y})=\mathcal{P}^{(1)}(\mathbf{y},\mathbf{y})=\mathcal{P}^{(1)}(\mathbf{y},\mathbf{y})=\mathcal{P}^{(1)}(\mathbf{y},\mathbf{y})=\mathcal{P}^{(1)}(\mathbf{y},\mathbf{y})=\a di Kabupatén Bandung<br>Kabupatèn Propinsi Jawa
                                                                                                                                                                                                                                                                                                                                                                  \mathcal{L}_{\mathcal{A}} and \mathcal{L}_{\mathcal{A}} are \mathcal{L}_{\mathcal{A}} . In the contribution of \mathcal{A}\langle \alpha_{\rm e}^{\rm e} \rangle\sim 10^{-1}\label{eq:2.1} \mathcal{L}=\mathcal{L}^{\frac{1}{2}}\log\left(\mathcal{L}^{\frac{1}{2}}\right) \leq \mathcal{L}^{\frac{1}{2}}\log\left(\mathcal{L}^{\frac{1}{2}}\right) \leq \mathcal{L}^{\frac{1}{2}}.\label{eq:2} \begin{split} \frac{d}{dt} \frac{d}{dt} \left( \frac{d}{dt} \right) & = \frac{d}{dt} \left( \frac{d}{dt} \right) \frac{d}{dt} \left( \frac{d}{dt} \right\sim 100\label{eq:10} \frac{\partial \rho_{\text{max}}}{\partial \rho_{\text{max}}}\left|\frac{\partial \rho_{\text{max}}}{\partial \rho_{\text{max}}}\right| = 2\pi \left(\frac{\partial \rho_{\text{max}}}{\partial \rho_{\text{max}}}\right)^2and a first companion of the state of th
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